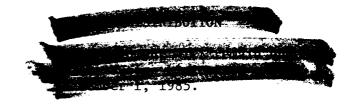
NASA Contractor Report 3818

Free-Jet Acoustic Investigation of High-Radius-Ratio Coannular Plug Nozzles

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National Aeronautics and Space Administration

Scientific and Technical Information Branch

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1.0 SUMMARY

This report, along with the companion Comprehensive Data Report, R81AEG212, summarizes the experimental and analytical results of a scale-model free-jet acoustic exploratory program performed by the General Electric Company under NASA-Lewis Research Center sponsorship on unsuppressed high-radius-ratio coannular plug nozzles with inverted velocity profiles. The nozzles selected for test were nozzles covering the range of geometry and class applicable to General Electric designs for dual-flow exhaust nozzles typical of a variable cycle engine (VCE) for advanced supersonic technology (AST).

In all, six high-radius ratio coannular plug nozzle models, along with a baseline conical nozzle, were tested for simulated flight acoustic evaluation in the General Electric Anechoic Free-Jet Acoustic Test Facility. The models tested were primarily at an outer nozzle radius ratio, $R_{\rm r}^{\rm o}$, of 0.853 and an inner-to-outer nozzle area ratio, $A_{\rm r}$, of 0.2. A model with $R_{\rm r}^{\rm o}=0.902$ and $A_{\rm r}=0.53$ was tested also. Some of the key geometry features studied were the influence of nozzle exhaust struts, a convergent-divergent flowpath on shock control, as well as effect of area ratio holding radius ratio fixed, and the effect of radius ratio holding area ratio fixed.

Some of the key results of this investigation were:

- In simulated flight, the high-radius-ratio coannular plug nozzle essentially maintained its jet noise and shock noise reduction feature observed under static conditions relative to a baseline conical nozzle.
- The presence of nozzle bypass struts will not significantly effect the acoustic characteristic of a General Electric-type nozzle design.
- A unique coannular plug nozzle spectral prediction method was evolved based on modern acoustic theories and significant static and simulated flight acoustic test results.
- Diagnostic acoustic and laser velocimeter tests were performed which led to observations regarding possible regions of the flow in which further coannular plus nozzle shock control research could evolve reduction of shock-cell noise.

2.0 INTRODUCTION

The General Electric Company under NASA Lewis Contract NAS3-18008 initiated an exploratory scale-model acoustic and aerodynamic performance test program to obtain parametric data on unsuppressed and suppressed coannular nozzles. That program was directed toward the development of high velocity jet noise technology for Advanced Supersonic Transport application. One of the findings of that program, results of which are in Reference 1, was that the unsuppressed coannular plug nozzle exhibited acoustic benefits with modest performance losses. A follow-on investigation was conducted under Contract NAS3-19777 (Reference 2) with the objective to determine the effects of key design variables of unsuppressed coannular plug nozzles through a systematic static acoustic and wind tunnel aerodynamic performance measurements. The variables considered were radius ratio, area ratio, inner stream plug geometry, inner and outer stream flow variables, and inner to outer stream velocity and weight-flow ratios. The measured data identified the mixed stream velocity V $_{\rm j}^{\rm mix}$, outer stream radius ratio $R_{\rm r}^{\rm o}$, inner-to-outer velocity ratio V $_{\rm r}$, and inner-to-outer stream area ratio Ar as the parameters that had influence on the measured jet noise data. The current study was initiated with the objective to confirm the observed coannular nozzle acoustic benefits under simulated flight conditions. In addition, effort has been made to develop a semiempirical spectral prediction method for coannular plug nozzles that will take into consideration the various noise generating mechanisms that have been identified in relevant studies (References 1-7).

To determine the effect of forward flight on the acoustic effectiveness of various coannular plug configurations, six coannular model nozzles along with a reference circular conic nozzle were tested under both static and simulated flight conditions. In addition, the influences of nozzle exhaust struts, area ratio and radius ratio on jet noise, and a convergent - divergent flowpath on shock-cell noise were investigated. The aerodynamic flow conditions for the coannular plug nozzle test points were selected to simulate a typical AST/VCE operating line and to yield an inverted velocity profile. Furthermore, the laser velocimeter was used with five of the test configurations to determine the jet plume mean and turbulent velocity distributions and to correalte these data with the acoustic results. The details of the configurations and scope of testing are summarized in Section 4.0 and the measured acoustic and LV data are presented and discussed in Section 5.0. Detailed acoustic and laser-velocimeter data are presented separately in the Comprehensive Data Report (Reference 8) of this program.

In addition to the relevance of the measured acoustic results in determining the effects of tested parameters on jet and shock-cell related noise of coannular plug nozzles, the data of this program were employed during the concept screening related to the selection of engine scale hardware for the YJ101/VCE tests. These engine tests are being conducted as a part of a multiphase, multiyear GE/NASA test-bed engine program to investigate key technology features applicable to an AST powerplant. Furthermore, selected model-nozzle

data, obtained during the tests are to be compared with the test-bed engine data in order to verify coannular nozzle acoustic data scaling procedures. Preliminary correlation of the full-scale engine and scale-model acoustic results and verification of the analytical prediction methodology developed during this program are reported in Reference 9.

3.0 TEST APPARATUS AND DATA REDUCTION PROCEDURES

All of the acoustic and laser velocimeter tests of this program were performed in the General Electric Anechoic Test Facility located at Evendale, Ohio. A brief description of the test apparatus, the data acquisition, and the data reduction procedures is presented in this section. Results of the tests conducted to determine (1) the acoustic characteristics of the anechoic chamber (e.g., inverse square law tests, background noise determination) and (2) the mean velocity and turbulence intensity distributions in the free jet, along with a detailed description of the aerodynamic/acoustic data acquisition and reduction systems are presented in the Comprehensive Data Report of this program (Reference 8).

3.1 GENERAL ELECTRIC ANECHOIC JET NOISE FACILITY

3.1.1 General Arrangement and Operational Range

The test facility, schematically and photographically shown in Figure 1, is a cylindrical chamber having a diameter of 13.1 m (43 ft) and a height of 21.95 m (72 ft). The inner surfaces of the chamber are lined with anechoic wedges made of fiberglass to yield a low frequency cutoff below 220 Hz and an absorption coefficient of 0.99 above 220 Hz.

The facility can accommodate model configurations up to 17.3 cm (6.8 in.) in diameter. The operating domains of this facility in terms of total temperature, pressure ratio, and jet velocity are indicated in Figure 2 for single and dual flow operation. The required streams of heated air that are produced by separate burners pass through acoustically treated plenum chambers for the suppression of flow and combustor noise.

A tertiary duct surrounds the model nozzles with the airflow in order to simulate a forward flight up to Mach 0.41. The tertiary air passes through a silencer plenum chamber before it is discharged through the 1.2 m (48 in.) free-jet exhaust. An overhead view of the tertiary exhaust surrounding a test conical nozzle is presented in Figure 3.

3.1.2 Acoustic Data Acquisition and Reduction

3.1.2.1 Acoustic Data Acquisition System

A schematic of the microphone data acquisition system used to obtain acoustic data during tests in the anechoic chamber is shown on Figure 4. This system is optimized for obtaining acoustic data up through 80 kHz 1/3-octave center frequency. The microphones used to obtain the far field data are the B&K 4135 6.4 mm (0.25 in.) condenser microphones. All the tests are conducted with the microphone grid caps removed to obtain the best frequency



ACOUSTIC WEDGES

0 110° 100°

COANNULAR MODEL NOZZLE

LASER TRACK

FREE JET NOZZLE (4 FT. DIA.) AIR SILENCER INLET

ENTRAINED FLOW

MODEL NOZZLE FLOW

00/

09

FREE JET

008

SCREEN AND HONEY-COMB FLOW STRAIGHTNERS

WEDGES AROUND FREE JET

006

ANECHOIC CHAMBER

(40' DIA x 73' HIGH)

MICROPHONE LOCATIONS (13)

MMMMMM

25.63 26.65 30.13 35.25 31.18 28.73 27.42 27.00 27.42

160 1150 1140 1130 1120 1100 90 80 70 60 60 60

- EXHAUST SILENCER

Radial Distance Fr.

Angle Deg. FREE JET PLUME (b) Photo

(a) Schematic

Figure 1. Free-Jet Arrangement in Anechoic Facility.

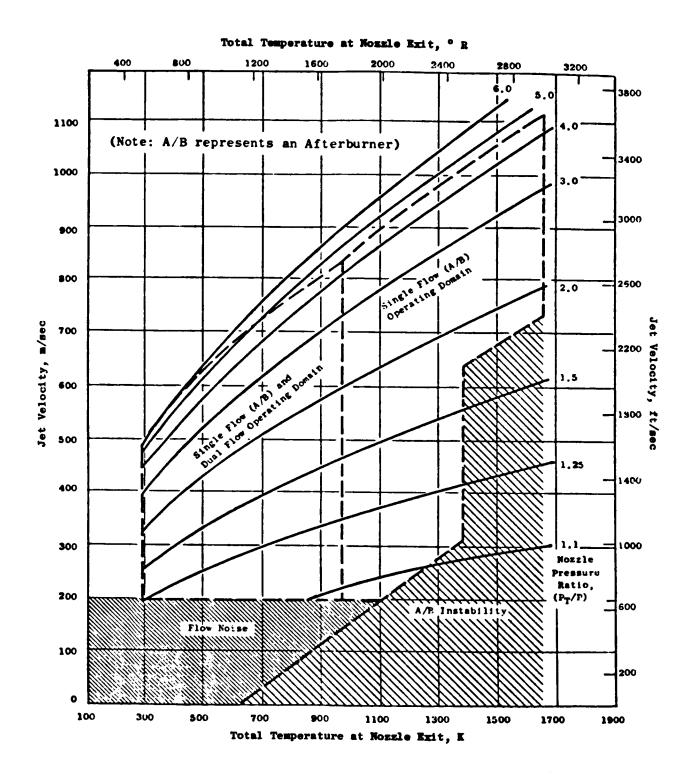


Figure 2. General Electric Anechoic Chamber Operating Domain.

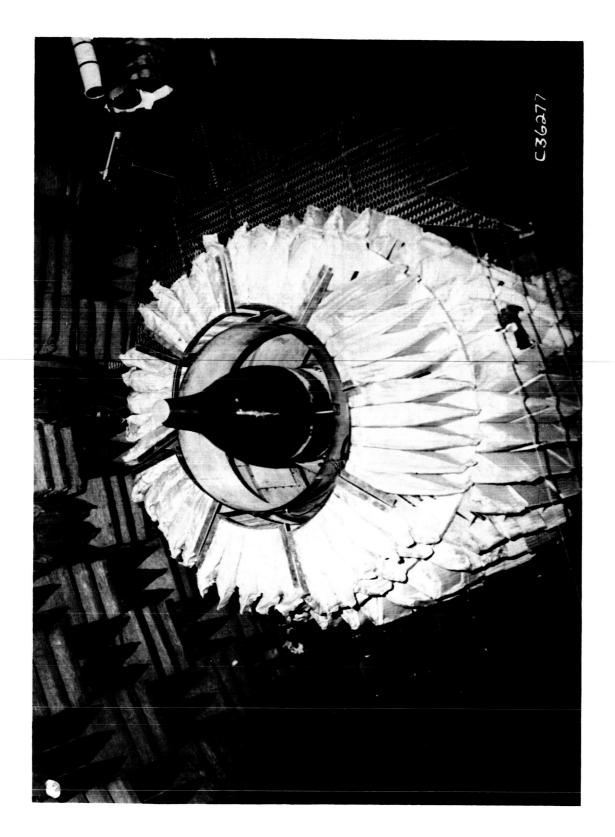


Figure 3. Overhead View of the Tertiary Exhaust with a Test Conical Nozzle.

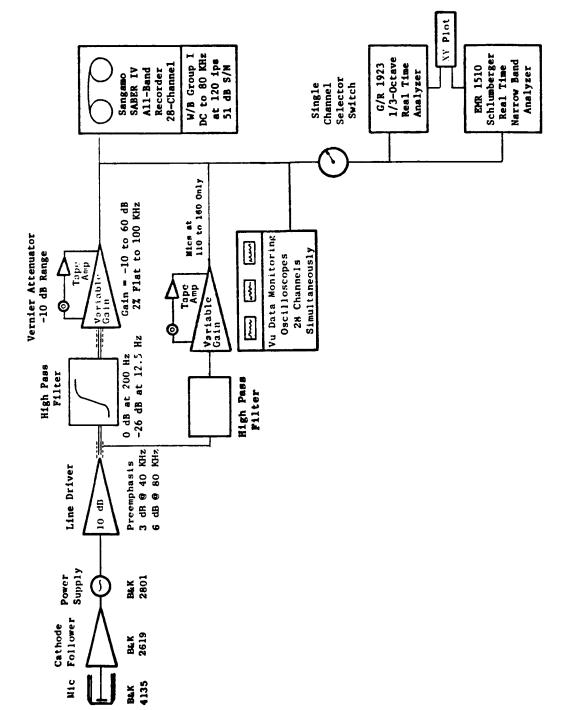


Figure 4. Acoustic Data Acquisition System.

response. The cathode followers are the transistorized B&K 2619 for optimum frequency response and lower inherent system noise characteristics relative to the 2615 cathode follower. All systems utilize the B&K 2801 power supply operated in the direct mode.

Power supply output is connected to a line driver adding 10 dB of amplification to the signal, as well as adding "preemphasis" to the high frequency portion of the spectrum. The net effect of this amplifier is a 10 dB gain at all frequencies, plus an additional 3 dB at 40 kHz and 6 dB at 80 kHz due to preemphasis. This procedure improves low amplitude, high frequency data. In order to remove low frequency noise, high pass filters with attenuations of approximately 26 dB at 12.5 Hz and decreasing to 0 dB at 200 Hz are installed in the system.

The tape recorder amplifiers have a variable gain from -10 dB to +60 dB in 10-dB steps and a gain trim capability for normalizing incoming signals. The prime system used for recording acoustic data is a Sangamo/Sabre IV, 28-track FM recorder. The system is set up for Wideband Group I (intermediate band double extended) at 120-in./sec tape speed. Operating at this speed provides a dynamic range that is necessary for obtaining satisfactory low-amplitude, high frequency acoustic signal. The tape recorder is set up for ±40% carrier deviation with a recording level of 8 volts peak-to-peak. During recording, the signal is displayed on a calibrated master oscilloscope, and the signal gain is adjusted to maximum without exceeding the 8-volt peak-to-peak level.

High pass filters are incorporated into the acoustic data acquisition systems to enhance the high-frequency data previously lost in the tape recorder electronic noise floor for microphones from 110° to 160°. The microphone signal below the 20-kHz 1/3-octave band is filtered out, and the gain is increased to boost the signal to noise ratio. For microphones from 110° to 160°, both filtered and unfiltered signals are recorded on tape. For data below 20 kHz, the unfiltered signal is used to calculate the sound pressure levels, while the filtered signal is employed for high frequencies. The entire jet noise spectrum at a given angle is obtained by computationally merging these two spectra.

3.1.2.2 Acoustic Data Reduction

Standard data reduction is conducted in the General Electric-AEBG Instrumentation and Data Room (IDR). As shown in Figure 5, the data tapes are played back on a CEC3700B tape deck with electronics capable of reproducing signal characteristics within the specifications indicated for Wideband Group I. An automatic shuttling control is included in the system. In normal operation, a tone is inserted on the recorder in the time slot designed for data analysis. The tape control automatically shuttles the tape, initiating an integration start signal to the analyzer at the tone as the tape moves in its forward motion. This motion continues until an "integration complete" signal is received from the analyzer at which time the tape direction is reversed;

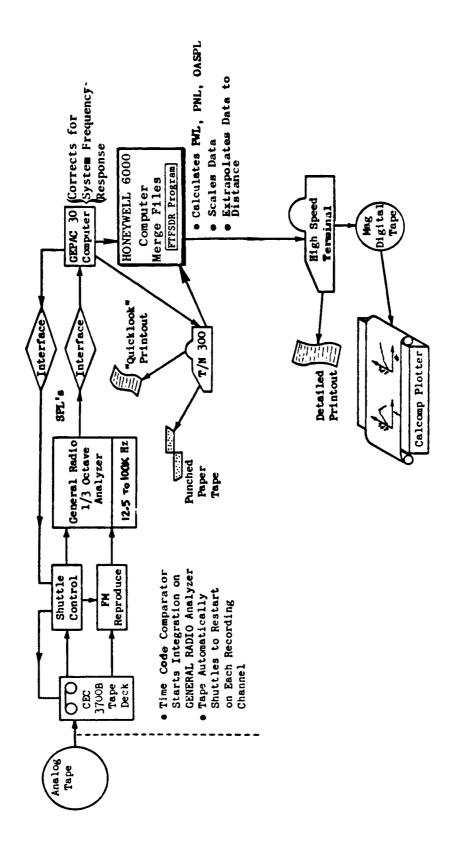


Figure 5. Acoustic Data Reduction System.

and at the tone, the tape restarts in the forward direction advancing to the next channel to be analyzed until all the channels have been processed. In addition, a time code generator is utilized to signal tape position as directed by the computer program control.

All 1/3-octave analyses are performed on a General Radio 1921 analyzer. Normal integration time is set for 32 seconds to ensure sufficient averaging time for the low frequency content. The analyzer has 1/3-octave filter sets from 12.5 Hz to 100 Hz, and has a rated accuracy of ±1/4 dB in each band. Each data channel is passed through an interface to the GEPAC 30 computer where the data are corrected for the frequency response of the microphone and the data acquisition system. Next, the data are corrected to standard day (15° C, 70% RH atmospheric attenuation conditions as recommended by Shields and Bass, Reference 10) and processed to calculate the perceived noise level (PNL) and the overall sound pressure level (OASPL) from the spectra. For calculation of acoustic power, scaling to other nozzle sizes, and/or extrapolation to different far field distances, the data are sent to the Honeywell 6000 computer for data processing. This is accomplished by transmitting the SPL via direct timeshare link to the 6000 computer through a 1200 Band Modem. In the 6000 computer, the data are processed through the Flight Transformed Full-Scale Data Reduction (FTFSDR) Program where the appropriate calculations are performed. The data printout is accomplished on a high-speed "remote" terminal. The FTFSDR Program also writes a magnetic tape for CALCOMP plotting of the data. Detailed descriptions of the acoustic data reduction system are given in the Comprehensive Data Report (Reference 8).

3.2 GENERAL ELECTRIC LASER VELOCIMETER

3.2.1 General Arrangement

The laser velocimeter (LV) used is a system developed under a USAF/DOT-sponsored program and reported in detail in References 11 and 12. The basic optical system is a differential Doppler, backscatter, single package arrangement that has the proven feature of ruggedness for the severe environments encountered in high velocity and high temperature jets. Figure 6 shows a photograph of the LV system installed in the General Electric Anechoic Jet Noise Facility and Figure 7 indicates a schematic arrangement of the laser package. The laser beams are projected from below the lens, forming an angle that keeps the major axis of the control volume ellipsoid to a minimum. The dimensions of the control volume are 0.535 cm (0.21 inch) for the major axis and 0.0518 cm (0.020 in.) for the minor axis. The range of the LV control volume from the laser hardware is 2.16 m (85 in.). The three steering mirrors and the beam splitter are mounted on adjustable supports made from the same aluminum alloy in order to eliminate temperature alignment problems.

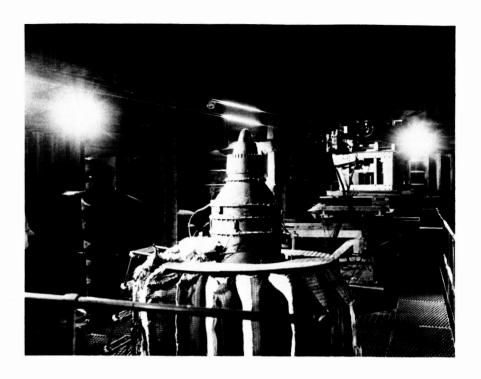


Figure 6. Two-Laser System in the GE Anechoic Jet Noise Test Facility.

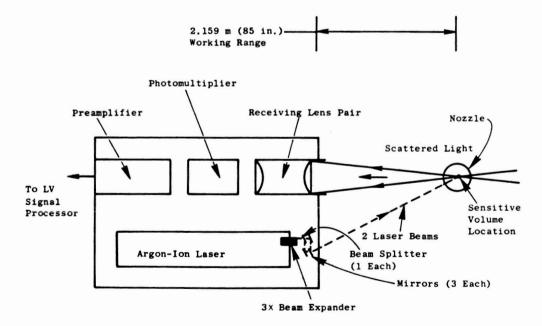


Figure 7. Laser Velocimeter Optics Package.

3.2.2 LV Actuator and Seeding

The remotely actuated platform has vertical, horizontal, and axial travel capabilities of 0.813 m (32 in.), 0.813 m (32 in.) and 5.79 m (228 in.), respectively. The resolution is ± 0.1588 cm (0.0625 in.) for each axis except for the last 5.28 m (208 in.) of axial travel, which has a resolution of ± 0.3175 cm (0.125 in.).

Seeding is by injection of nominal $1-\mu m$ diameter aluminum oxide (Al₂0₃) powder into the air supply to the burners and to the exit region of the tertiary duct so as to seed the entrained air. The seeding equipment used is described in Reference 11 (Chapter V, Section 3). However, the air supply to the fluidized bed column is currently heated to about 394° K (250° F) to prevent powder aggregation by moisture absorption.

3.2.3 Signal Processing and Recording

The LV signal processor is a direct-counter (time-domain) type similar to that reported in Reference 11, but with improvements. These improvements result in a lowered rate of false validations and improved linearity and resolution. Turbulent velocity probability distributions (histograms) are recorded by a NS633 pulse-height analyzer with 256 channels. The data acquired from the LV are transmitted to a minicomputer system (PDP 11/45) for storage on disk/tape and data reduction.

The processing capabilities of the General Electric LV system are as follows:

- Velocity range 35 to 5000 ft/sec
- Random error for single particle accuracy (error associated with system inaccuracies such as fringe spacing, linearity, stability, burst noise) - 0.75%
- Bias error for mean velocity 0.5%
- False data rejection capability (possibility of accepting bad data) - 0.0002%

The GE system uses a 16-fringe control volume where all of the 8 center fringes are used in the data acceptance/rejection testing.

3.2.4 Laser Velocimeter Data Reduction and Typical Test Results

The concept of using LV measurements for obtaining routine mean and turbulent velocity profiles can be described as follows. Two beams of monochromatic light intersect at a point in space and set up a fringe pattern of

known spacing. The flow is seeded with small particles which scatter the light while passing through the control volume. The scattered light is collected and the laser signal processor determines the time used by the particles to pass through each fringe. A knowledge of the distance and the time used by each validated particle enables the construction of the usual histogram. Then, by statistical techniques, the mean value (corresponding to the mean velocity) and the standard deviation (corresponding to the turbulent velocity) are constructed.

As with any large number of data samples, guidelines for estimating the accuracy of the measured delay are needed. Tables I and II, respectively, provide estimates of the percent of error associated in measuring the mean and turbulent velocities for a 95% confidence level as a function of the number of data samples and level of turbulence. Between 2000 to 5000 data samples are taken during a routine measurement. For simple and quick diagnostictype information, this number of samples is sufficient.

Although the principle of LV measurement is well known, the practical aspects of designing a reliable electronic processing unit in order to monitor valid particles are arduous. Earlier investigators have had great difficulty in performing measurements even in low velocity jets, and hence the successful use of a LV system in the heated supersonic jets of this program represents a major achievement. References 13 through 16 list some of the reference materials in which General Electric has demonstrated the capabilities of the LV system for measuring mean and turbulent velocities in high temperature jet exhaust plumes of supersonic conical and coannular nozzles.

3.3 DESCRIPTION OF THE FLIGHT TRANSFORMATION TECHNIQUE

3.3.1 Objective and Concept

The objective of the General Electric free-jet transformation process is to employ far field SPL spectra at various angles to the jet axis (typically) for $20^{\circ} \leq \theta_{\rm I} \leq 150^{\circ}$ in increments of 10°) obtained in a free-jet experiment and transform them to yield SPL spectra as would be measured in a true moving frame experiment.

The concept employed is that with area ratios of 50:l or so and with the primary nozzle exhaust plane displaced aft of the free-jet plane sufficiently to permit aquisition of acoustic data in the inlet arc (e.g. up to $\theta_{\rm I}$ = 150°), proper aerodynamic simulation of the effects of forward flight is achieved; but that in terms of the acoustic simulation of the effects of uniform flow over the primary jet plume noise sources, the free jet achieves this only to a limited extent. In other words, the free jet achieves the effect of the right source mix but radiating into an environment that more nearly approaches a static environment than the environment of sources shrouded by either a finite or infinite extent of uniform nonturbulent flow. (The basis of several previous investigations has been to assuem that a well-defined region of uniform, nonturbulent flow surrounds the sources. This well-defined region is taken as

Table I. Estimated Percent Error in the LV Measurement of Mean Velocity with 95% Confidence.

Number of Data Samples		۷'/	v,	
	0.2	0.1	0.05	0.025
10	14.1	7.0	3.5	1.76
20	9.3	4.7	2.3	1.20
30	7.4	3.7	1.9	0.93
40	6.3	3.2	1.6	0.80
60	5.0	2.6	1.3	0.65
120	3.6	1.8	0.9	0.45

Table II. Estimated Percent Error for LV Turbulent Velocity Measurements With 95% Confidence.

Number of Data Samples	Percent Error
20	31.5
40	21.8
60	17.8
120	12.6
240	9.12
480	6.45
960	4.56
5,000	2.0
25,000	0.89

a doubly infinite cylinder of constant circular section equal to the cross section of the free-jet exhaust plane.) The acoustic sources in a free jet, of course, do not radiate into a completely static environment and hence some propagation effects of the free-jet flow do have to be accounted for.

Based on the above picture, the broad outline of the procedure adopted is as follows. Defining as the "static" directivity, the directivity pattern (in various frequency bands) that the sources (of the primary jet exhaust plume altered by the effects of relative velocity due to imposition of the free jet) may be expected to produce if they radiated into a quiescent environment, we first deduce this static directivity from the measured free-jet experimental data by correcting the latter for propagation effects of the free jet. Since the free-jet flow field includes intensely turbulent shear layers through which the sound field of the sources must pass before it reaches the far field microphones (located in the quiescent ambient), some degree of empiricism (especially for the high frequency sound) is involved in attempting to account for these propagation effects.

Once such a static directivity is extracted, it still remains to deduce what the noise signature of the source distribution would be if the source distribution was not stationary relative to the ambient but moving relative to the ambient at the flight velocity. A multiple decomposition procedure suitable for the broad band jet noise problem which attempts to synthesize the static directivity by ascribing it to a mix of uncorrelated singularities was developed in order to enable the prediction of the flight noise. Once such a decomposition is completed, we simply apply the dynamic exponent applicable to each singularity to derive the flight noise signature.

In summary, the method starts with directivities from the free-jet experiment in various third octave bands, corrects these directivities for free-jet propagation effects in a frequency-dependent manner to retrieve the static directivity, synthesizes the static directivity by a suitable mix of uncorrelated singularities, and finally applies the dynamic effect appropriate to each singularity to predict the flight noise. It is an inherent feature of the method that it works separately with each third octave band directivity pattern. The final flight predictions can then be summed to yield either OASPL or PNL directivities or simply displayed as flight SPL spectra at various angles to the jet axis. (Doppler shift effects on the frequency are fully accounted for.)

3.3.2 Algorithm Description

A detailed algorithm description is shown in Figure 8 along with the applicable nomenclature. Complete description and discussion of this procedure can be found in Reference 17.

3.3.3 <u>Further Details</u>

The recommended procedure for transformation of free-jet noise to flight noise consists of extracting the "basic" directivity from the measured free-jet

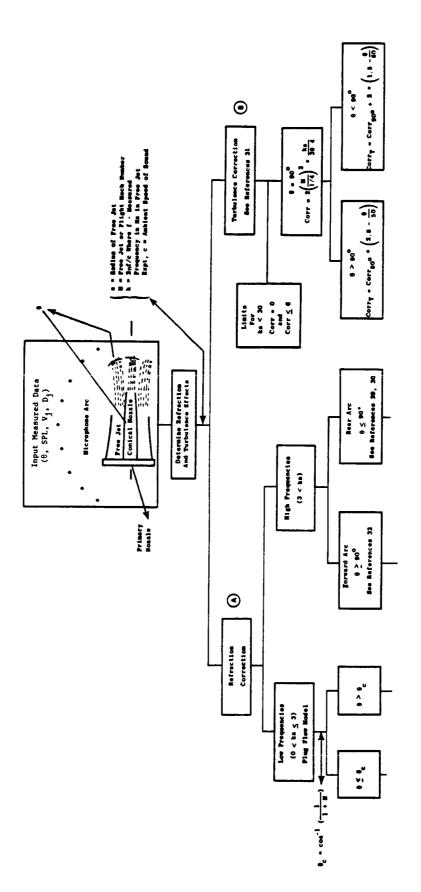


Figure 8. Algorithm Description.

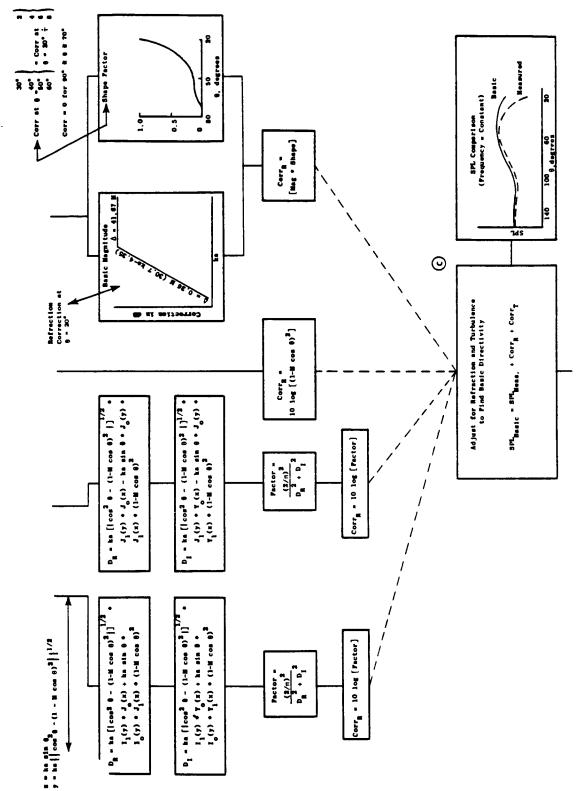


Figure 8. Algorithm Description (Continued).

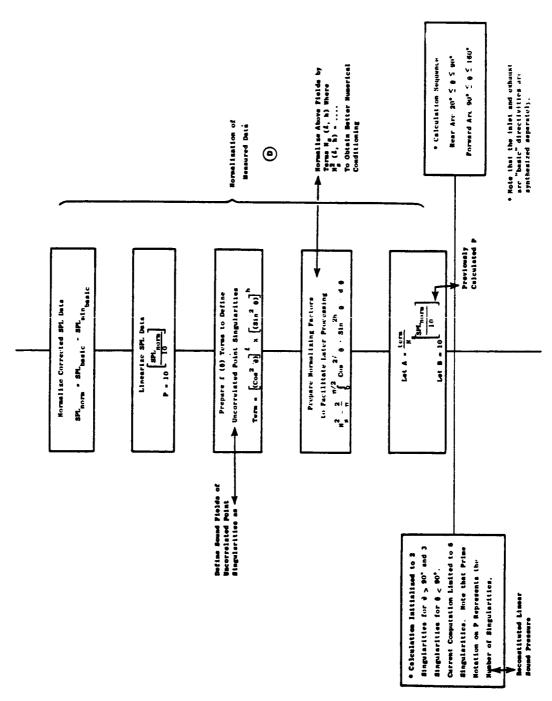


Figure 8. Algorithm Description (Continued),

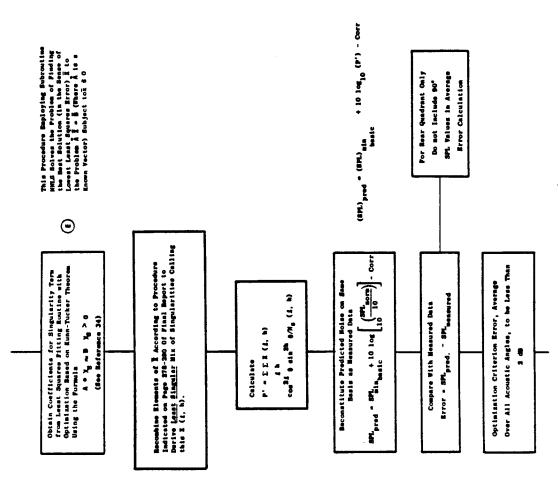
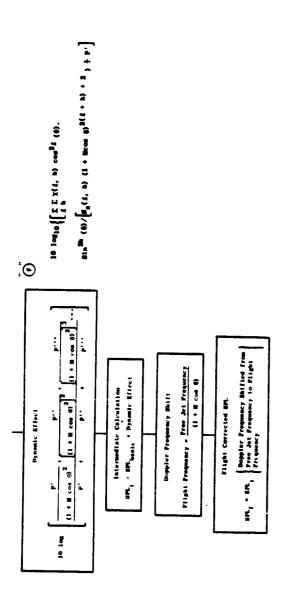


Figure 8. Algorithm Description (Continued).



List of Symbols

Numerical Value Varies with Level of Singularity Considered Mach Number n (Free Jet Velecity) + (Ambient Speed of Sound) Hormalising Pactor, Panction of Singularity Linear Sound Presente	0.0001 Nicrobara Subseriat for Defraction Comments	Singularity Subscript	Subscript for Turbulence Correction	Acronys Used to Identify a Unique Algebraic Grouping	Bessel Punction Argument, x . hs aim 0	A Vector Derived from Least Squarus Pitting, Punction of Singularity	Bessel Puzztion Argument, y a ha cos 0 = (1 - H cos 0) 2 Bessel Puzztion of the Second Kind of Order m, Argument x	Angle from the Jet Azis Referred to the Exhaust Critical Angle that Defines the Jet Zone of Silence .	con (1/1 + H)	
7 # 5 A	Ĭ.		•	Term	= 1		, <u>, , , , , , , , , , , , , , , , , , </u>	• •		
Am Imput Matrix to the Loast Squares Pitting Procedure Am Imput Vector to the Loast Squares Pitting Procedure Am Acrespa used to identify the Mefrection and/or Turbulence Correction	Beal Boot of Demainator Term in Solution of the Sound Prasoure for the Plux Flow Bodel	Inselbary Seet of Denominator Term in Solution of Ihr Sound Pressury for the Flow Hodel	Subscript for Flight Currected SPL	Mmerical Value Varies with Level of Singularity Considered	Subscript on SPI, to Identify un Intermediate Calculation	Modified Bessel Purchion of the First Kind of Order m. Argument in is a	Modified Bessel Punction of the Pirst Eind of Order m. Argument in in in y	Beneal function of the First Kind of Order α_s Argument in is α_s	Reseal Panelion of the Pirsi Eind of Order n, Argument in 10 γ	Frequency Parimates . (Free Jet Prequency Band of Interest in Redishs per Second) : (Free Jet Redius in Feet) ; (Ambiest Speed of Second)
• • •	- *	٠-			-	(E)	3.	÷.	£.	:

Figure 8. Algorithm Description (Concluded).

data and then applying the "dynamic" effects to determine the noise in flight. The basic directivity is the directivity that the sources associated with the primary nozzle plume would create, if they radiated into a static rather than the free-jet environment.

Two phenomena are involved that change the directivity of the noise radiated by the sources associated with the jet plume when the jet is exhausting into a free-jet environment as opposed to a static environment. These are:

a. Refractive Effects of the Free-Jet Flows

To deduce the refractive effects of the free-jet flow, the following procedure is adopted:

1. At low frequencies $(k_0c < 3)$, the plug flow model solution for a point pressure source is used

$$p' \propto (1 - M_{fj} \cos \theta)^{-2}$$

2. At high frequencies $(k_0c > 3)$, the asymptotic high frequency solution for a pressure source is used

$$p' \simeq (1 - M_{fj} \cos \theta)^{-1}$$

At these values of the frequency parameter ($k_{O}c < 3$), the exhaust arc was used to deduce the refractive effect following the method due to Schubert (Reference 18). In this method:

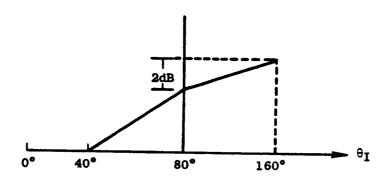
- First, the refractive dip in dB along the jet exhaust axis is determined as being product of the jet Mach number and the frequency parameter.
- Then, a shape factor that is essentially Mach number and frequency indpendent is used to determine the refractive dip at other angles. For the range $3 < k_{o}c < 6$, Ribner's results were used with a linear extrapolation in the range $6 > k_{o}c > 1.25$.
- Based on experimental data, the refractive dip in the exhaust arc for $k_0c > 6$ was considered independent of k_0c , but still linearly proportional to $M_{\mbox{fi}}$.
- b. Absorptive Effects of the Fine Grain Turbulence in the Shear Layer of the Free-Jet

This relates to the fact that fine-grained turbulence in the shear layer of the free jet can absorb sound, expecially at high frequencies. This correction is based on Crow's theory that states that the effective absorption

coefficient is proportional to the frequency, distance the sound traveled in the shear layer, and the square of the Mach number.

^aabsorption coefficient
$$_{\alpha}$$
 f $_{\rm fjl}^2$

Based on the path length that the sound has to traverse, the absorption coefficient is assumed to vary with θ I as shown in the following sketch:



The absorption was calculated assuming an eddy viscosity

$$\epsilon_{\rm eddy\ viscosity}$$
 = 70 μ for $M_{\rm fj}$ = 0.25 and f = 50 kHz.

This yields corrections for $k_{\text{O}}c > 30$. The actual expressions used were

where M = free-jet Mach number

$$k = \frac{2 \pi f}{c_a}$$

a = Radius of the free jet

$$Corr_T \mid_{\theta_I} > 90^{\circ}$$
 = $Corr_T \mid_{\theta_I} = 90^{\circ} \times (1.5 - \frac{180 - \theta_I}{180})$

$$\operatorname{Corr}_{\mathbf{T}} \left| \begin{array}{c} \theta_{\mathbf{I}} > 90^{\circ} \end{array} \right| = \operatorname{Corr}_{\mathbf{T}} \left| \begin{array}{c} \theta_{\mathbf{I}} = 90^{\circ} \end{array} \right| \times (2.8 - \frac{180 - \theta_{\mathbf{I}}}{50})$$

From the measured free-jet data, the refraction and turbulences absorption corrections are added to obtain the basic directivity of the sources.

The basic directivity obtained above is assumed to be generated by a set of singularities F_0 , F_x , F_v , etc., such that the sound field is a solution to

$$\nabla^2_p + k_0^2_p = F_0 \delta(x) \delta(y) \delta(z) + F\alpha'(x) \delta(y) \delta(z) + Fy \delta(X) \delta(y) \delta(z)$$

where F_0 , F_x , F_y are mutually uncorrelated, so that they contribute to the far field only additively. As the mean square pressure of any singularity is symmetric about both $\theta = 0^\circ$ and $\theta = 90^\circ$, the inlet and exhaust arc are synthesized separately.

The procedure adopted to determine the dynamic effect is as follows:

- 1. From the basic directivity pattern, obtain the normalized SPL's based on the least singular fit in both the forward and aft quadrants.
- 2. Determine the linearized levels using the equation

$$\frac{1}{p^2} = 10 \frac{SPL - SPL_{min}}{10}$$

- 3. Decide on a level of fitting using the criterion that the data ought to be reconstructed to within an error of 2 dB on the average. Then, assuming that the data ought to be reconstructed with the least singular distribution of uncorrelated sources possible, the problem simplifies to one of solving a least squares problem of the type find \bar{x} to minimize $|\bar{r}| = (A\bar{x} b)$ subject to nonnegative constraining $\bar{x} \ge 0$. This is done using an algorithm based on the Kuhn-Tucker theroem of optimization theory.
- 4. The singularities obtained using the Kuhn-Tucker theorem are then combined to obtain the least singular decomposition of the sources.

5. The appropriate dynamic effect is then applied to each singularity type to determine the correction that is applied to the measured free-jet data corrected for refraction and turbulence absorption. If the mean square of the sound pressure is obtained by adding the singularities as

$$P_{\theta}^{'2} = F_{o} c^{6} + F_{1}c^{4} s^{2} + F_{2} c^{2} s^{4} + F_{3} s^{4}$$

where $C = \cos \theta$

 $S = Sin \theta$

the dynamic effect is calculated using the relation

Dynamic Effect = 10
$$\log_{10} \frac{P_F^{'2}}{P_S^{'2}} = 10 \log_{10} \frac{\frac{F_0C^6}{k^8} + \frac{F_1C^4S^2}{k^8} + \frac{F_2C^2S^4}{k^8} + \frac{F_3S^4}{k^6}}{P_S^{12}}$$

where $k = (1 + M_{fj} \cos \theta)$

6. The levels are then corrected to:

7. Doppler frequency shift results in the flight frequency given by:

$$f_F = \frac{f_j}{1 + M_{f_i} \cos \theta}$$

8. Hence $SPL_F = SPL_i$ Doppler shifted from free jet to flight.

Thus, using the above transformation, the free-jet data can be transformed into flight data. Further discussion of this procedure is found in Reference 17.

4.0 CONFIGURATION DESCRIPTION AND SCOPE OF TESTING

To determine the effect of forward flight on the acoustic effectiveness of various coanular plug nozzle configurations, six coannular plug nozzles along with a reference conical nozzle were tested during this program in the General Electric anechoic facility. Furthermore, the laser velocimeter system was used with five of these configurations to determine the jet plume mean velocity and turbulence intensity distributions and to correlate these data with the acoustic results. The details of the nozzle configurations and the aerodynamic flow conditions of the tests are presented in this section.

4.1 CONFIGURATION DESCRIPTION

Table III summarizes the significant geometric parameters of the test configurations. Models 1A through 4 basically are geometrically scaled versions of the General Electric VCE/AST test-bed nozzle designs. In particular, Models 1A and 2 (both having a C-D outer termination and an inner-to-outer area ratio $A_{\rm r}=0.2$) are identical in all respects except that Model 1A has eight internal struts similar to the VCE/AST test-bed design while Model 2 has no struts. Model 3 ($A_{\rm r}=0.2$) is similar to Model 2 except that it has a convergent termination in the outer stream. Finally, Model 4 with an area ratio $A_{\rm r}=0.53$ is obtained from Model 3 by removing a spacer. The design features of these four nozzles are such that (1) the influence of the eight internal struts in the outer stream on the far field acoustic measurements, (2) the influence of a convergent and C-D terminated outer nozzle on the measured shock noise, and (3) the effect of the area ratio (test range = 0.2 + 0.53 with an outer stream radius ratio $R_{\rm r}^{\rm O}=0.853$) on the acoustic effectiveness of the VCE/AST design can be determined.

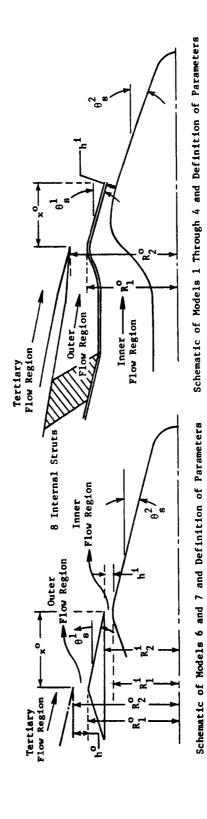
The Model 5 configuration is a single stream conical reference nozzle. The remaining nozzle configurations, namely, Models 6 and 7, were selected for flight simulated tests for the following reasons:

- They were tested earlier under static conditions (Reference 2, referred to therein as Configurations 6 and 7) over a wide range of aerodynamic flow conditions and, hence, a large data base is available on these two models for a comparison with the corresponding flight data from this study.
- These nozzles have a small area ratio and a large outer stream (Reference 2) radius ratio. The static tests have indicated that these geometric features are key coannular nozzle suppression parameters.

Detailed drawings of the test hardware are presented in the Comprehensive Data Report of this contract (Reference 8).

Table III. Significant Geometric Details of Test Configurations.

Model	h°, in.	h ⁱ , in.	Model h ⁰ , in. h ¹ , in. R ₁ , in.	R ⁰ , in.	R ₁ , in.	R2, in.	్జ	A°, in. 2	R_2^0 , in. R_1^1 , in. R_2^1 , in. R_2^0 A°, in. R_2^0 A°, in. R_2^0 , in. R_2^0 , in. R_2^0 , in. R_2^0 in. R_2^0 in.	A ⁱ /A ^o	x°, fn.	x°/hº	-t- -	9 ₂	Internal Struts	Term Outer	Termination Outer Inner
41		0.675 0.200	3.918	4.593		i	0.853	18.049	18.049 3.610	0.20	3.417 5.06	90.5	52	13.	Yes	Q-D	Convergent
2	0.675	0.200	3.918	4.593	-	i	0.853	18.049	3.610	0.20	3.417 5.06		12.	.51	2	G-S	Convergent
c	0.675	0.200	3.918	4.593	1		0.853	18.049	3.610	0.20	3.971	5.88	13.	15.	8	Conv.	Convergent
4	0.675	0.583	3.918	4.593		i	0.853	18.049	9.566	0.53	3.971	5.88	.51	15.	9	Conv.	Convergent
•		ł	ŀ	2.547	1	1	1	20.380	-	i	ł	1	i	i	No	Conv.	Convergent
•	0.426	0.311	3.918	4.344	2.858	3.168	0.902	11.057	5.878	0.53	3.089	7.25	.11	15.	No	Conv.	Convergent
1	0.675	0.311	3.918	4.593	2.858	3.168	0.853	18.049	4.878	0.33	3.063	4.54	:	.51	£	Conv.	Convergent



4.2 SCOPE OF TESTING

The aerodynamic flow conditions for the coannular plug nozzle test points were selected to simulate a typical VCE/AST operating line and to yield an inverted velocity profile wherein the outer annulus flow is at a higher velocity and temperature than the inner stream. To the extent possible, identical inner and outer stream conditions were set during those tests that involved determining the effect of different geometries and velocity ratios on the acoustic characteristics of the nozzles.

4.2.1 Acoustic Tests

The total number of static and flight acoustic tests performed with the seven selected nozzle configurations was 196. Details of the tests are given in Appendix A. These tables list the inner, outer, and mixed stream conditions of the test points along with the PNL measured at $\theta_{\rm I}=50^\circ$, 70° , 90° , 110° , 130° , and 140° , and the OAPWL. PNL data have been scaled and extrapolated to a 1400 in. 2 nozzle exhaust area and a 2400 ft sideline. Detailed acoustic test results, including spectral data for each of the test points, are presented in the Comprehensive Data Report.

4.2.2 Laser Velocimeter Tests

The aerodynamic flow conditions of the LV tests are presented in Appendix B.

5.0 TEST RESULTS AND DISCUSSION

The analyses of the acoustic and laser velocimeter measurements obtained with the annular plug nozzle configurations of this program are discussed in this section. Descriptions of the nozzle configurations and the range of test conditions were presented in Section 4.0.

This section consists of three main subsections. Subsection 5.1 contains a discussion of the static and simulated flight acoustic data with emphasis on the influences of various coannular plug nozzle geometries and aerodynamic flow conditions. The measured shock noise data of coannular plug nozzles also are discussed in this subsection. Subsection 5.2 contains a discussion of the mean and turbulent velocity measurements taken with the LV. In conclusion, Subsection 5.3 describes a unique coannular spectral prediction method developed as a part of this contract effort.

5.1 ACOUSTIC TEST RESULTS

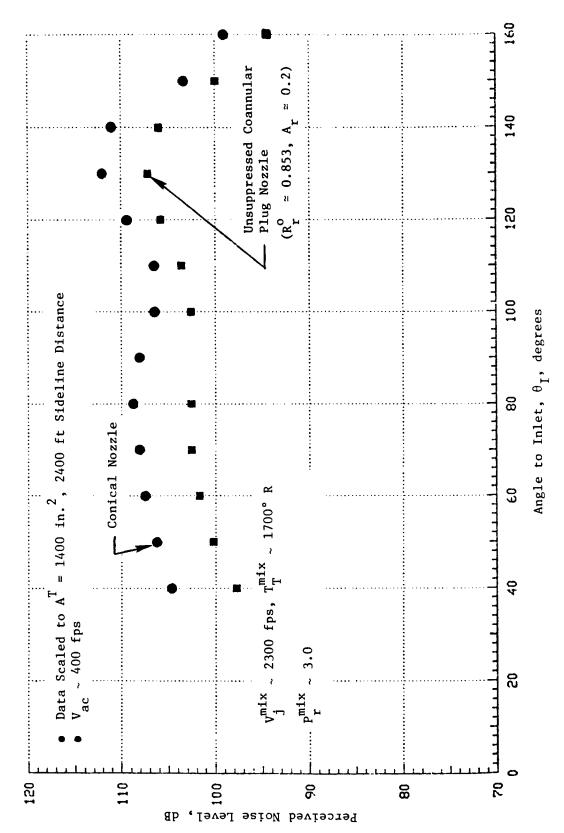
5.1.1 Verification of High-Radius-Ratio Coannular Plug Nozzle Jet Noise Reduction in Simulated Flight

Earlier experimental investigations reported in References 2 and 17 showed that significant jet and shock noise reductions were obtained for high-radius-ratio coannular plug nozzles relative to a conical nozzle at the same specific thrust in a static environment. A key objective of this investigation was to verify this important noise reduction feature under a simulated flight environment.

In the course of the discussion of results presented in this section, a number of flight influences will be illustrated for different geometric or thermodynamic flow variations. In addition, general results that verify in simulated flight the high-radius-ratio coannular plug nozzle jet and shock noise reductions previously measured statically are presented.

The first illustration of the verification of high-radius-ratio coannular plug nozzle total jet noise reduction in flight is a plot of the measured PNL directivity for a conical nozzle and a coannular plug nozzle of radius ratio 0.853, an inner-to-outer area ratio of approximately 0.2, and at the same specific thrust for a simulated flight speed of approximately 390 fps. The coannular plug nozzle chosen for the comparisons is the nozzle with struts (Model 1A) which is representative of the similitude YJ101 nozzle configuration. This model nozzle has been designed based on one-dimensional Mach number simulation of the exhaust nozzle flowpath of the baseline coannular YJ101 nozzle. Details of the engine nozzle and summary of the measured data are in Reference 9. The comparison of this model σ

conical nozzle is presented in Figure 9. The results are scaled to a typical supersonic cruise engine size of 1400 in. 2 at a 2400 foot sideline distance. The sample case is for a typical AST/VCE



Verification of High Radius Ratio Coannular Plug Nozzle PNL Directivity: Jet and Shock Noise Reduction at Typical Takeoff Sideline Engine Cycle and Flight Conditions. Figure 9.

takeoff sideline engine cycle operating condition (V_j^{mix} ~ 2300 fps and p_r^{mix} ~ 3.0). This result shows that, in simulated flight, the high-radius-ratio coannular plug nozzle has maintained jet noise and shock noise reduction relative to a conical convergent circular nozzle at equivalent specific thrust (V_j^{mix}) and nozzle pressure ratio (p_r^{mix}). The peak angle jet noise reduction measured for this case is 5 PNdB at θ_I = 130°, and 6 PNdB shock noise or forward radiated noise at θ_I = 60°. The OASPL directivity for this case is shown in Figure 10. For this measurement, the peak jet noise angle is 140°, but the relative jet noise reduction at the peak angle is also 5 dB. In the forward quadrant at θ_I = 60°, the shock noise or forward quadrant noise reduction on a OASPL basis is as high as 8 dB.

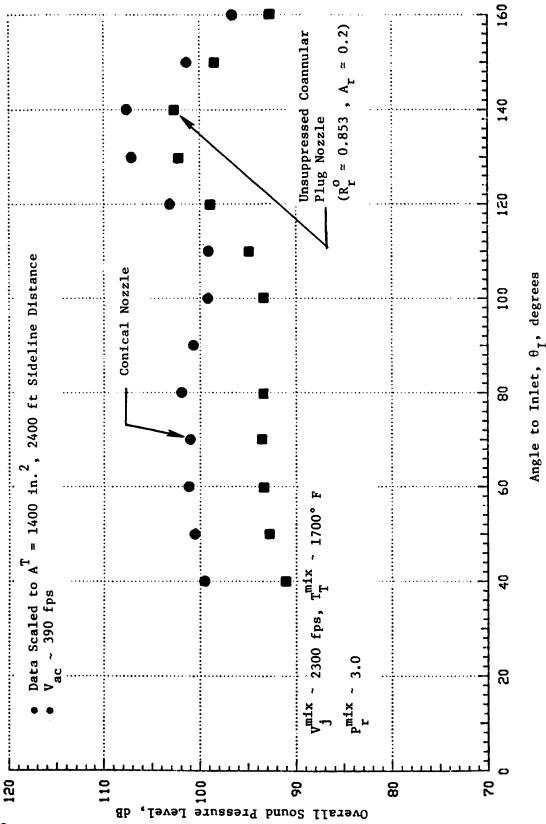
Comparisons of the coannular plug nozzle acoustic spectral distribution with those of conical nozzle for observation locations from angle to the inlet, $\theta_{\rm I}$, of 40° to 160° are shown in Figure 11. These test measurements show that at a simulated flight condition of ~390 fps, the forward quadrant shock noise is considerably reduced – up to 12 dB on a peak spectral basis. In the aft quadrant, where jet mixing noise is the dominant exhaust noise mechanism, the high-radius-ratio coannular plug nozzle noise reduction is observed to occur nearly over all the measured frequency bands. Figures 12 through 17 illustrate the general static and simulated flight results measured for the typical takeoff sideline AST/VCE engine cycle condition.

Figures 12 and 13 show the static-to-simulated flight PNL and OASPL directivity characteristics for the circular conical convergent exhaust nozzle. What can be observed from these comparisons is that in the aft microphone observation angles the jet mixing noise has been reduced inflight. In the forward quadrant, or shallow observation angles, a shock associated noise "lift" is observed. The measurements indicate a 4 PNdB noise reduction due to flight at the peak aft quadrant noise angle, but a corresponding 4 PNdB forward quadrant amplification of noise at $\theta_{\rm T} = 60^{\circ}$.

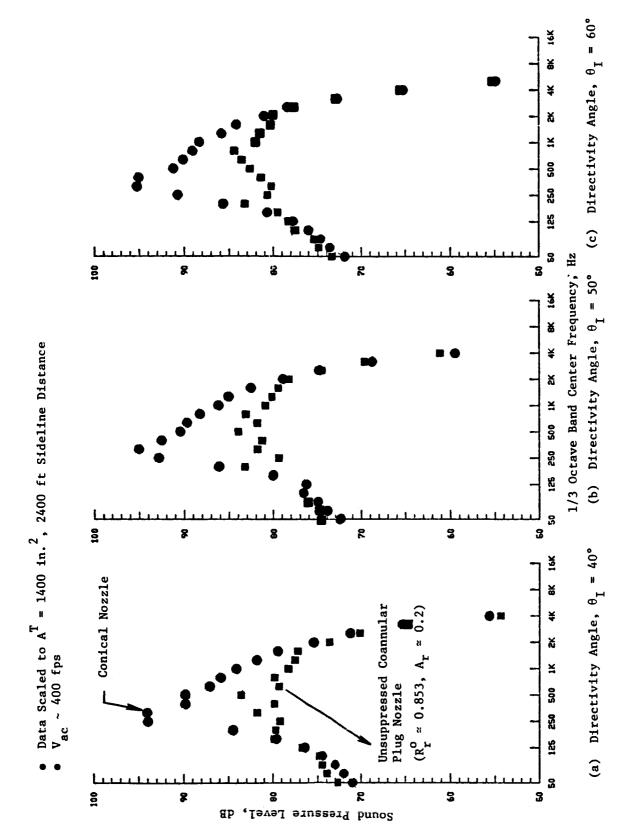
The spectral static-to-flight comparisons for the conical nozzle are shown in Figure 14. These results show that in the forward quadrant there is a Doppler shift of the shock associated noise toward higher frequencies associated with a general noise amplification. In the aft jet noise measurement angles, there is a general reduction in the noise signature over all frequency bands.

A comparison set of static-to-simulated flight acoustic measurements for the high-radius-ratio coannular plug nozzle is shown in Figures 15 through 17. As we generally observed with the conic nozzle, the high-radius-ratio coannular plug nozzle also shows for this case that in the aft quadrant at the peak noise angle there was a reduction in the jet noise with flight. In the forward quadrant (e.g., at $\theta_{\rm I}=60^{\circ}$), an amplification of the shock associated noise is observed again. At the peak jet noise angle, the coannular plug nozzle jet noise was reduced by 3.5 PNdB due to flight, while the forward quadrant lift was observed to be about 4.5 dB at $\theta_{\rm T}=60^{\circ}$ (Figure 15).

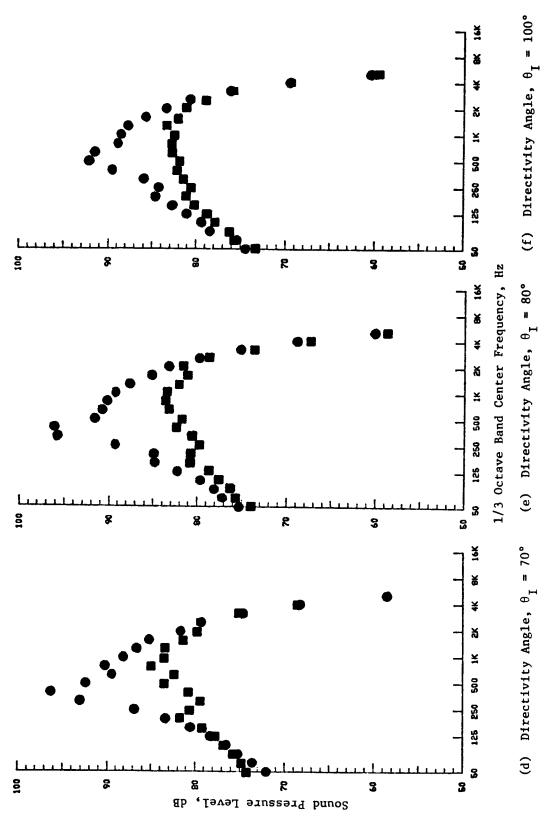
The spectral static-to-simulated-flight acoustic test results for the unsuppressed coannular plug nozzle are shown in Figure 17. In the forward obser-



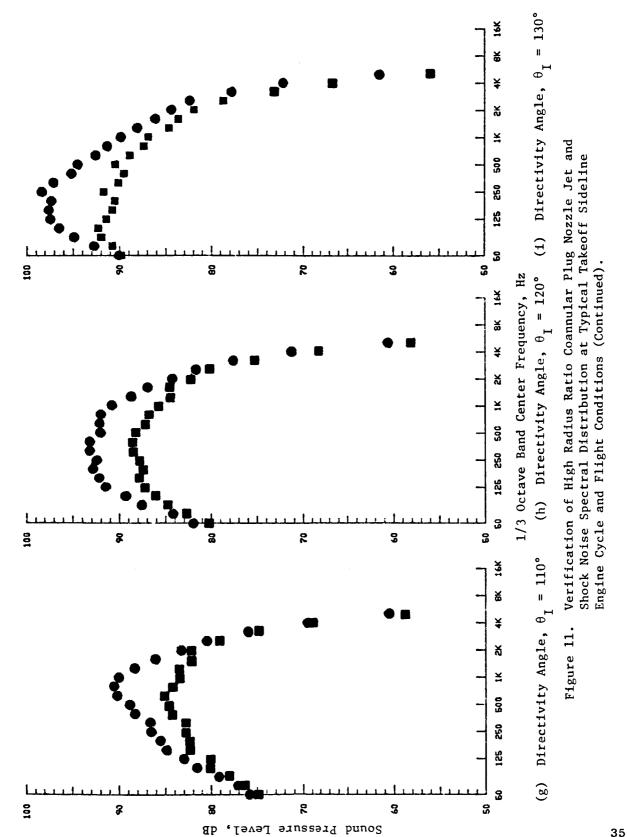
Directivity: Jet and Shock Noise Reduction at Typical Takeoff Verification of High Radius Ratio Coannular Plug Nozzle OASPL Sideline Engine Cycle and Flight Conditions. Figure 10.

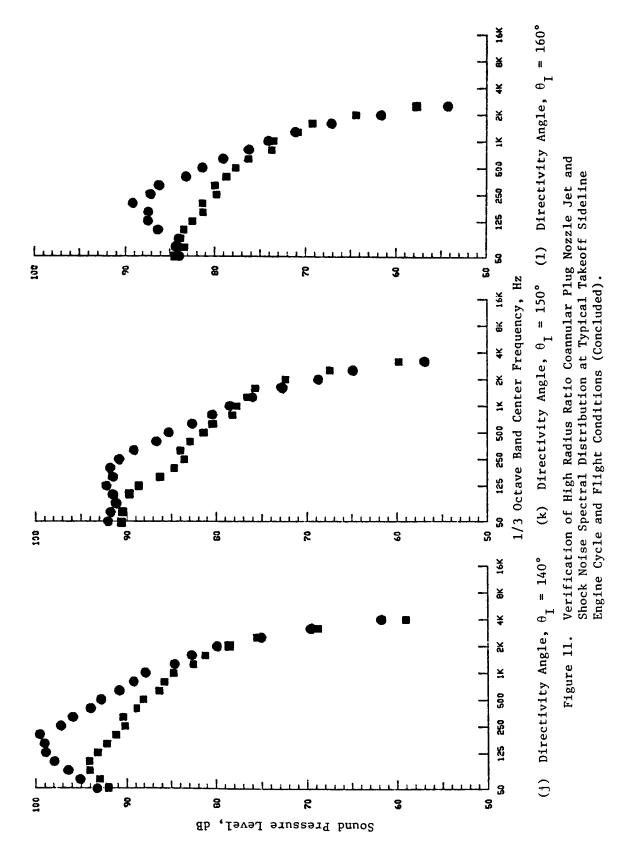


Verification of High Radius Ratio Coannular Plug Nozzle Jet and Shock Noise Spectral Distribution at Typical Takeoff Sideline Engine Cycle and Flight Conditions. Figure 11.



Verification of High Radius Ratio Coannular Plug Nozzle Jet and Shock Noise Spectral Distribution at Typical Takeoff Sideline Engine Cycle and Flight Conditions (Continued): Figure 11.





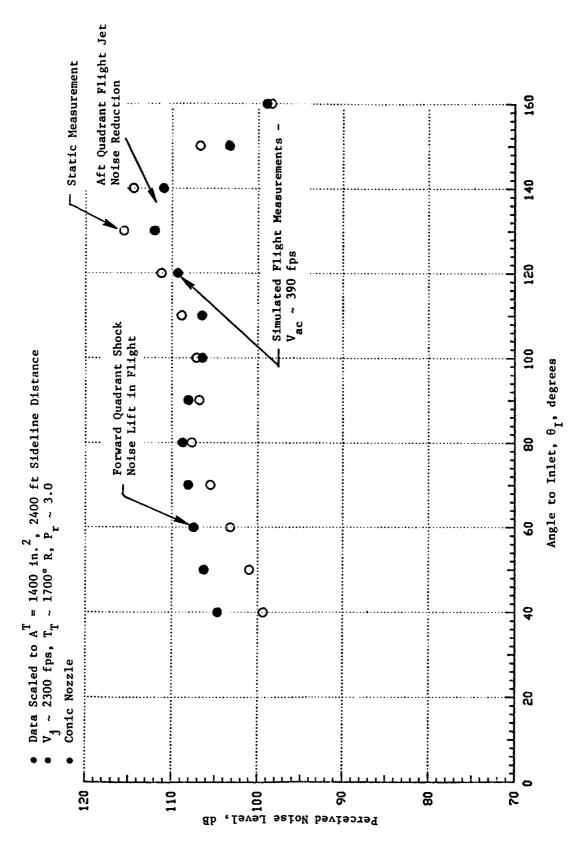
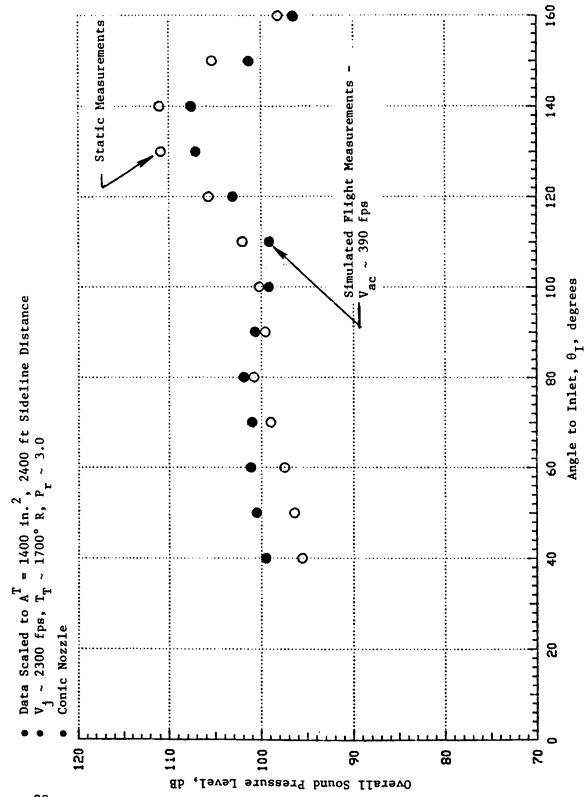
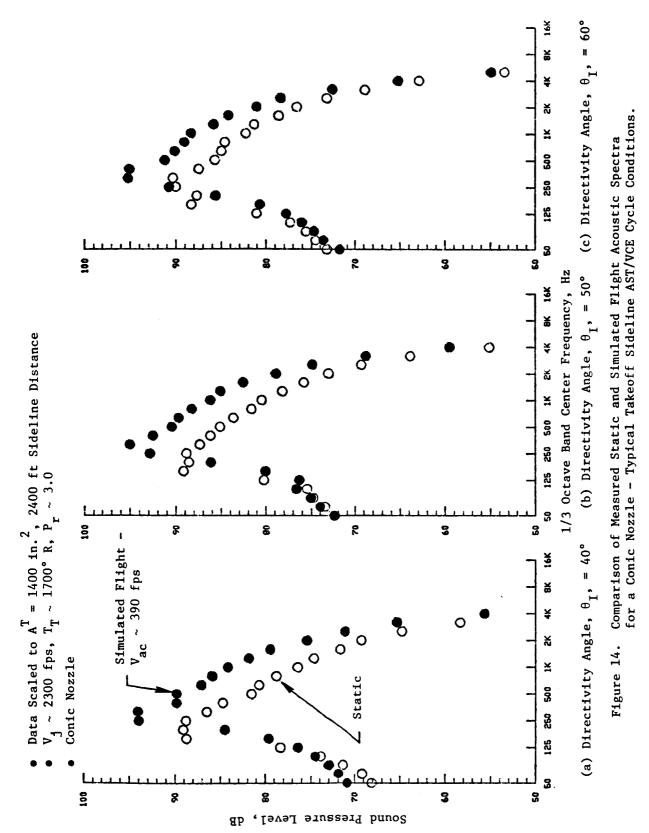
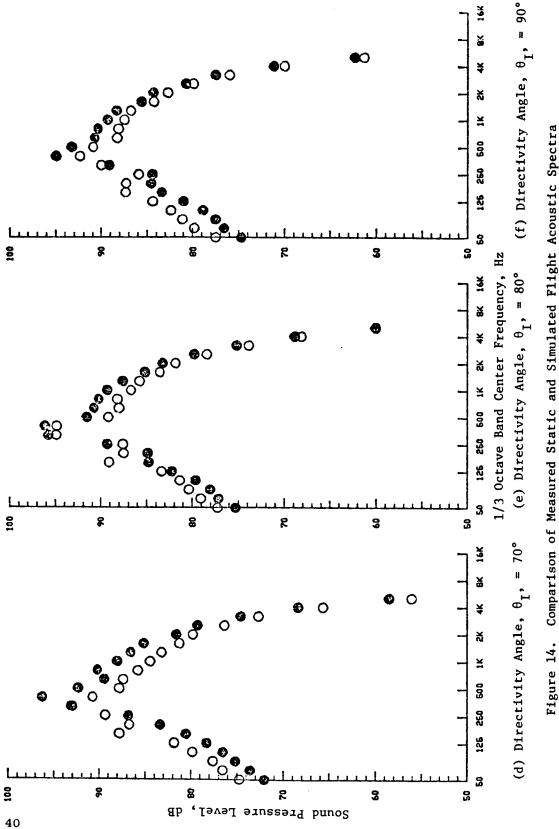


Figure 12. Comparison of Measured Static and Simulated Flight PNL Directivity of a Conic Nozzle at a Typical Takeoff Sideline AST/VCE Cycle Condition.

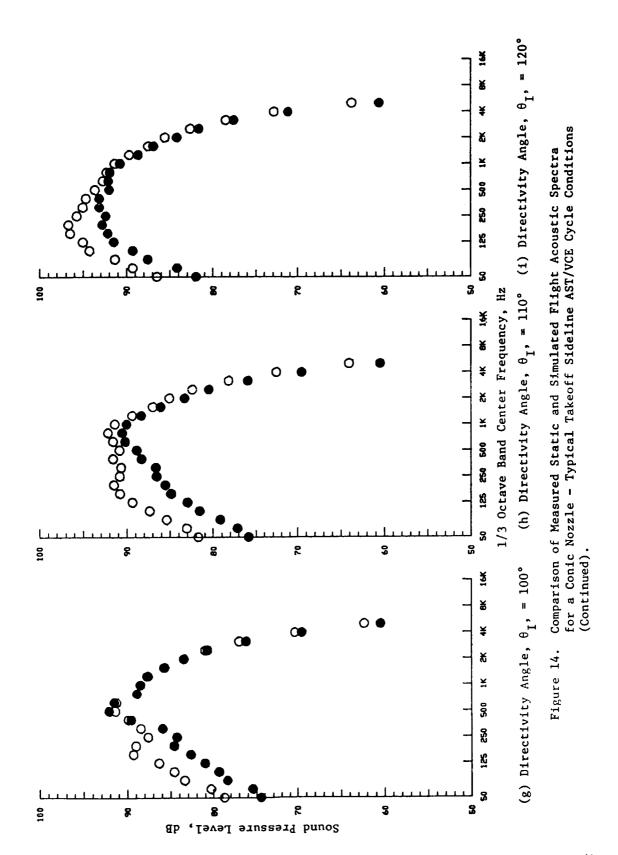


Directivity of a Conic Nozzle at a Typical Takeoff Sideline Comparison of Measured Static and Simulated Flight OASPL AST/VCE Cycle Condition. Figure 13.





for a Conic Nozzle - Typical Takeoff Sideline AST/VCE Cycle Conditions (Continued).



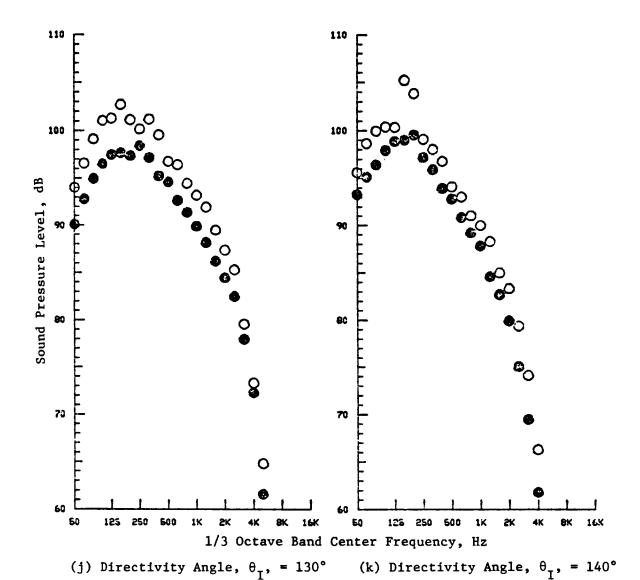
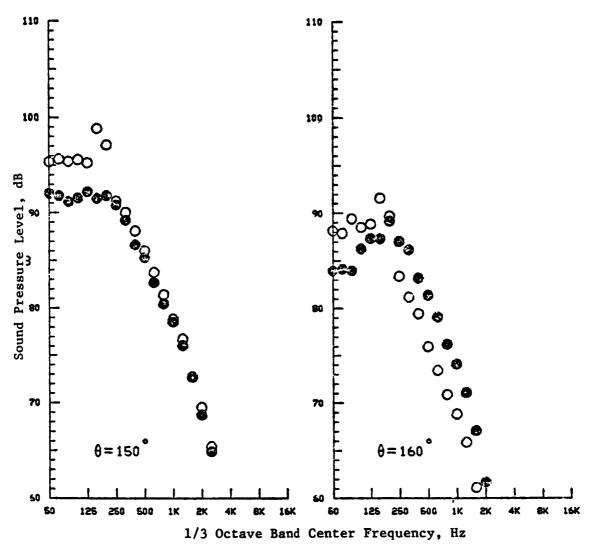
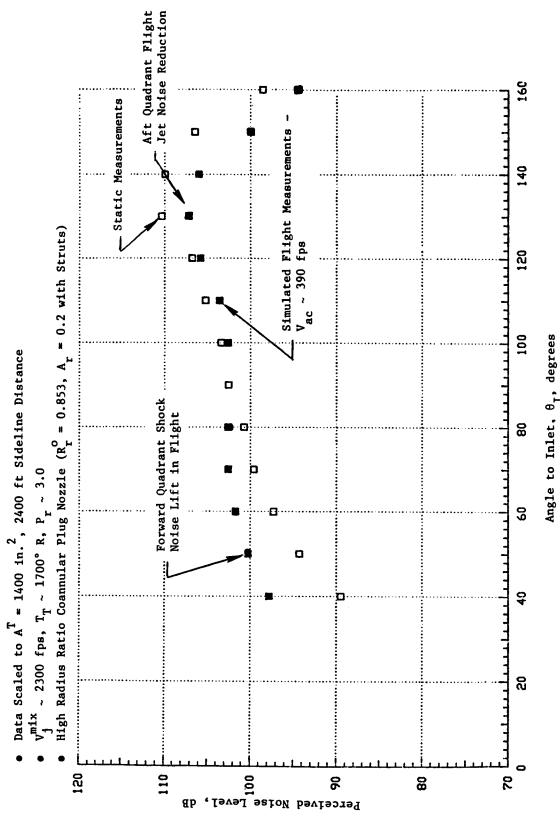


Figure 14. Comparison of Measured Static and Simulated Flight Acoustic Spectra for a Conic Nozzle - Typical Takeoff Sideline AST/VSE Cycle Conditions (Continued).

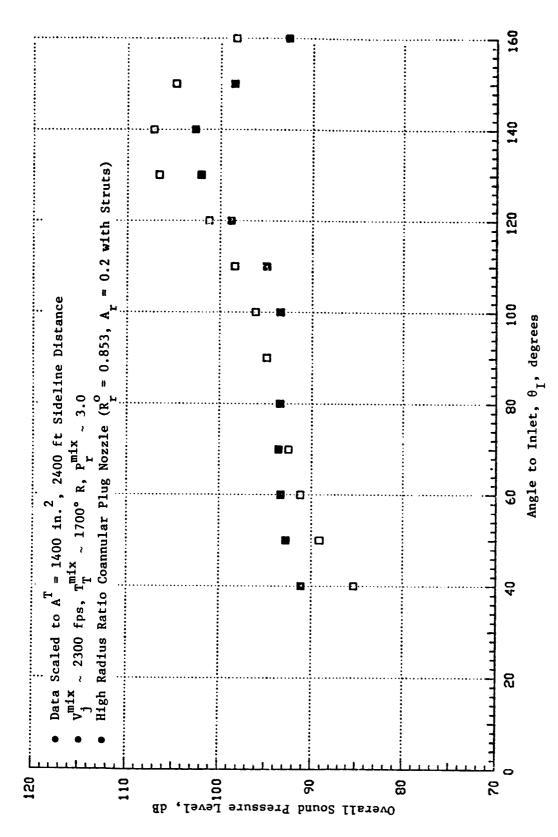


(1) Directivity Angle, $\theta_{\rm I}$, = 150° (m) Directivity Angle, $\theta_{\rm I}$, = 160°

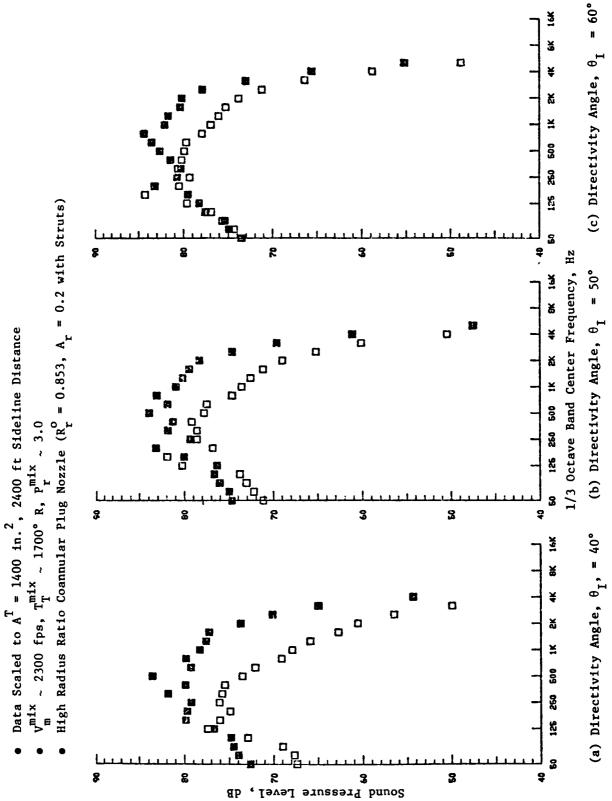
Figure 14. Comparison of Measured Static and Simulated Flight Acoustic Spectra for a Conic Nozzle - Typical Takeoff Sideline AST/VCE Cycle Conditions (Concluded).



Comparison of Measured Static and Simulated Flight PNL Directivity of a Coannular Plug Nozzle at a Typical Takeoff Sideline AST/VCE Cycle Condition. Figure 15.



Directivity of a Coannular Plug Nozzle at a Typical Takeoff Comparison of Measured Static and Simulated Flight OASPL Sideline AST/VCE Cycle Condition. Figure 16.



Coannular Plug Nozzle at Typical Takeoff Sideline AST/VCE Cycle Conditions. Figure 17. Comparison of Measured Static and Simulated Flight Acoustic Spectra for a

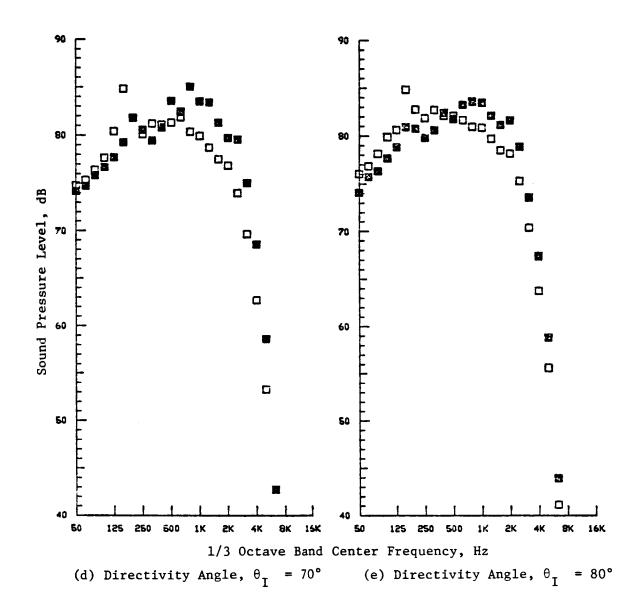
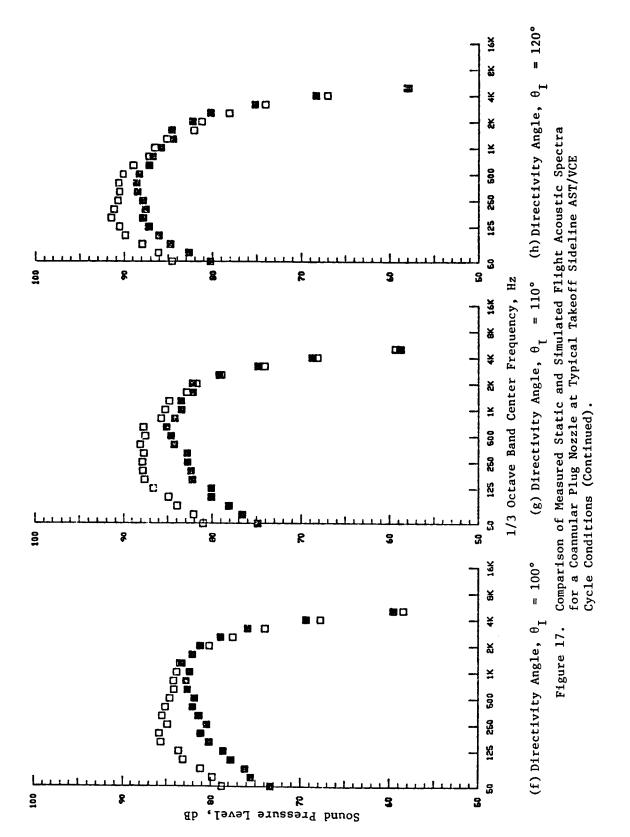


Figure 17. Comparison of Measured Static and Simulated Flight Acoustic Spectra for a Coannular Plug Nozzle at Typical Takeoff Sideline AST/VCE Cycle Conditions (Continued).



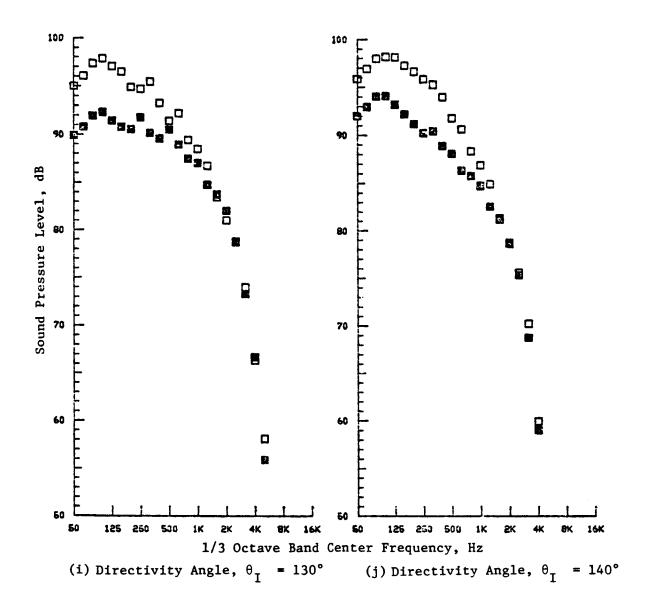


Figure 17. Comparison of Measured Static and Simulated Flight Acoustic Spectra for a Coannular Plug Nozzle at Typical Takeoff Sideline AST/VCE Cycle Conditions (Continued).

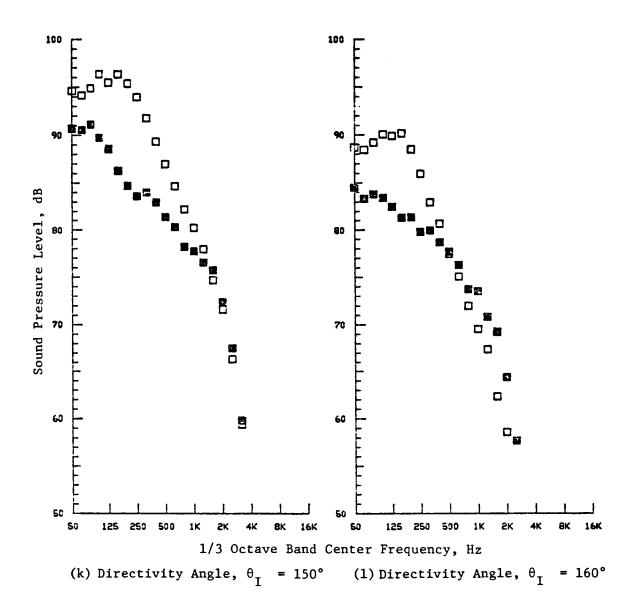


Figure 17. Comparison of Measured Static and Simulated Flight Acoustic Spectra for a Coannular Plug Nozzle at Typical Takeoff Sideline AST/VCE Cycle Conditions (Concluded).

vation angles, the amplification of the coannular shock associated noise over the entire spectral range is observed. In the aft quadrant, the flight noise reduction is observed to be primarily associated with the lower frequency bands of the scaled data (50 Hz to 2000 Hz); while at the higher frequencies, the static and flight spectrum are about the same or actually slightly higher for the flight conditions.

As a final general observation from this data set, Figure 18 shows a relative PNL and OASPL differences between the static and flight measurements in the forward quadrant for the conic and high-radius-ratio coannular plug nozzle in order to evaluate the forward quadrant shock noise lift. Shown on the figure is a predicted dynamic effect [-40 log (1 - $\rm M_{ac}$ Cos $\theta_{\rm I})]$ for shock noise. In general, the results seem to be in fair agreement with the simple dynamic effect prediction for the conic nozzle and the coannular plug nozzle measurements.

5.1.2 Influence of Geometry on Coannular Plug Nozzle Acoustics

5.1.2.1 Influence of Internal Struts

In order to study the influence of the eight internal struts in the outer stream of the VCE/AST test-bed nozzle on the measured far-field noise data, tests were conducted with Models 1A and 2 that are geometrically scaled versions of the test-bed nozzle but designed with and without struts, respectively. The measured static and simulated flight ($V_{ac} = 390 \, \text{ft/sec}$) normalized PNL_{max} data that are scaled to a 2400-ft sideline, 59° F standard day and a 1400 in. 2 nozzle exhaust area are presented in Figure 19 as a function of $10 \, \log(V_j^{\text{mix}}/c_a)$. An examination of this figure indicates that the PNL_{max} data with and without struts generally agree. Further, the normalized PNL_{max} data obtained from the YJ101 engine static tests with a nozzle similar to Model 1A are presented in Figure 19. Good agreement between the engine and model data is noted.

The PNL directivity patterns of Models 1A and 2 at a typical takeoff condition ($V_{:}^{mix}$ = 2250 ft/sec, P_{r}^{o} = 3.2), measured under static and simulated flight conditions, are presented in Figure 20. Even though the data with/without struts reasonably agree in the aft quadrant, Figure 20 indicates that the nozzle with struts is less noisy in the forward quadrants. Such is the case from the data presented in Figure 21 comparing the frequency spectra at θ_{I} = 60° and 130° for Models 1A and 2 and at the typical takeoff condition. Furthermore, Figure 21 indicates that the forward angle SPL of Model 2 (without struts) is usually greater than that of Model 1A (With Struts) for all 1/3-octave band frequencies that are greater than 200 Hz.

Observations similar to those made earlier at the typical takeoff condition were noted at the underexpanded ($v_j^{\text{mix}} \sim 2380 \text{ ft/sec}$, $P_r^0 = 3.8$) test case. However, no significant differences in the with/without strut data were observed for the overexpanded cutback case ($v_j^{\text{mix}} \sim 1920 \text{ ft/sec}$, $v_r^0 = 2.3$).

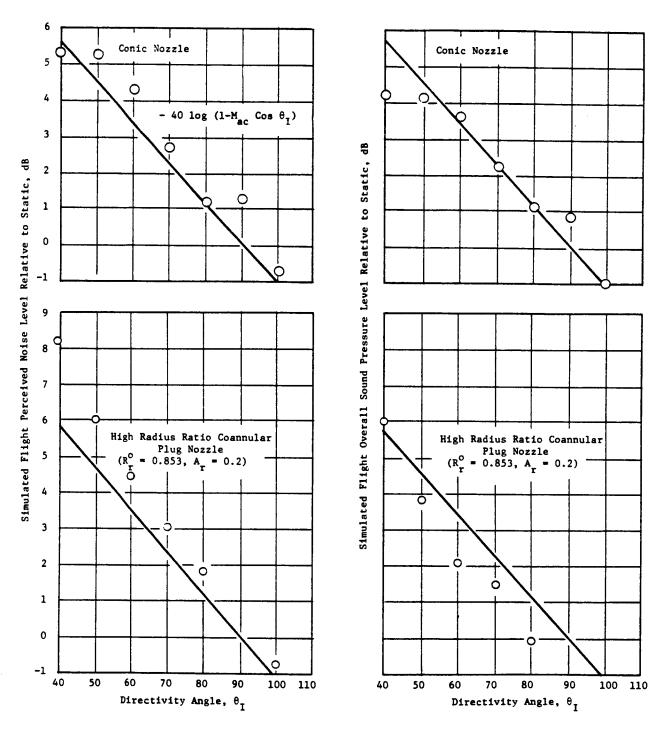
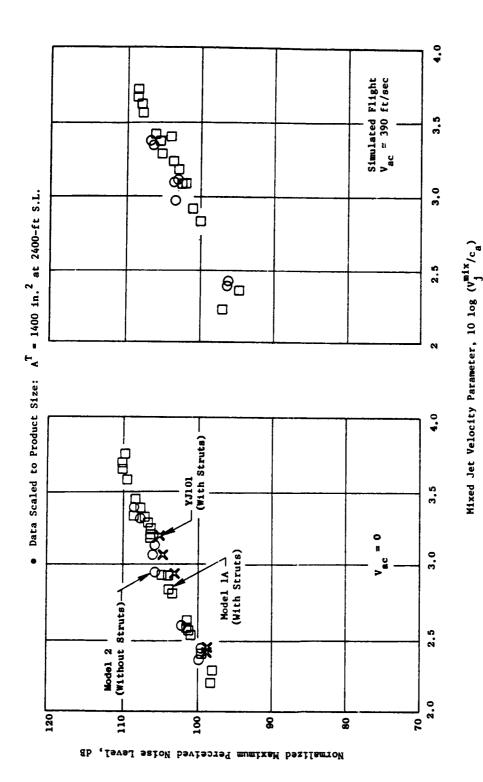


Figure 18. Simulated Flight PNL and OASPL Relative to Static Levels in Forward Quadrant.



Normalized Maximum Perceived Noise Level Results for Model 1A (With Struts) and Model 2 (Without Struts). Figure 19.

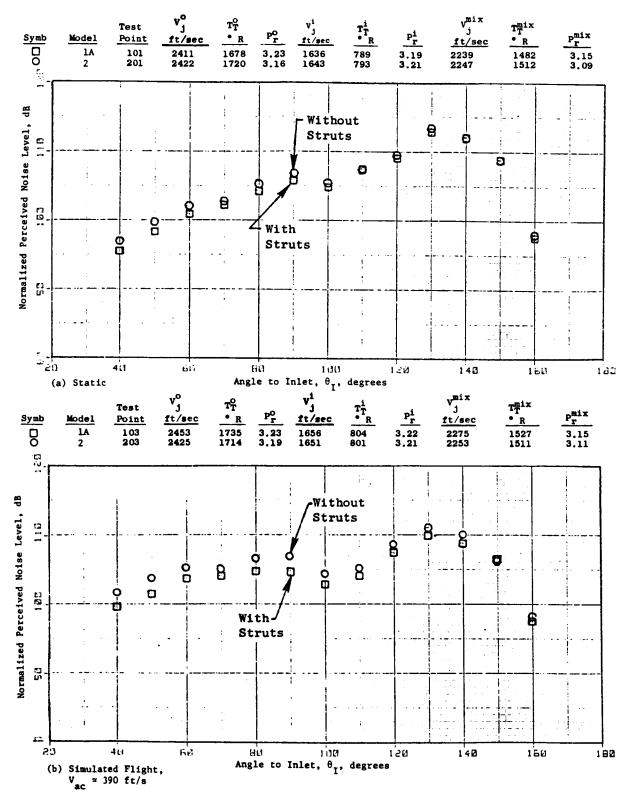


Figure 20. Comparison of Normalized PNL Directivity of Nozzles With and Without Struts; Typical Takeoff, $V_j^{mix} \simeq 2250$ Ft/S, $P_r^o = 3.2$.

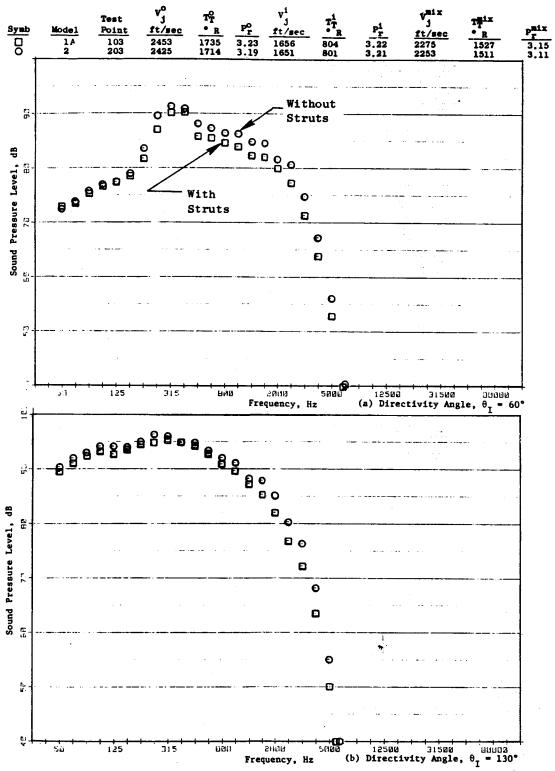


Figure 21. Comparison of Frequency Spectra of Model 1A (With Struts) and Model 2 (Without Struts), Simulated Flight V = 390 Ft/S at Typical Takeoff.

5.1.2.2 Influence of Outer Stream Termination

To determine the effect of forward flight on the acoustic effectiveness of a convergent-divergent (C-D) termination in the outer stream of the VCE/AST test-bed scale model, tests were conducted with a Model 2 nozzle having a C-D outer termination that is designed for a typical takeoff condition (P_r° = The measured static and flight acoustic data are presented in Figures 22 and 23 and are compared with the corresponding data obtained with Model 3 which has its outer stream terminated at the throat in order to yield a convergent termination. An examination of the figures indicates no significant differences, under both static and simulated flight conditions, in the forward quadrant acoustic data of C-D and convergent terminated nozzle configurations. But, in the aft quadrant, the convergent terminated configuration is observed to be beneficial, particularly during simulated flight. This aft quadrant benefit with the convergent terminated nozzle and with no significant differences in the forward quadrant data of the two nozzles is observed to be not only at the C-D optimum design pressure ratio but also at the other overexpanded/underexpanded test cases. This is made clear from the data presented in Figure 24 that compares the PNL $_{60}$ and the normalized PNL $_{
m max}$ data of Model 2 with that of Model 3 over a range of $V_i^{mix} = 2380 \div 1900$ ft/sec and $P_r^0 = 3.8 + 2.3.$

5.1.2.3 Influences of Outer Stream Radius Ratio and Nozzle Area Ratio

The objectives of this study are to determine, under static and simulated flight conditions, (1) the effect of the outer stream radius ratio $R_{\rm r}^{\rm O}$ for a given nozzle area ratio $A_{\rm r}$ (defined = $A^{\rm i}/A^{\rm O}$) and (2) the effect of nozzle area ratio for a constant outer stream radius ratio on the acoustic characteristics of coannular plug nozzles. The measured data are presented in this section.

(1) Radius Ratio Effect

The configurations employed for this study are:

Mode 1	A _r	$\frac{R_r^0}{}$	Outer Termination				
4	0.53	0.853	Convergent				
6	0.53	0.902	Convergent				

The tests included a series where the inner-to-outer stream velocity ratio was varied. This was achieved by holding the outer stream velocity constant at $V_j^0 \sim 2300$ ft/sec and regulating the inner stream velocity V_j^i so as to obtain velocity ratios of 0.16 to 0.70.

The measured normalized PNL_{max} data are summarized in Figure 25. The data indicate that, under both static and simulated flight conditions and for a given area ratio, an increase in the radius ratio of the outer annular nozzle

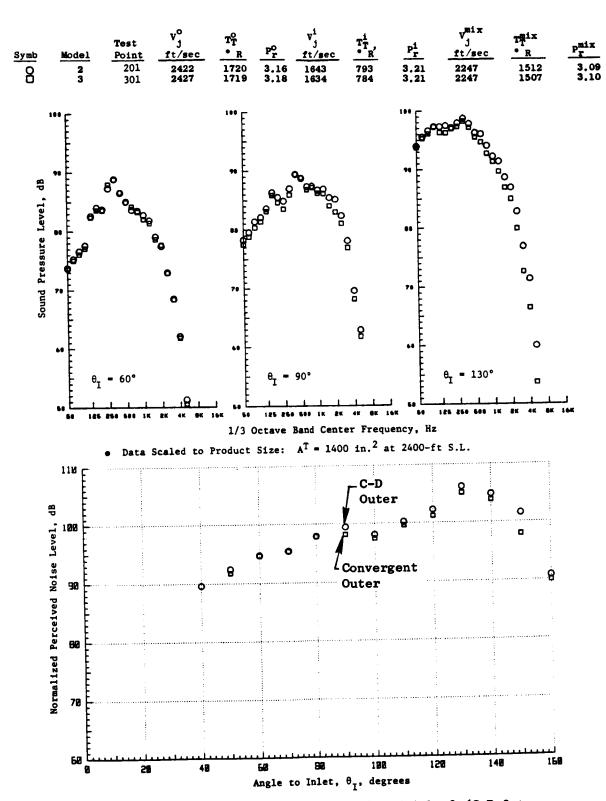


Figure 22. Comparison of Acoustic Data for Models 2 (C-D Outer Termination) and 3 (Convergent Outer Termination) at Typical Takeoff Condition at Design $P_r^0 = 3.21$ (Static).

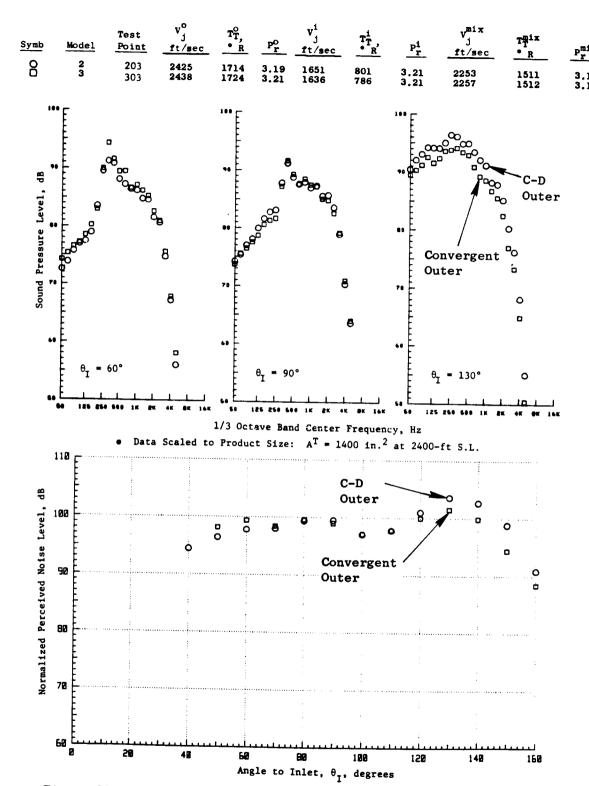
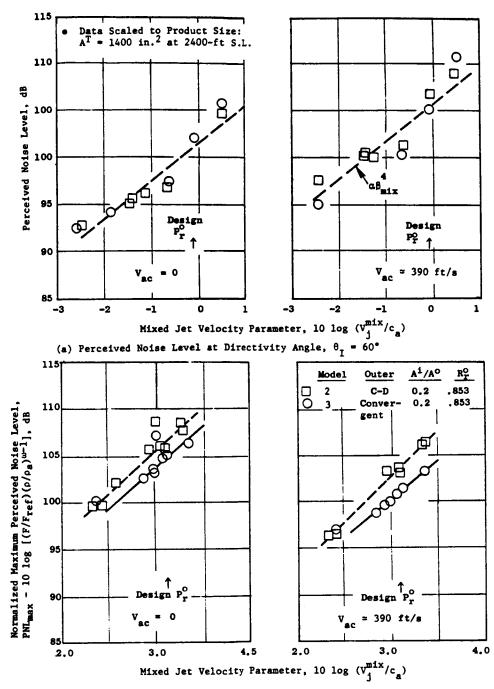
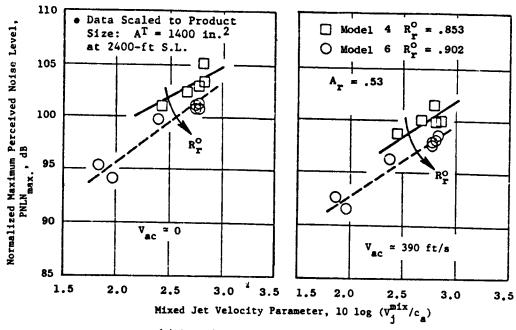


Figure 23. Comparison of Acoustic Data for Models 2 (C-D Outer Termination) and 3 (Convergent Outer Termination) at Typical Takeoff Condition at Design $P_r^0 = 3.21$, Simulated Flight $V_{ac} \simeq 390$ Ft/S.

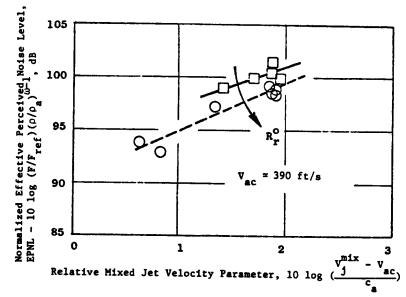


(b) Normalized Maximum Perceived Noise Level

Figure 24. Perceived Noise Levels as a Function of Mixed Jet Velocity Parameter for Models 2 (C-D Terminated) and 3 (Convergent Terminated).







(b) Normalized Maximum Perceived Noise Level

Figure 25. Influence of Outer Stream Radius Ratio on Normalized PNL and EPNL; Inner-to-Outer Area Ratio = 0.53.

is significantly beneficial in the test velocity range of $V_{i}^{mix} = 1800$ to 2200 ft/sec. This is indicated also on a normalized EPNL basis in Figure 25.

The data obtained from the tests involving variation in the inner-to-outer-stream velocity ratio are presented in Figure 26. The figure indicates that the PNL_{max} benefit due to a high outer stream radius ratio is maintained over the test velocity ratio range and under both static and simulated flight conditions.

Typical frequency spectra (at $\theta_{\rm T}=60^{\circ}$, 90°, and 130°/140°) and normalized PNL directivities that compare Models 4 and 6 static and flight data at vmix = 2100 ($V_{\rm T} \approx 0.2$, 0.7) and 1900 ft/sec ($V_{\rm T} \approx 0.4$) are presented in Figures 27 through 29. An examination of the PNL directivities indicates that, while under static conditions, only aft angles benefit from a higher outer radius ratio; beneficial acoustic characteristics are observed at all angles in simulated flight. In general, the higher radius ratio nozzle results in a lower SPL at all frequencies.

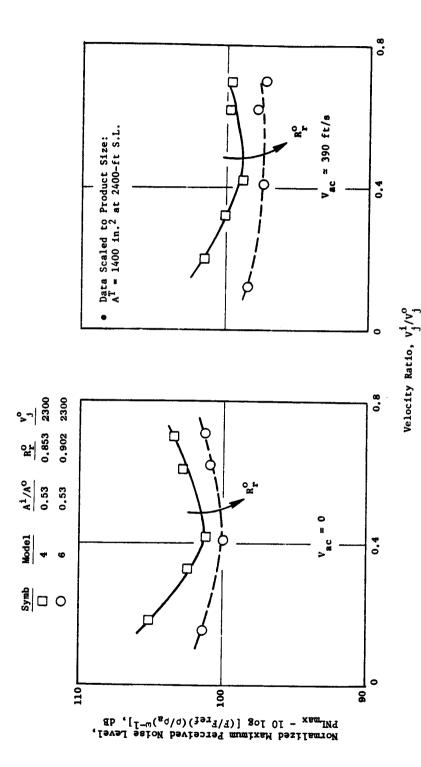
(2) Area Ratio Effect

The configurations employed for this study are:

Model	R°	A _r	Outer Termination
3	0.853	0.20	Convergent
7	0.853	0.33	Convergent
4	0.853	0.53	Convergent

The measured normalized PNL_{max} data of this study are presented in Figure 30 along with YJ101 engine static data obtained with a convergent terminated coannular plug nozzle having $A_r = 0.2$ and $R^0 = 0.853$. These data indicate that, for a fixed outer stream radius ratio and V_j^{mix} , PNL_{max} decreases with a decrease in the nozzle area ratio. Also, over the test velocity range of $V_j^{mix} = 1800 \div 2400$ ft/sec and under static and simulated flight situations, the configuration with $A_r = 0.2$ yielded the lowest measured aft quadrant, PNL_{max} data.

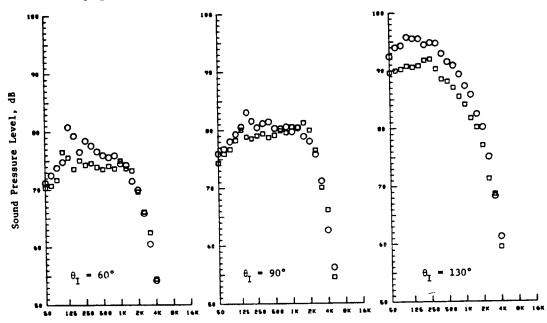
Typical frequency spectra ($\theta_I = 60^\circ$, 90° , and $130^\circ/140^\circ$) and normalized PNL directivities of Models 3, 4, and 7 are presented in Figures 31 and 32. While in Figure 31 the data are plotted from tests having identical outer and mixed flow conditions ($V_j^0 \sim 2300$ and $V_j^{mix} \sim 2100$ ft/sec), the data presented in Figure 32 were obtained from tests having constant outer and inner conditions (and, hence, a constant V_r). An examination of the data indicates that a change in the area ratio had no significant effect on the front quadrant acoustics for all of the test cases. In the aft quadrant, although the nozzle with the smaller area ratio (0.2) yielded beneficial results for given outer and mixed stream conditions (due to a favorable V_r), the nozzle with the larger area ratio (0.53) yielded beneficial results for given inner and outer conditions (due to the smaller V_j^{mix}).



Influence of Outer Stream Radius Ratio on Normalized Maximum Perceived Noise Level at Different Inner-to-Outer Stream Velocity Ratios; Inner-to-Outer Stream Area Ratio = 0.53. Figure 26.

Symb	M od e l	Test Point	v ^o j ft/sec	$T_{\mathbf{T}}^{\mathbf{O}}$ • R	P°	v ⁱ j ft/sec_	$\mathbf{T}_{\mathbf{R}}^{\mathbf{i}}$	P _r	v ^{mix} j ft/sec	TT R	Pmix
0	4	481	2313	1736	2.78	425 371	980 637	1.06	2104 2039	1552 1584	2.39 2.35

Data Scaled to Product Size: A^T = 1400 in.² at 2400-ft S.L.



1/3 Octave Band Center Frequency, Ha

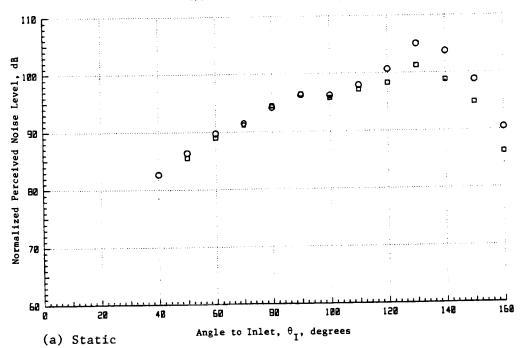
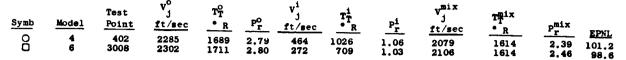
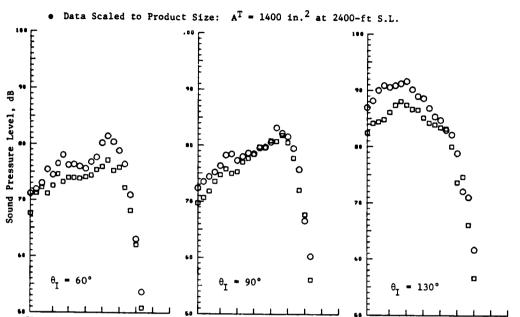


Figure 27. Comparison of Model 4 ($R_r^0 = 0.853$) and Model 6 ($R_r^0 = 0.902$) Frequency Spectrum and PNL Directivity at $V_r^{mix} \simeq 2100$ ft/s, $V_r^0 \simeq 0.2$, $A_r^0 = 0.53$.





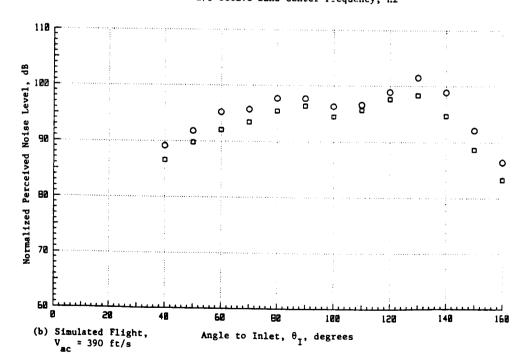
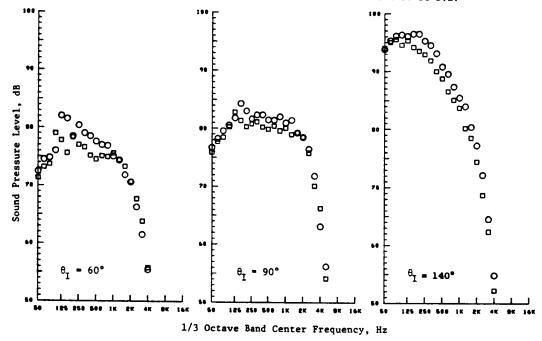


Figure 27. Comparison of Model 4 ($R_r^0 = 0.853$) and Model 6 ($R_r^0 = 0.902$) Frequency Spectrum and PNL Directivity at $V_j^{\text{mix}} \simeq 2100$ Ft/Sec, $V_r \simeq 0.2$, $A_r = 0.53$ (Concluded).

Data Scaled to Product Size: $A^{T} = 1400 \text{ in.}^{2}$ at 2400-ft S.L.



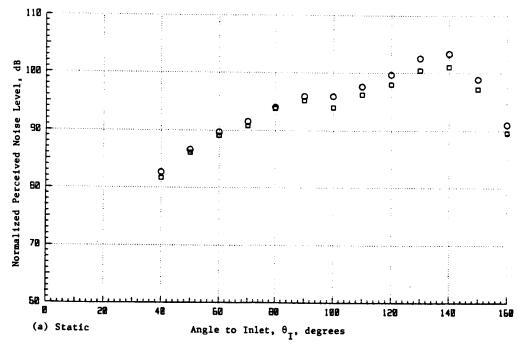


Figure 28. Comparison of Model 4 (R_{r}^{O} = 0.853) and Model 6 (R_{r}^{O} = 0.902) Frequency Spectrum and PNL Directivity at $V_{j}^{mix} \simeq 2100$ Ft/Sec, $V_{r} \simeq 0.7$, A_{r} = 0.53.

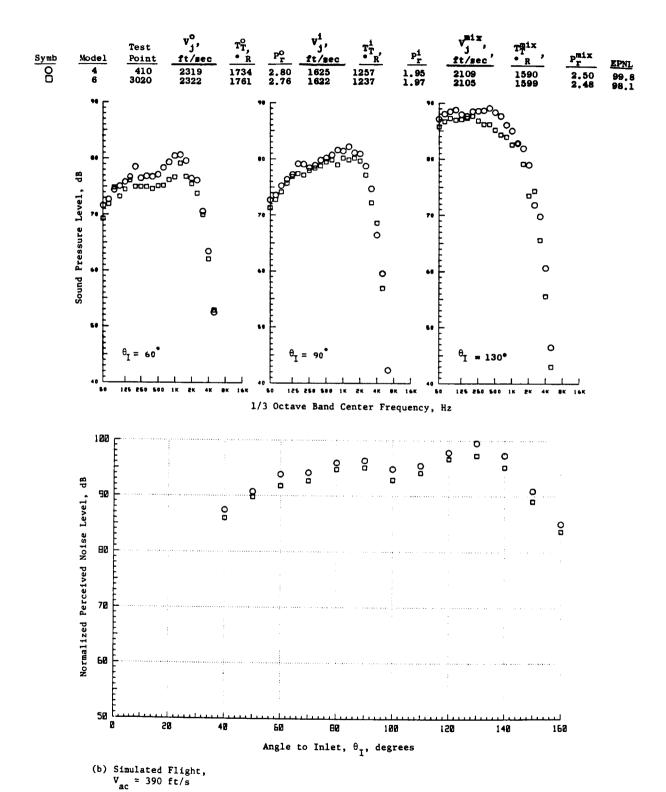


Figure 28. Comparison of Model 4 ($R_r^O = 0.853$) and Model 6 ($R_r^O = 0.902$)
Frequency Spectrum and PNL Directivity at $V_j^{mix} \simeq 2100$ Ft/Sec, $V_r \simeq 0.7$, $A_r = 0.53$ (Concluded).

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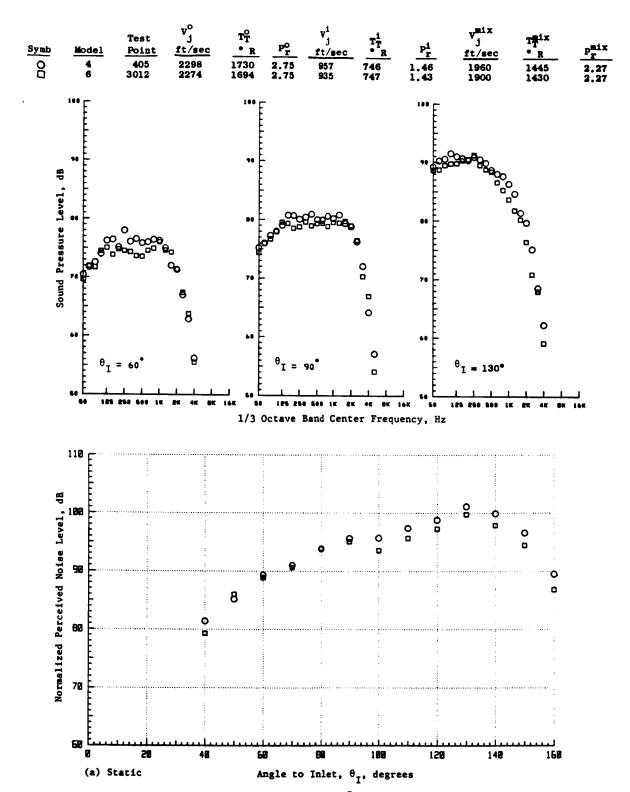


Figure 29. Comparison of Model 4 (R_r^0 = 0.853) and Model 6 (R_r^0 = 0.902) Frequency Spectrum and PNL Directivity at $V_r^{mix} \approx 1900$ Ft/Sec, $V_r \approx 0.4$, $A_r = 0.53$.

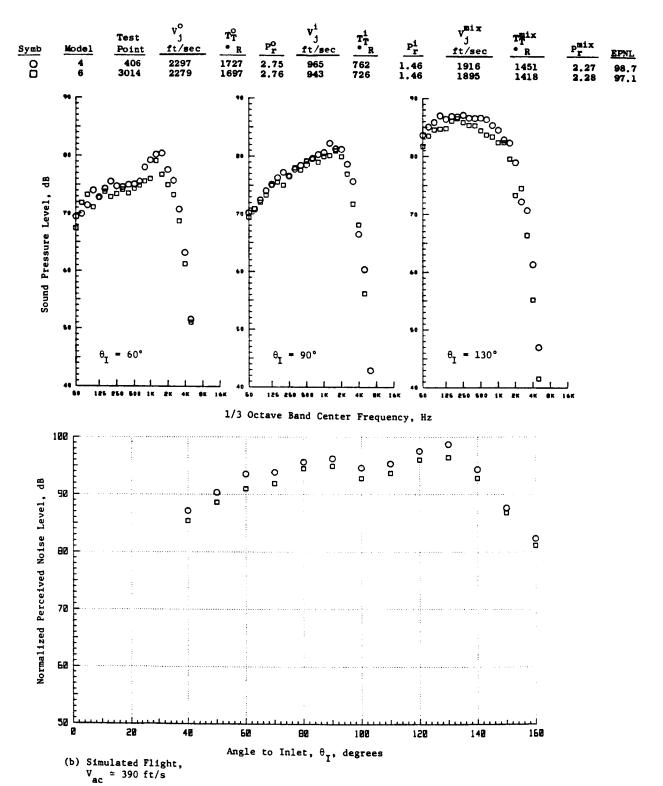


Figure 29. Comparison of Model 4 ($R_r^O = 0.853$) and Model 6 ($R_r^O = 0.902$)

Frequency Spectrum and PNL Directivity at $V_J^{mix} \simeq 1900$ Ft/

Sec, $V_r \simeq 0.4$, $A_r = 0.53$ (Concluded).

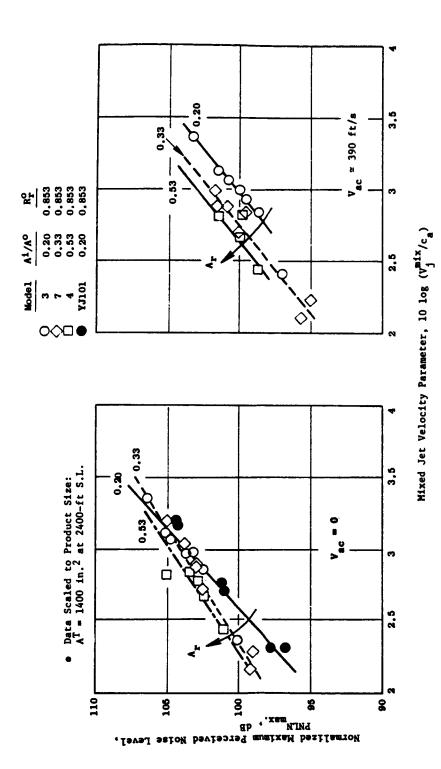


Figure 30. Effect of Area Ratio on PML for a Given Radius Ratio.

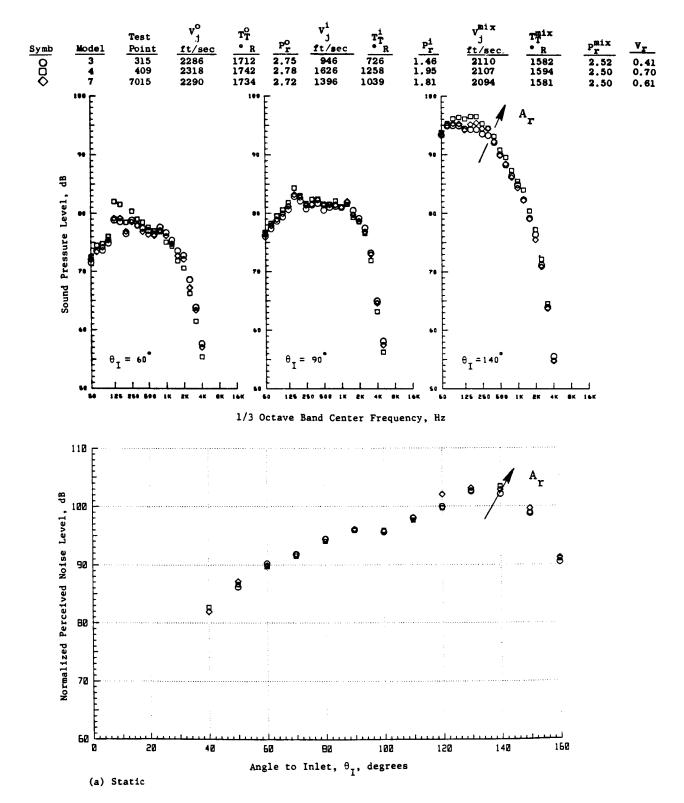


Figure 31. Effect of Area Ratio for a Given Radius Ratio (0.853) and Fixed Outer and Mixed Conditions.

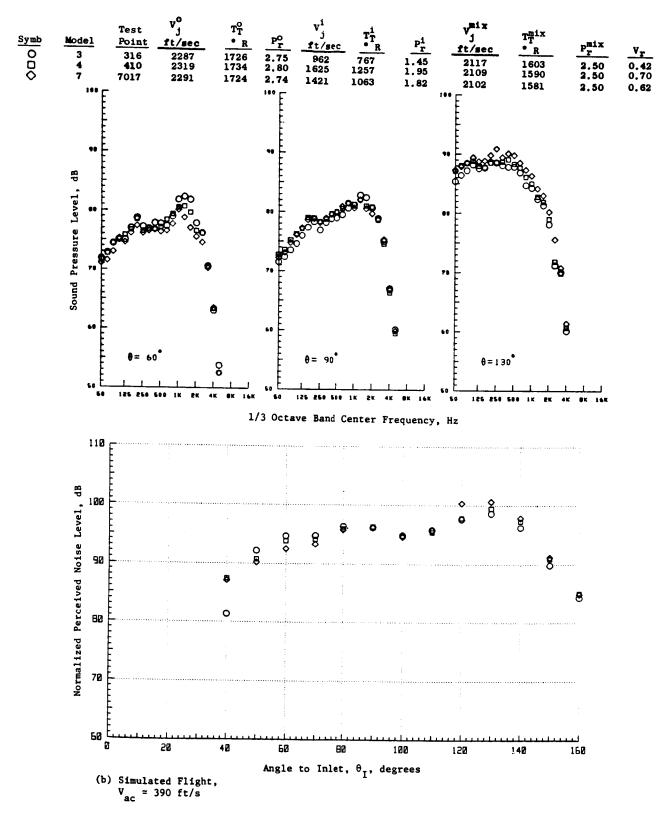


Figure 31. Effect of Area Ratio for a Given Radius Ratio (0.853) and Fixed Outer and Mixed Conditions (Concluded).

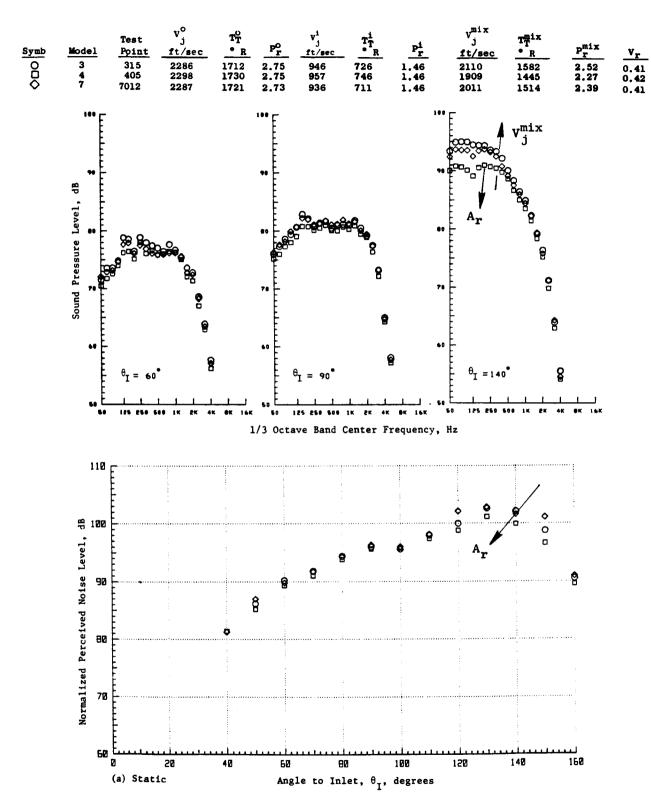


Figure 32. Effect of Area Ratio for a Given Radius Ratio (0.853) and Fixed Inner and Outer Conditions.

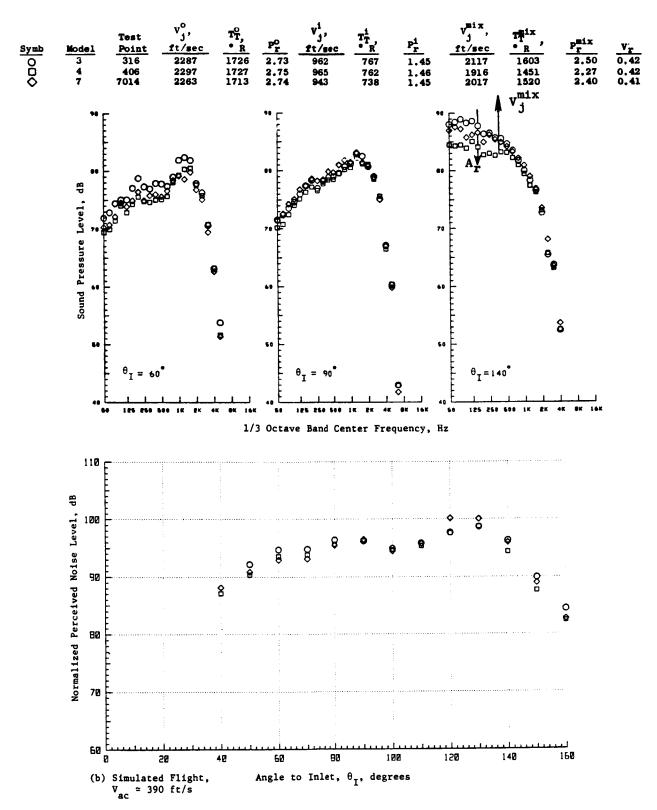


Figure 32. Effect of Area Ratio for a Given Radius Ratio (0.853) and Fixed Inner and Outer Conditions (Concluded).

5.1.3 Influence of Flow Variables on Static and Simulated Flight Acoustics of High-Radius-Ratio Coannular Plug Nozzles

Within this subsection, several aspects regarding the influence of the free stream and the various nozzle flow variables on the acoustics of a typical high-radius-ratio coannular plug nozzle are discussed. Discussion of the acoustic data measured with Model 7 ($R_{\rm r}^{\rm o} \simeq 0.853$, $A_{\rm r} \simeq 0.2$) will cover (1) variation in free-stream velocity holding $V_{\rm j}^{\rm mix}$ or $V_{\rm j}^{\rm o}$ fixed, (2) static and free-jet measurements with variations in $V_{\rm j}^{\rm mix}$ holding $V_{\rm j}^{\rm o}$, $T_{\rm T}^{\rm o}$ and $T_{\rm T}^{\rm mix}$ fixed, (3) static and free-jet measurements of variations with $V_{\rm j}^{\rm o}$ holding $V_{\rm j}^{\rm mix}$, $T_{\rm T}^{\rm o}$ constant, (4) nozzle temperature effects, and (5) inner-to-outer velocity ratio effects.

5.1.3.1 Influence of Free-Jet Velocity for a Constant V_j^0 , T_j^{mix} , and T_j^0

Free-Jet Velocity Influence on PNLmax and PNL60

Several tests with variations in the specific thrust (V_j^{mix}) were run while holding the outer stream velocity (V_j^0) , the outer stream static temperature (T_j^{mix}) , and the mixed stream static temperature (T_j^{mix}) nearly constant. These correspond to Test Points 7101 through 7114 having the test conditions defined in Section 4.2.1. Figures 33 and 34 illustrate the static and simulated flight acoustic test results, respectively, at θ_I that corresponds to the maximum noise angle in the aft quadrant and at θ_I = 60° for the following range of test conditions: V_j^{mix} ranges from 1500 to 2400 ft/sec; V_j^0 = 1700, 2100, and 2500 ft/sec; T_j^0 = 1200° R, T_j^{mix} = 1100° R.

Figure 33 illustrates the influence of flight on the normalized PNL_{max} plotted against 10 log10 (V_j^{mix}/C_a). All test results are scaled to a typical product engine size. The basic trend of the data is that flight decreases the peak angle noise at all the tested conditions. At the lower V_j^{mix} conditions, the influence of flight on peak angle noise reduction is observed to be larger than at the higher V_j^{mix} conditions.

In the forward quadrant where shock noise exists, the influence of flight is to enhance the noise. These results are shown in Figure 34. The range of mixed pressure ratios is $P_{\rm r}^{\rm mix} \approx 1.9$ to 3.5. Roughly, a 3- to 4.5-PNdB forward quadrant lift is observed over the range of test conditions. It is likely that these measurements are influenced by jet mixing noise. However, there should not be any influence of temperature since the static temperatures for all the tests were approximately constant. Later in Section 5.1.4, some of these issues concerning the shock noise will be discussed.

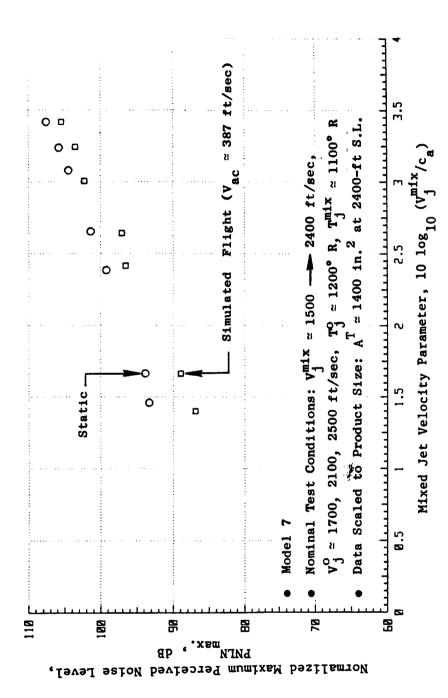


Figure 33. Influence of Free-Jet Velocity on Normalized PNL max.

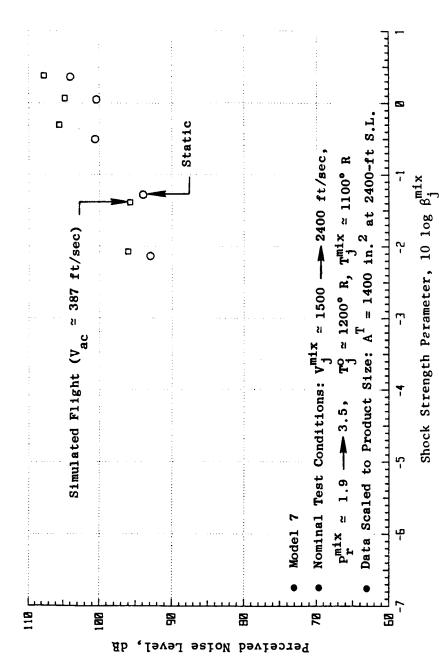


Figure 34. Influence of Free-Jet Velocity on PNL60.

Free-Jet Velocity Influence on Directivity and Spectra

For the full range of test conditions, the static and simulated flight acoustic PNL directivity, SPL spectra at $\theta_{\rm I}$ = 60°, 90°, and the maximum noise angle are shown in Figures 35 through 38. Figures 35 and 36 show the static and simulated flight PNL directivities, and Figures 37 and 38 show the corresponding SPL spectra. In general, the results show coannular plug nozzle trends described earlier in this section. But, it is worth noting the strong influence of flight on shock noise. The results indicate that in-flight shock noise may be influencing the coannular plug nozzle acoustic spectra up to $\theta_{\rm I} = 110^\circ$. This observation is viewed more clearly when comparing the static and simulated flight data on one graph. Figures 39 and 40 compare static and simulated flight results when $V_{\rm j}^{\rm mix} = 2.56$; Figures 41 and 42 compare static and simulated flight results when $V_{\rm j}^{\rm mix} = 2.400$ ft/sec and $P_{\rm r}^{\rm mix} = 2.56$;

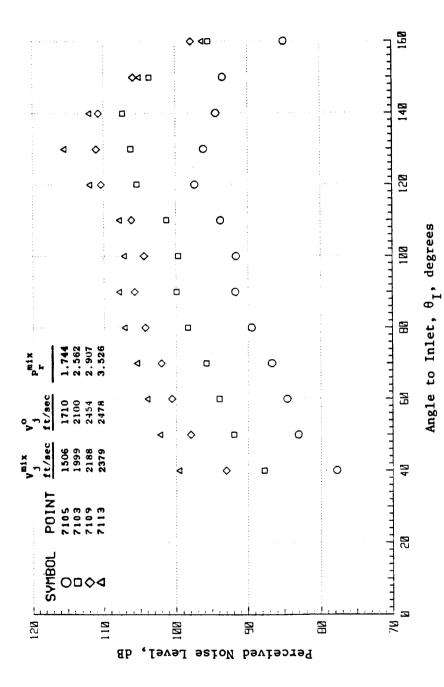
As an example of subcritical flow conditions, Figures 43 and 44 are shown ($V_j^{mix} = 1500 \text{ ft/sec}$, $P_r^{mix} = 1.75$). As a rule, the subcritical tests show a greater range of beneficial flight effects on aft angle directivity and spectra as compared with the supercritical tests shown in Figures 39 to 42. The flight benefit extends beyond $\theta_I = 70^\circ$ on a directivity basis. The simulated flight spectra shown at $\theta_I = 90^\circ$ and the maximum jet noise angle show reduction at nearly all frequencies.

5.1.3.2 Influence of Free-Jet Velocity for a Constant
$$V_j^{mix}$$
, T_j^o , and T_j^{mix}

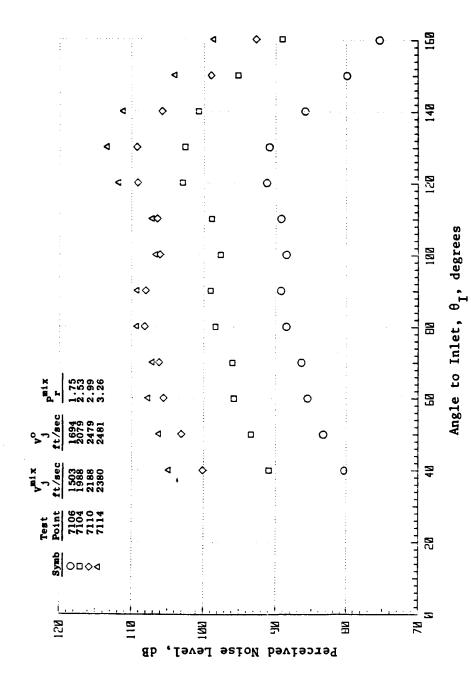
General Results

A series of test points was run where the specific thrust V_j^{mix} , mixed static temperature T_j^{mix} , and the outer stream temperature T_j^{o} were held approximately constant ($V_j^{mix} = 2250$ ft/sec, $T_j^{mix} = 1250^{\circ}$ R, $T_j^{o} = 1275^{\circ}$ R) while the outer/inner stream flows were appropriately varied, as well as the freestream velocity. Figures 45 and 46 show typical static test results, and Figures 47 and 48 show typical simulated flight results at $V_{ac} = 387$ ft/sec.

The results of Figures 45 and 46 show that, although the outer stream velocity ranged from approximately 2300 to 2500 ft/sec (with V_j^{mix} fixed), the jet mixing noise in terms of coannular plug nozzle PNL levels are relatively the same. On a spectral basis, Figure 46 shows that at $\theta_I \approx 60^\circ$ and 90° the shock noise spectral signature differs with the various combinations of P_r^0 and P_r^1 , even though P_r^{mix} (or 10 $\log \beta^{mix}$) is maintained, approximately constant. In simulated flight (Figures 47 and 48), the trends noted above statically are observed also. However, for the test point with the highest outer stream pressure ratio ($P_r^0 \approx 3.3$), greater shock noise amplification is observed. It is also worth mentioning here that for $\theta_I > \approx 120^\circ$ the static



Comparison of Static PNL Directivities at $v_j^{\text{mix}} \approx 1500$, 2000, 2200, and 2400 Ft/Sec. Figure 35.



Comparison of Simulated Flight PNL Directivities at $v_j^{mix} \approx 1500, 2000, 2200,$ and 2400 Ft/Sec. Figure 36.

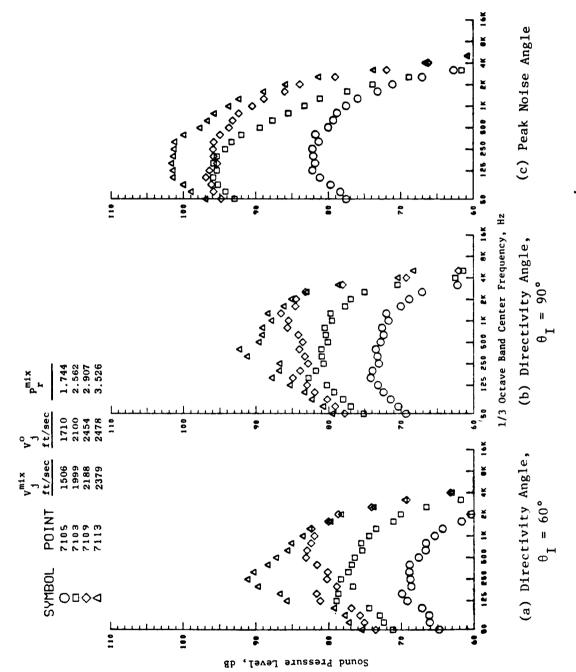
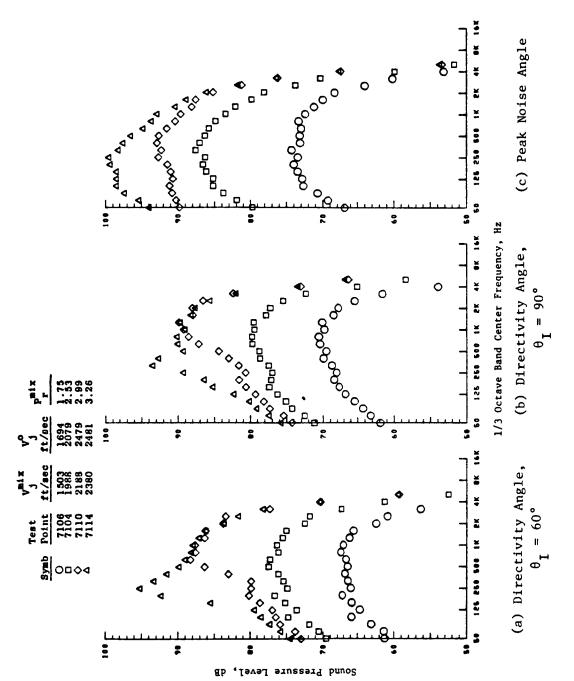
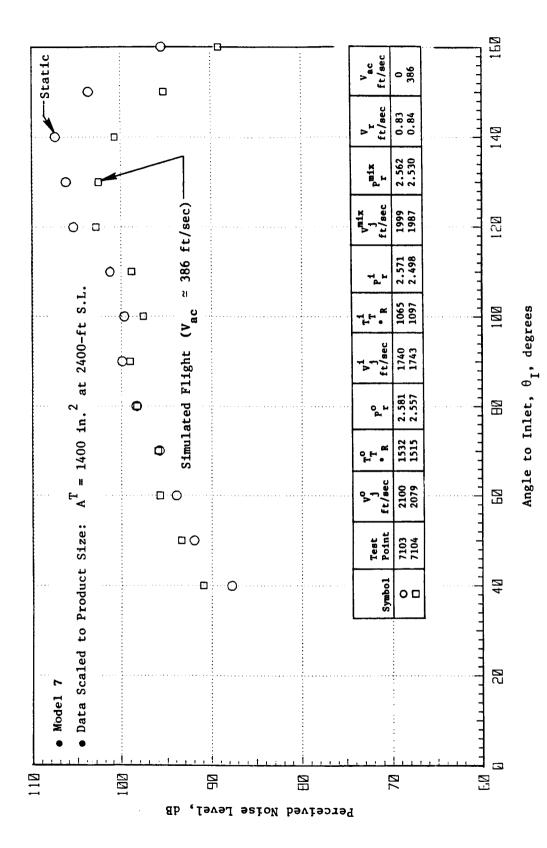


Figure 37. Comparison of Static SPL Spectra at $V_j^{mfx} \approx 1500$, 2000, 2200, and 2400 Ft/Sec.



Comparison of Simulated Flight SPL Spectra at $v_j^{mix} \approx 1500$, 2000, 2200, and 2400 Ft/Sec. Figure 38.



Comparison of Static and Simulated Flight PNL Directivity at $v_{j}^{\text{mix}} \approx 2.000$ Ft/Sec, r Figure 39.

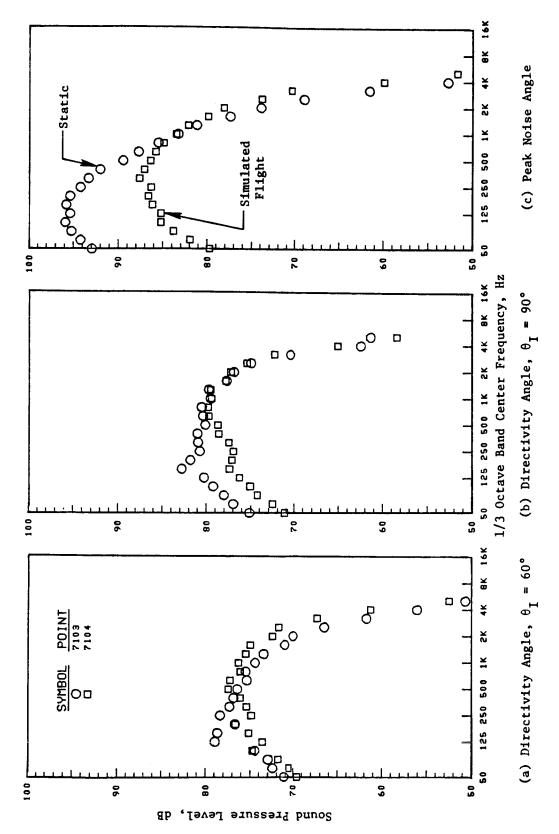
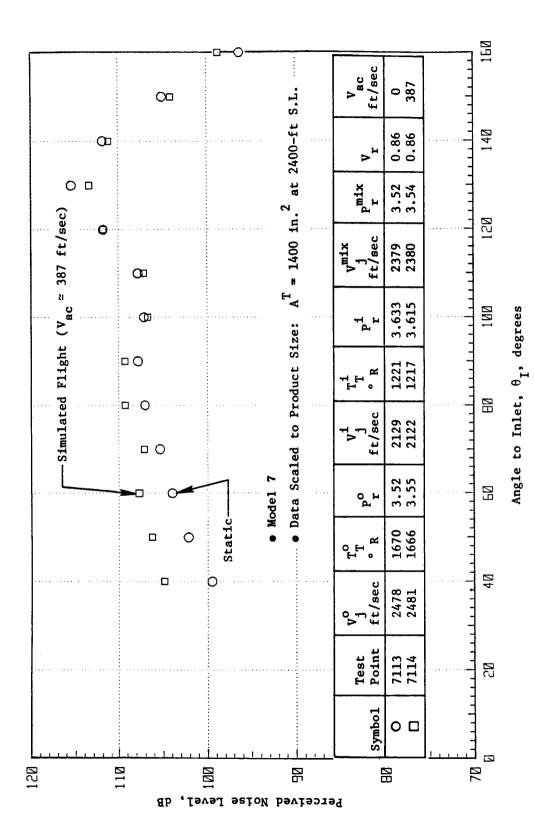
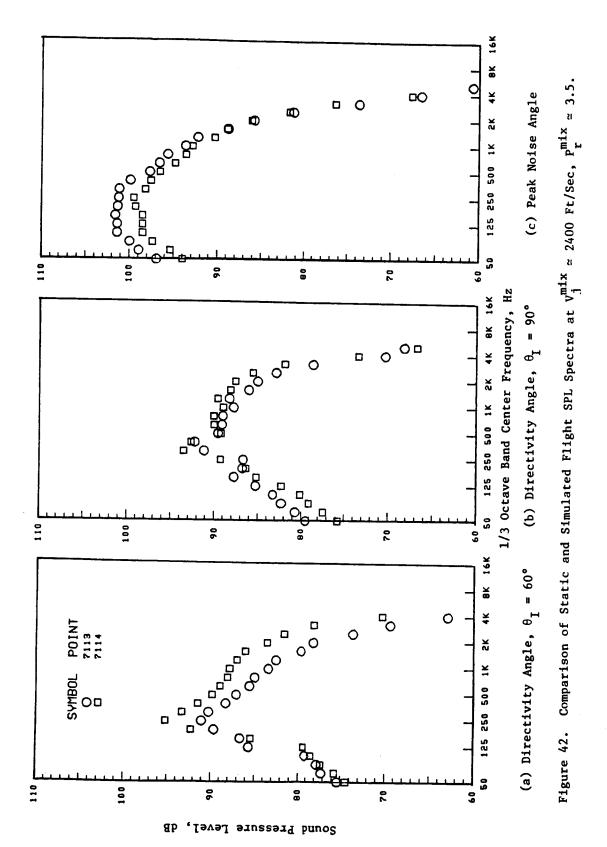
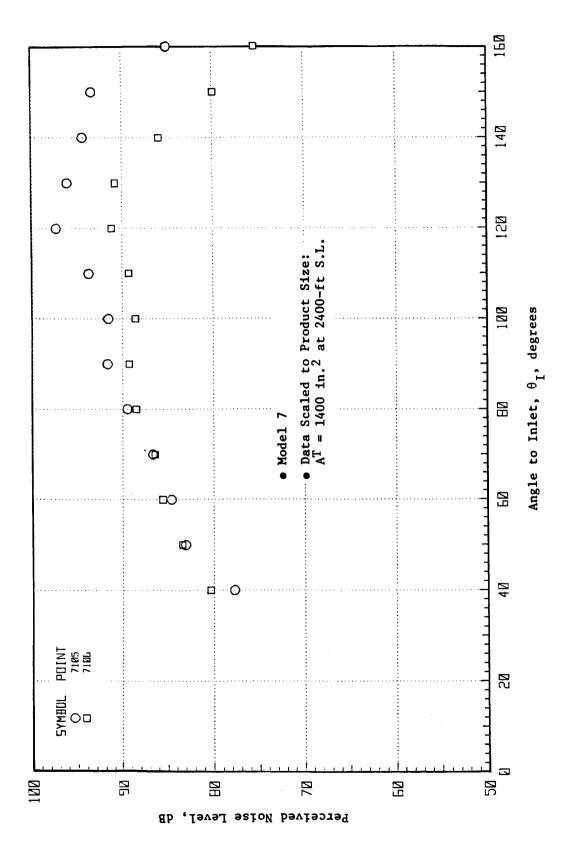


Figure 40. Comparison of Static and Simulated Flight SPL Spectra at $v_{
m J}^{
m mix} \simeq 2000$ Ft/Sec, $v_{
m r}^{
m mix} \simeq v_{
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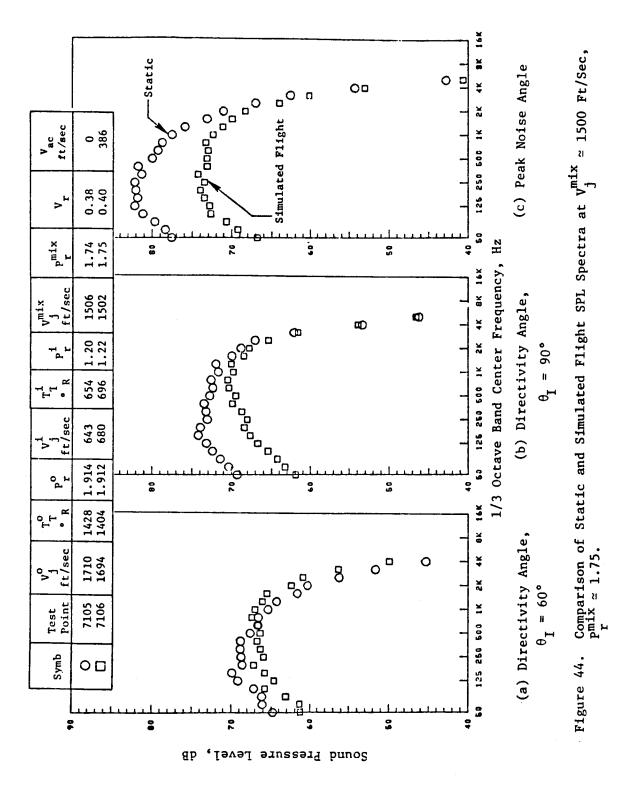


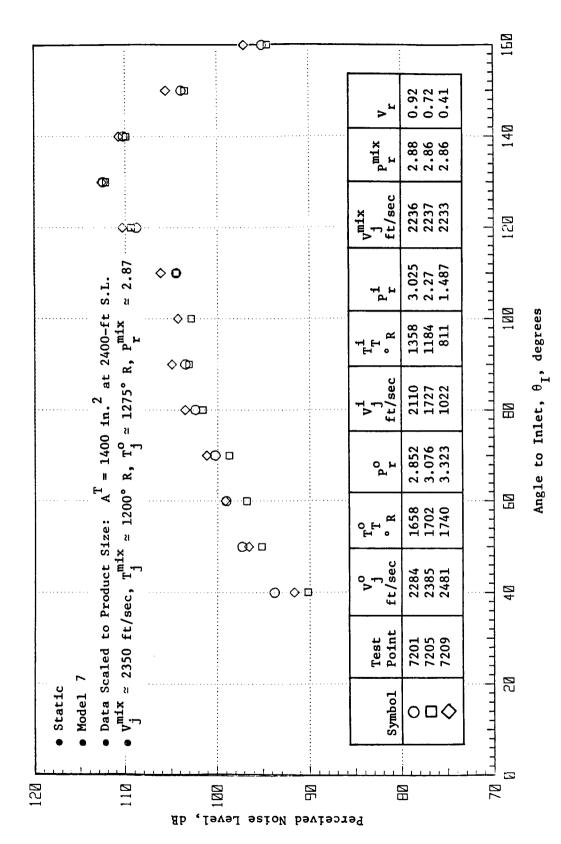
Comparison of Static and Simulated Flight PNL Directivity at $v_j^{mix} \approx 2400$ Ft/Sec, $v_j^{mix} \approx 3.5$. Figure 41.



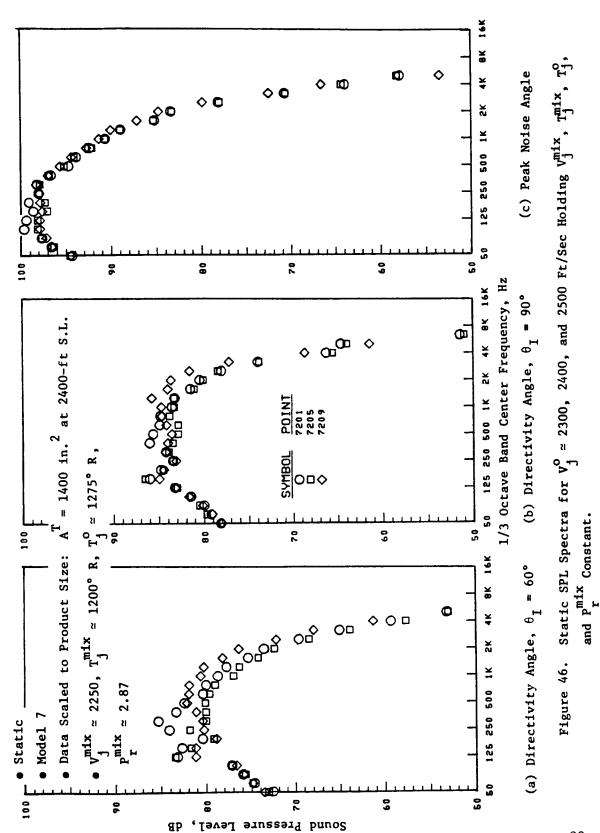


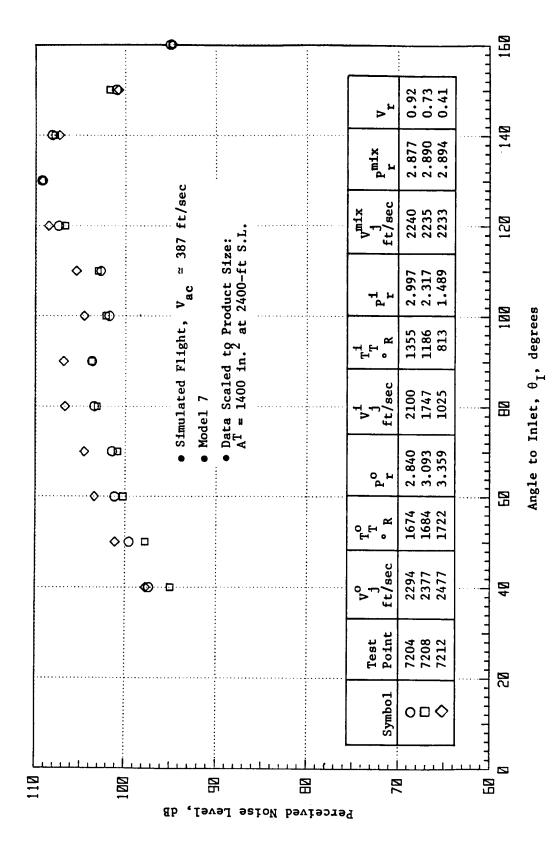
Comparison of Static and Simulated Flight PNL Directivity at $v_j^{mix} \approx 1500$ Ft/Sec, $v_r^{mix} \approx 1.75$. Figure 43.



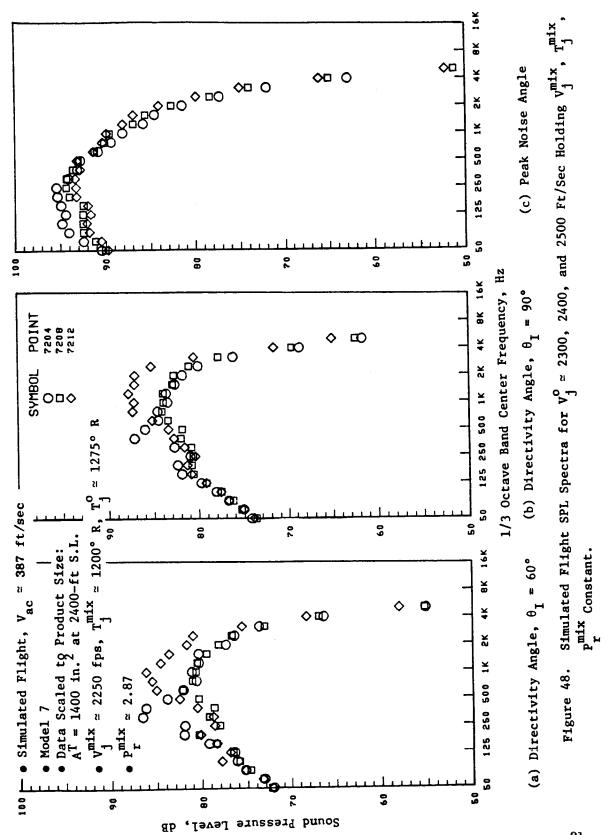


Static PNL Directivity for $V_j^o \approx 2300$, 2400, and 2500 Ft/Sec Holding V_j^{mix} , V_j^o , and V_r^{mix} Constant. Figure 45.





Simulated Flight PNL Directivity for $V_c^o \approx 2300$, 2400, and 2500 Ft/Sec Holding V_j^{mix} , I_j^o , and I_r^m Constant. Figure 47.



and simulated flight aft angle overall noise levels remained relatively insensitive to inner-to-outer velocity ratio (V_r) changes from 0.41 to 0.92. Further studies on velocity ratio effects are discussed in Subsection 5.1.3.5, but the general trend is that for a low inner-to-outer stream area ratio (A_r) coannular plug nozzle suppression is maintained over a broad range of V_r at a fixed specific thrust.

Variation With Free-Jet Velocity

During this test series, acoustic free-jet data were obtained with Model 7 at $V_{ac} \simeq 0$, 150, 300, and 400 ft/sec. Figures 49 and 50 illustrate these results. Shown in Figure 49 is the PNL directivity at the four free-jet speeds. Figure 50 shows the corresponding comparison of the spectral characteristics at $\theta_I = 60^\circ$, 90°, and at the maximum noise angle. The nozzle flow conditions are $V_j^{mix} \simeq 2250$ ft/sec, $P_r^{mix} \simeq 2.87$, and $V_j^0 \simeq 2400$ ft/sec, and $V_r \simeq 0.72$.

Figure 49 shows that as the flight velocity is increased the aft angle noise progressively decreases. However, in the forward quadrant, an opposite trend is observed, i.e., as the flight velocity is increased, so does the forward quadrant noise. From the spectral results shown in Figure 50, it is observed that as forward speed is increased the spectral level at θ_{max} decreases over the frequency range of 50 to 4 kHz. Thereafter, the classical flight effect is not observed. The maximum effect of flight is observed at the lower frequencies. At $\theta_{\rm I} \simeq 90^\circ$ and 60° , observations show that at the lower frequencies (< 2 kHz), the jet noise decreases with an increase in forward speed. At the higher frequencies, an amplification of coannular plug nozzle noise is observed. This is due to the shock-noise forward quadrant amplification with an increase in flight velocity. Other flight tests performed with $V_j \simeq 2300$ and 2500 ft/sec yielded acoustic results similar to the above discussed data.

5.1.3.3 Variation in V_j^0 Holding V_j^{mix} , T_j^{mix} , T_j^0 Constant

In order to study the effects of outer stream velocity on coannular plug nozzle acoustics, a series of test points were run where V_j^0 was varied with V_j^{mix} , T_j^{mix} , and T_j^0 held approximately constant. Under these conditions, it is of interest to examine the effect on the static and flight measured high frequency noise spectra.

Static Test Results

Figures 51 through 54 illustrate the results for two specific thrust conditions ($V_j^{mix} = 1900$ and 2250 ft/sec) for outer stream velocities of V_j^0 2000, 2100, 2200, 2300, 2400, and 2500 ft/sec.

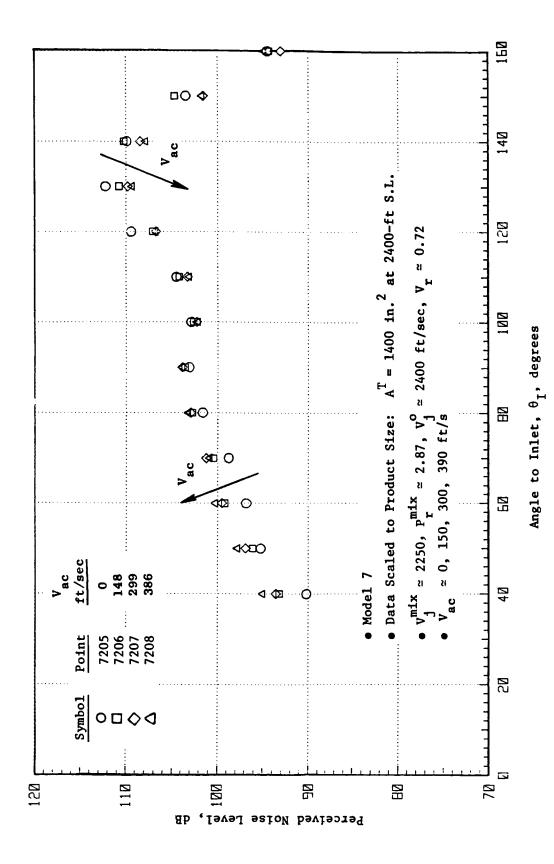
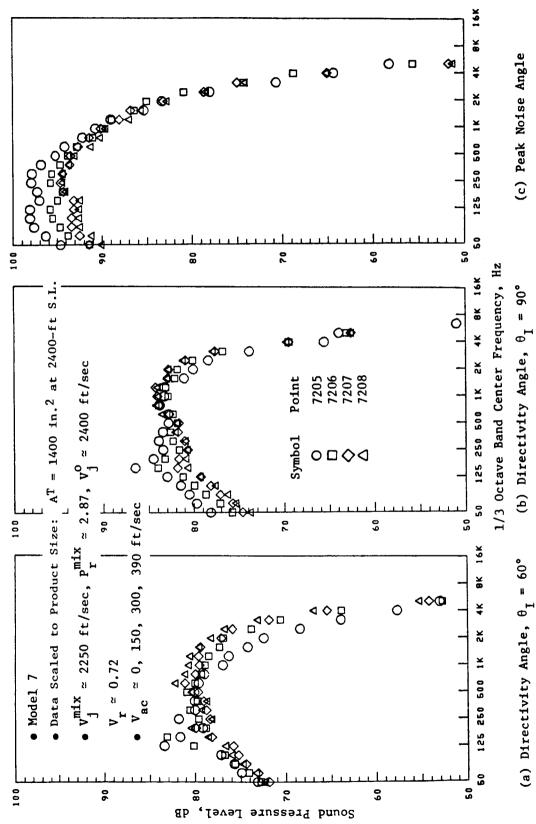
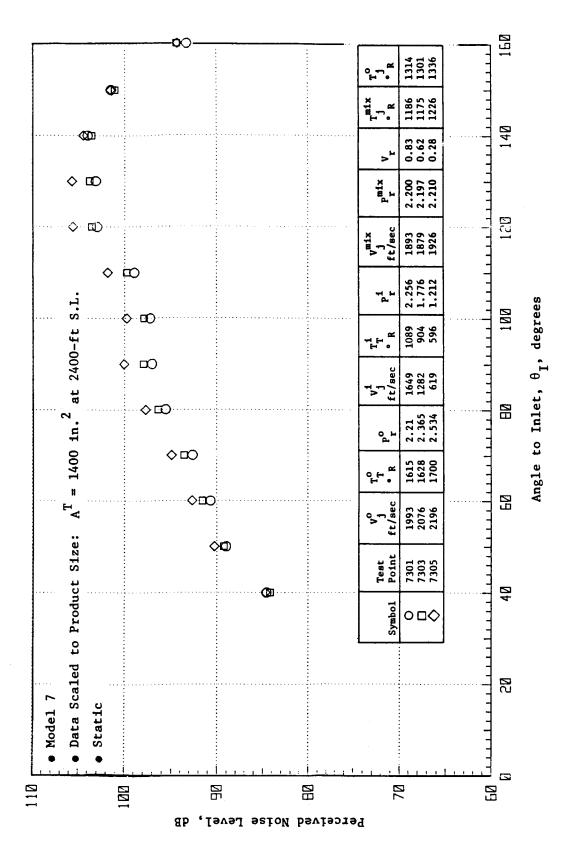


Figure 49. Influence of Free-Jet Velocity on Coannular Plug Nozzle PNL Directivity.

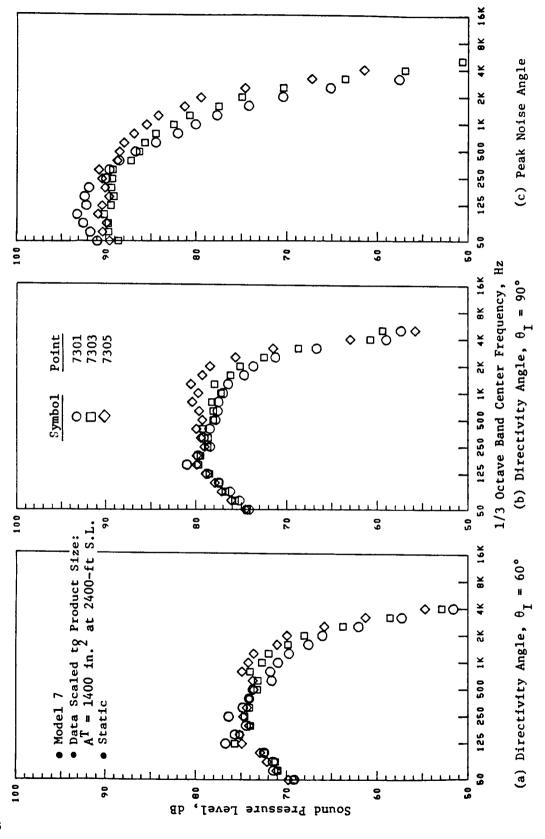


Influence of Free-Jet Velocity on Coannular Plug Nozzle SPL Spectra.

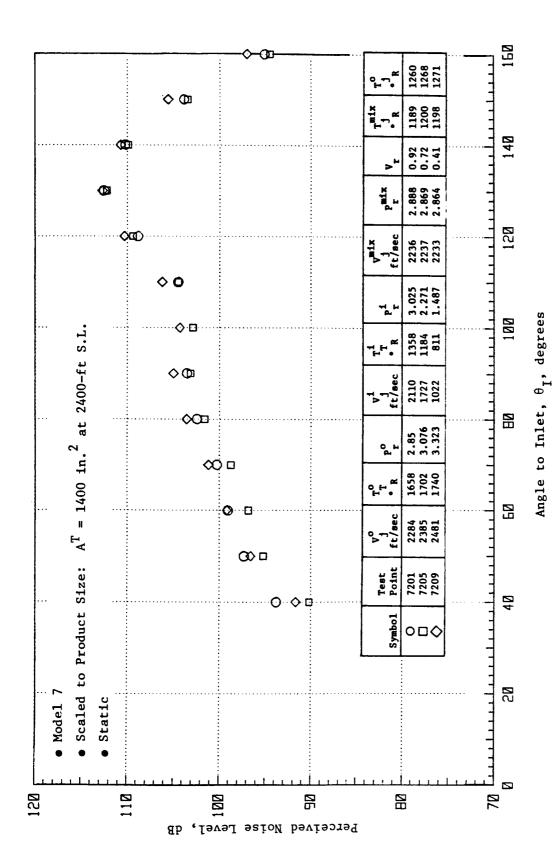
Figure 50.



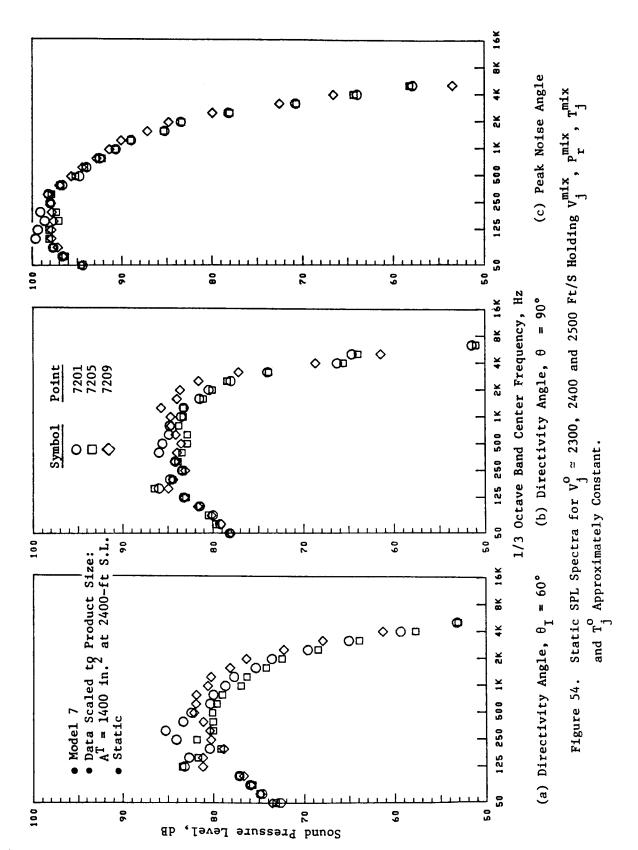
Static PNL Directivity for $v_j^o \approx 2000$, 2100 and 2200 ft/s Holding v_j^{mix} , v_j^{mix} and v_j^o Approximately Constant. Figure 51.



 T_4^{m1x} Static SPL Spectra for $v_j^o \approx 2000$, 2100 and 2200 ft/s Holding v_j^{mix} , v_m^{mix} and T_j^o Approximately Constant. Figure 52.



Static PNL Directivity for $V_j^o \approx 2300$, 2400 and 2500 ft/s Holding V_j^{mix} , P_m^{ix} , T_{j}^{mix} and T_{j}^{o} Approximately Constant. Figure 53.

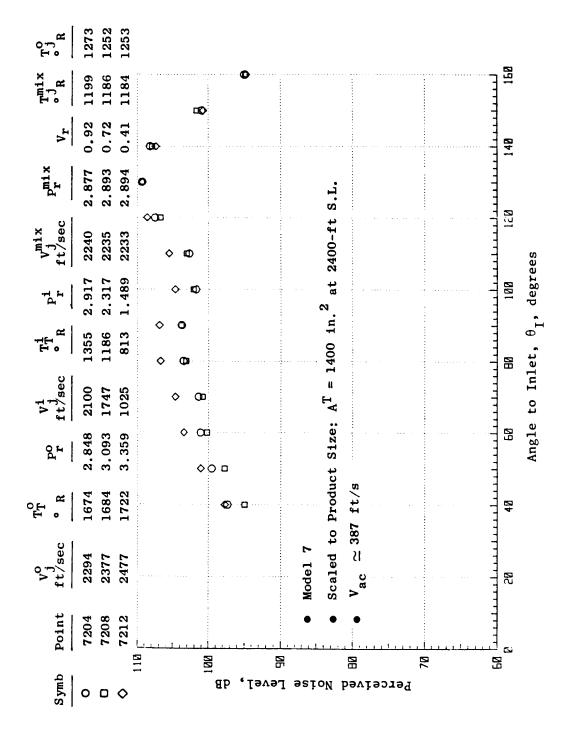


Shown in Figures 51 and 52 are the PNL directivity and SPL spectra for the case where $V_i^{mix} \approx 1900$ ft/sec and V_j^0 is 2000, 2100, and 2200 ft/sec. The results show that while the PNL levels appear to differ the same for the three test points the differences can be due mainly to the variation in V_r and to a smaller extent in V_j^{mix} . On a spectral basis, Figure 52 demonstrates that as V_j^0 increases, the high frequency noise (f > 2000 Hz) increases while the low frequency noise remains nearly the same. As V_j^0 was increased so was P_r^0 , although P_r^{mix} remained approximately constant. The change and/or increase in P_r^0 is observed particularly for the case where P_r^1 was subcritical, and the outer stream pressure ratio, P_r^0 , was at its highest value for this test point ($P_r^0 \approx 2.534$). The amount of influence of shock noise versus jet noise is not easily assessed here. This effect will be discussed separately in Section 5.1.4.

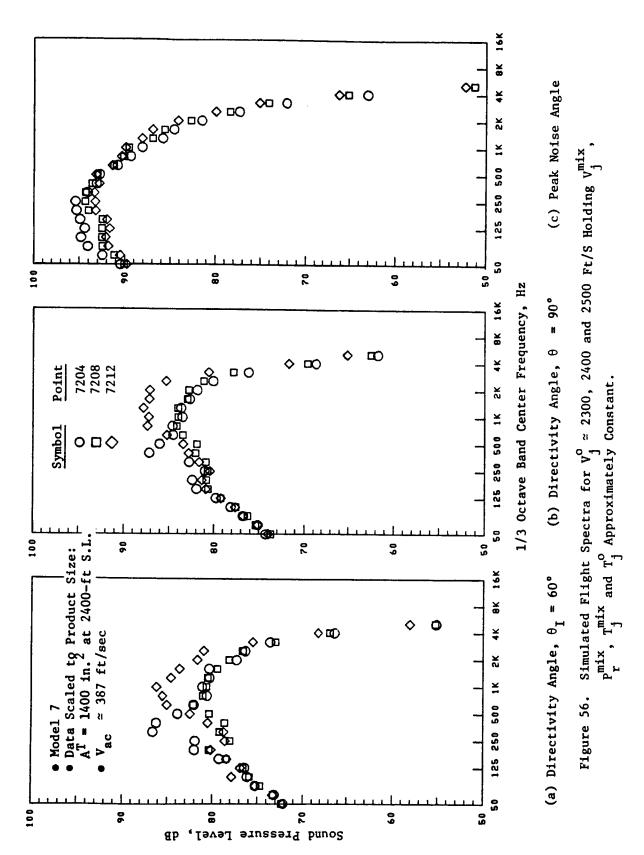
Figures 53 and 54 illustrate the results when $V_j^{mix} \simeq 2250$ ft/sec and V_j^0 is ~ 2300 , 2400, and 2500 ft/sec. At this higher V_j^{mix} condition, the V_j^0 variation does not indicate the strong influence of V_j^0 on the high frequency noise spectrum at θ_{max} as was observed in the previous case. Nonetheless, it is observed that the low frequency noise remained approximately constant and the high frequency noise increased with an increase in V_j^0 . In the forward quadrant, the change in shock noise spectrum with increases in P_r^0 is seen again. However, a significant change and an amplification are noticed when the inner stream is subsonic and the outer stream is fully supercritical ($P_r^0 \simeq 3.3$). For this case, which is representative of a typical AST takeoff sideline condition, it appears that it would be acoustically preferable to run at a reduced outer stream velocity with both streams supercritical rather than at a high outer stream velocity and inner stream subcritical for a given specific thrust and mixed pressure ratio.

Simulated Flight Acoustic Results

The simulated flight acoustic test results are illustrated in Figures 55 and 56. Results shown here are for the case where $V_j^{\text{mix}} \approx 2250$ ft/sec, $P_r^{\text{mix}} \approx 2.88$, and $V_j^0 \approx 2350$, 2400, and 2500 ft/sec, with T_j^{mix} and T_j^0 approximately constant at 1200° and 1275° R, respectively. The results show that increasing the outer stream velocity at a fixed specific thrust does tend to increase the high frequency portion of the θ_{max} spectra. However, PNL_{max} is relatively insensitive to the variations in V_j^0 and is governed by changes in V_j^{mix} and contributions from lower frequency noise. In the forward quadrant and for the case when the inner stream was subsonic and the outer stream supersonic, an appreciable increase in shock noise occurred. The benefit of running both streams at supercritical conditions rather than just the outer stream supercritical (for a given P_r^{mix}) is reaffirmed and even more true in flight. Still, the effect of these changes on the shock structure for the two cases is not understood at this time.



Simulated Flight PNL Directivity for $V_i^o \simeq 2300$, 2400 and 2500 Ft/S Holding $oldsymbol{\mathsf{v}}_{\mathbf{j}}^{\mathsf{mix}}, oldsymbol{\mathsf{r}}_{\mathbf{j}}^{\mathsf{mix}}$ and $oldsymbol{\mathsf{T}}_{\mathbf{j}}^{\mathsf{o}}$ Approximately Constant. Figure 55.



The last illustrative example of test results for this subsection is presented in Figures 57 and 58. These figures compare the static and simulated flight acoustic measurements at $V_j^0 \approx 2387$ ft/sec, $V_j^{mix} \approx 2250$ ft/sec, $P_r^{mix} \approx 2.87$, and $V_r \approx 0.72$ (typical VCE takeoff sideline conditions). The data show that the aft quadrant jet noise is reduced and the forward quadrant shock-associated noise is amplified with increase in flight velocity. By and large, a significant reduction in low-frequency jet noise occurs with flight, offset to some extent by an increase in the high frequency shock-associated noise.

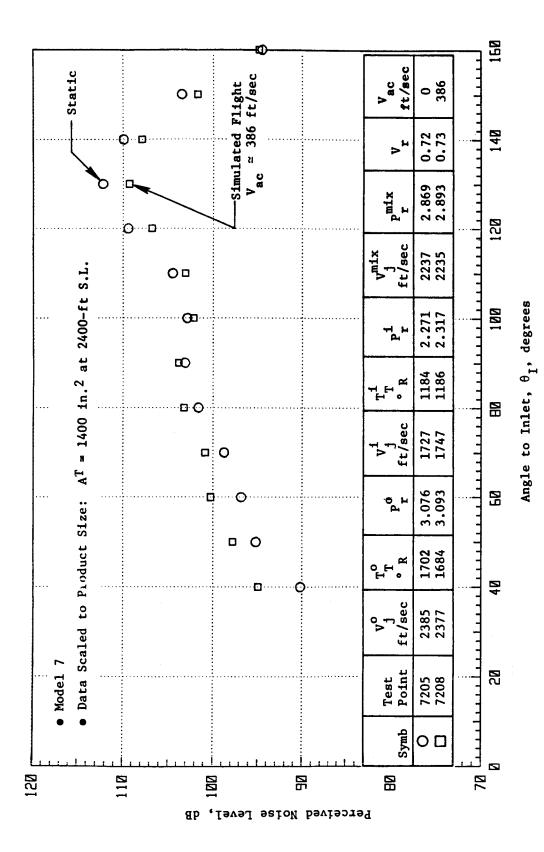
5.1.3.4 Temperature Effects

In order to assess the influence of temperature on coannular plug nozzle static and simulated flight acoustics, a few test points were run in which the outer stream velocity V_j^0 , the inner stream static temperature T_j^i , and the inner stream pressure ratio P_j^i were fixed and the outer stream temperature T_j^i was varied. These tests were performed on Model 7. The nominal conditions were as follows:

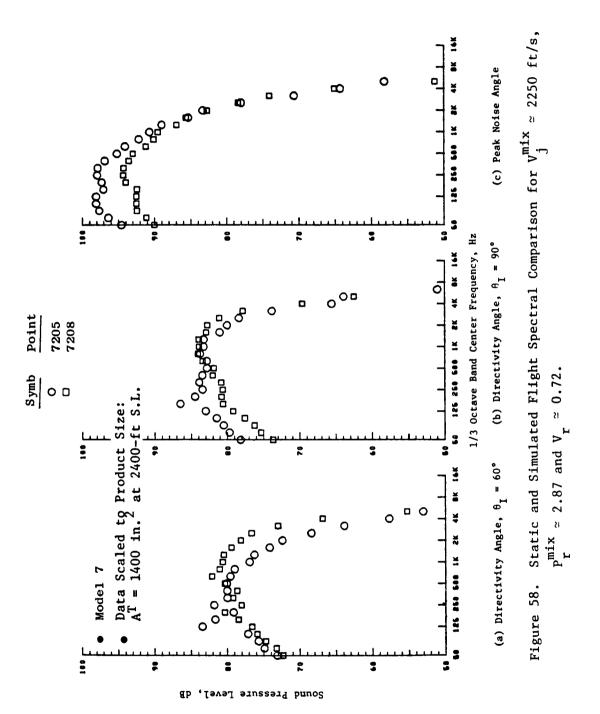
Test Point	V ^o j, ft/sec	T _T , R	T ^o , R	P _r ^o	V ⁱ j, ft/sec	T_{T}^{i} , R	$\frac{\mathtt{T}_{\mathbf{j}}^{\mathbf{i}}}{\mathtt{R}}$	_P _r _	V ^{mix} ft/sec	T ^{mix} ,	P _r ^{mix}
7501	2000	1200	860	3.05	1340	830	680	2.01	1840	825	2.80
2000	2000	1730	1440	2.08	1350	830	680	2.02	1790	1205	2.03

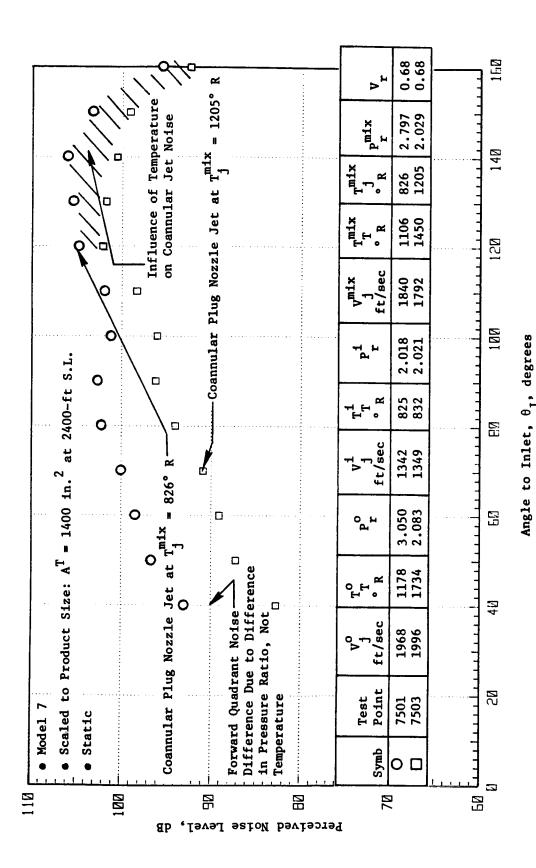
Thus, with V_j^0 fixed and V_j^{mix} differing only slightly, the expected influences due to temperature changes are on the jet mixing noise in the aft quadrant. However, the forward quadrant noise will be influenced significantly by the large change in the pressure ratio. Figures 59 through 61 illustrate the measured results.

Figures 59 and 60 show the static test measurements for the two test points. Figure 59 compares the static directivities. The results show that by increasing the static mixed stream temperature from $T_j^{mix} \approx 826$ to 1205° R, but V_j^{mix} changing from 1840 to 1790 ft/sec, alters the angle of maximum noise from $\theta_I = 140^\circ$ to 120°. Associated with this shift in θ_{max} , there is also a reduction in PNLmax of ≈ 4 dB. An 80 log10 V_j^{mix} and a 20 log T_j^{mix} -type of correlation could account for this change in PNL. In the forward quadrant, there exists a 10-dB noise reduction. Based on a 40 log β^{mix} , there would be a predicted 13.8 dB reduction. Since there is also a significant amount of jet mixing noise in the forward quadrant, it is expected that the jet mixing noise is holding up some of the forward quadrant noise levels. To better illustrate the temperature effect on the aft quadrant noise, SPL spectra results at $\theta_I = 130^\circ$, 140°, and 150° are presented in Figure 61. Figure 60 illustrated the influence of temperature at θ_{max} , but θ_{max} was different for the two test conditions. Figure 61 compares the results at equal values of

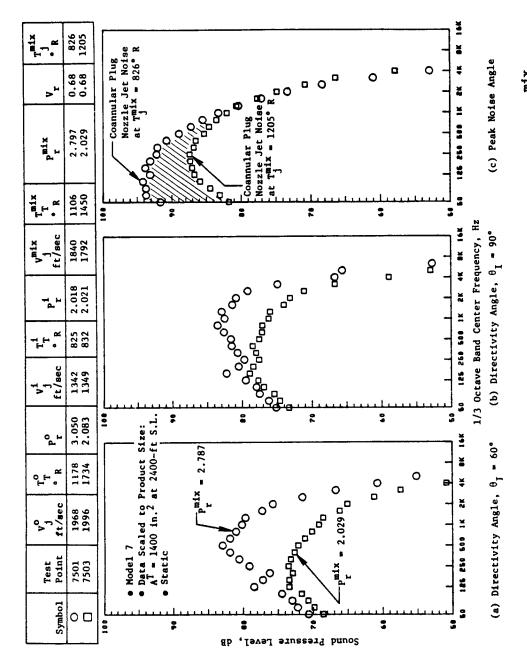


Static and Simulated Flight PNL Directivity Comparison for $V_4^{mix} \simeq 2250$ ft/s, $P_r^{mfx} \approx 2.87$ and $V_r \approx 0.72$. Figure 57,

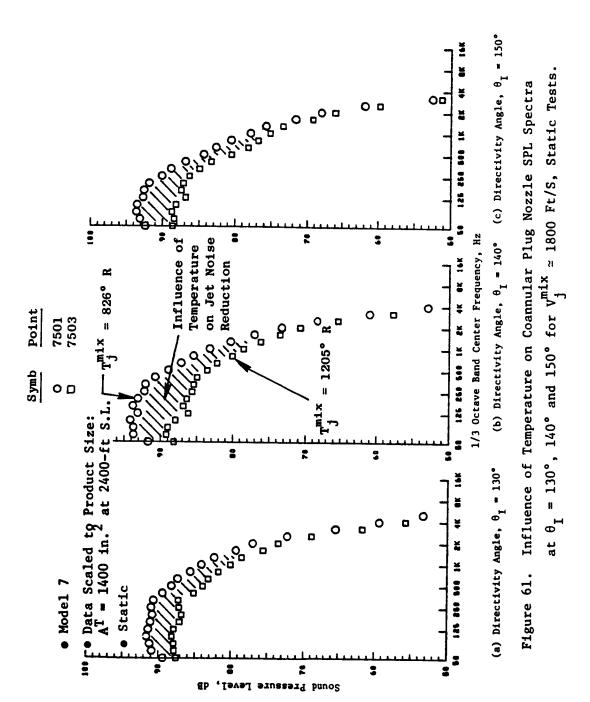




Influence of Temperature on Coannular Plug Nozzle Jet Noise Directivity for $v^{mix} \simeq 1800$ Ft/Sec, Static Tests. Figure 59.



Influence of Temperature on Coannular Plug Nozzle SPL Spectra for $v_1^{mix} \approx 1800$ Ft/Sec, Static Tests. Figure 60.



 $\theta_{\rm I}$ and it indicates that the reduction in the aft quadrant jet noise is observed at all frequencies and that the relative amount of reduction increases as $\theta_{\rm I}$ increases.

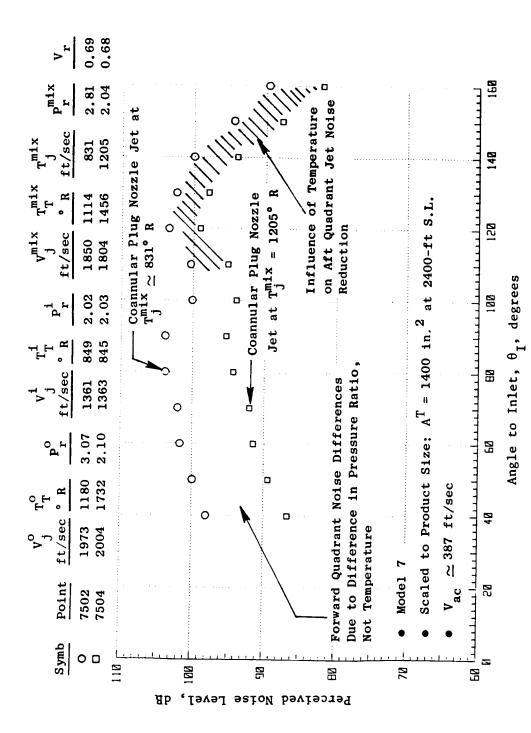
The corresponding simulated flight acoustic results are presented in Figures 62 through 64. The PNL directivity results are shown in Figure 62, which indicates that the angles corresponding to a peak PNL in the aft quadrant are now the same. However, the actual peak angle for the $T_j^{mix} \approx 825^\circ$ R test case now occurs at $\theta_1 \approx 90^\circ$. This is due to the strong shock noise contribution for this test case, whereas the heated test $(T_j^{mix} \approx 1205^\circ$ R) is barely supercritical and has relatively little coannular shock noise. The next observation is that the amount of peak aft quadrant noise reduction due to increased temperatures is somewhat greater in simulated flight than was observed statically. The simulated flight SPL spectra comparisons at $\theta_1 = 130^\circ$, 140°, and 150° that are shown in Figure 64 illustrate the broad range of spectral reductions due to temperature effects, as was observed in the static test results.

5.1.3.5 Effect of Ratio of Inner-to-Outer Velocities

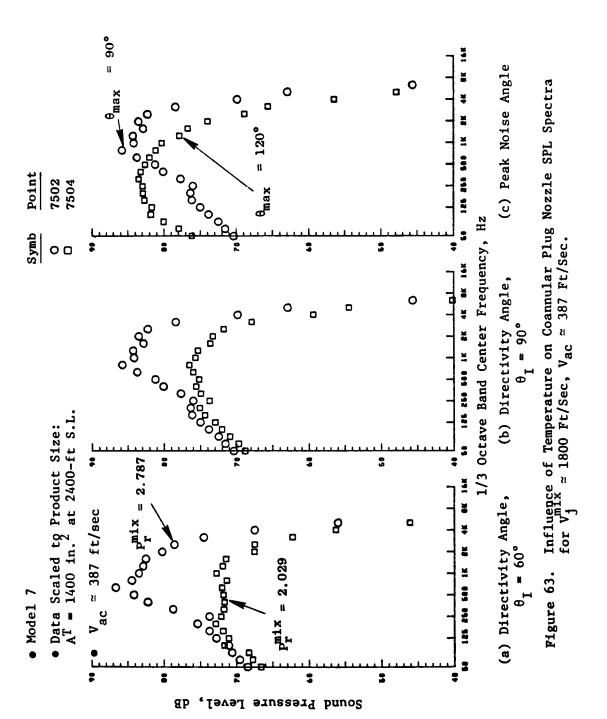
In order to determine the effect of the velocity ratio on the noise characteristics of coannular plug nozzles (Models 3, 4, 6, and 7), tests were conducted where the ratio of inner-to-outer stream velocity was varied. This was achieved by holding the outer stream velocity constant at $V_j^0 \approx 2300$ ft/sec and regulating the inner stream velocity V_j^1 so that velocity ratios of 0.1 to 0.7 were obtained. The normalized PNL_{max} static and flight data (scaled to 2400 ft sideline and 1400 in.² exhaust area) that were measured during these tests are presented in Figure 65(a). An examination of this figure indicates that the measured PNL_{max} for each of the test configurations is a minimum at a velocity ratio between 0.4 and 0.5. The data further indicate that the variation in the PNL_{max} with the velocity ratio is more significant for Models 3, 4, and 7 ($R_r^0 = 0.853$) when compared to that of the high-radius-ratio Model 6 ($R_r^0 = 0.902$).

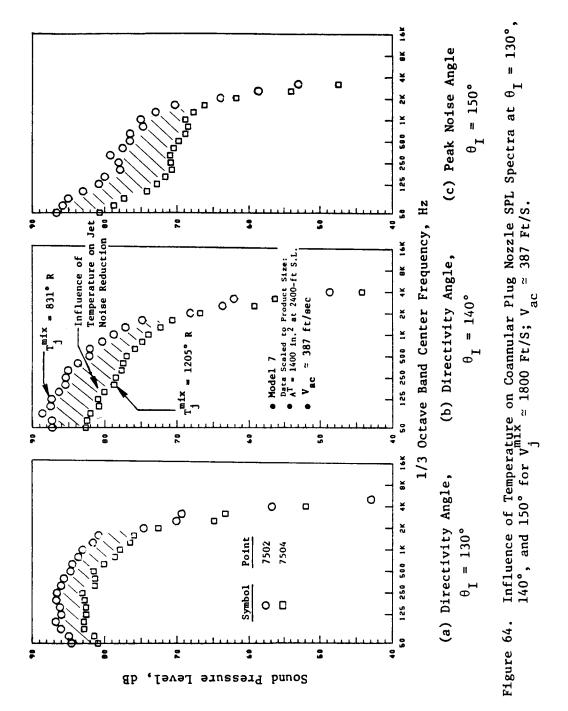
The corresponding mixed stream velocities are presented in Figure 65(b). An examination of the data along with the acoustic data indicates that both the PNL_{max} and the mixed stream velocities vary with the velocity ratio in an identical manner for each of the test configurations. A similar observation has been made during an earlier NASA-supported static study (Reference 2) on a family of coannular plug nozzles. Moreover, after accounting for the different values of V_j^{mix} , it was shown during that study that the noise levels of nozzles with an area ratio <1 are not affected significantly by a variation in the velocity ratio.

For the purpose of making similar analyses with the static and flight data of this study, a linear regression of normalized PNL_{max} as a function of 10 $log~(V_1^{mix}/c_a)$ for each of the four configurations was conducted and



Influence of Temperature on Coannular Plug Nozzle Jet Noise Directivity for $V_4^{mix}\simeq 1800~{\rm Ft/Sec}$, Simulated Flight at V_{ac} ≈ 387 Ft/Sec. Figure 62.





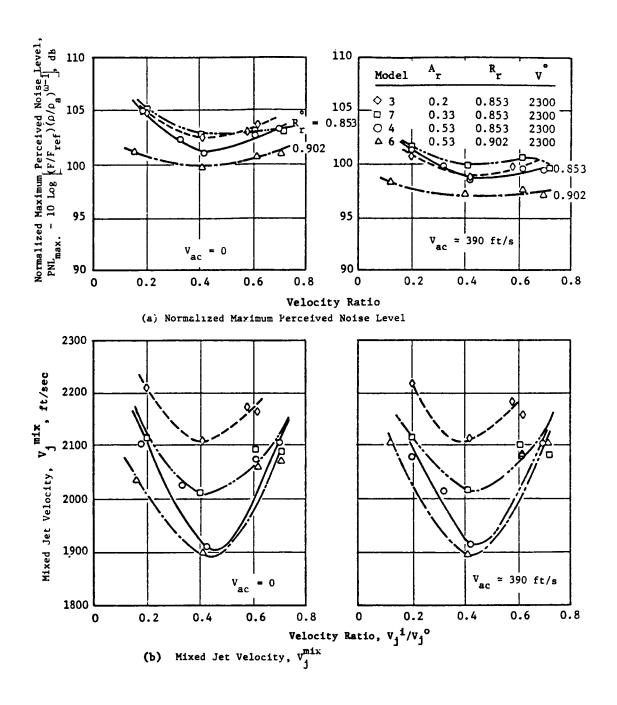


Figure 65. Effect of Velocity Ratio on the Acoustic Characteristics of Coannular Nozzles (Models 3, 4, 6, and 7).

the following expressions obtained (data are summarized in Figure 30 under the area ratio study):

$$PNLN_{max.} = \alpha_1 + \alpha_2 \left[10 \log \left(V_m^{mix} / c_a \right) \right]$$

the coefficients A and B are

$\frac{\texttt{Model}}{\texttt{V}_{\texttt{ac}}}$	3		4		5		6	
	<u>0</u>	400	<u>0</u>	<u>400</u>	<u>0</u>	<u>400</u>	<u>0</u>	<u>400</u>
α_1	85.050	79.520	82.98	90.610	81.960	81.040	82.260	74.740
α_2	6.304	6.956	7.34	3.407	7.016	5.981	7.311	8.272
$\sigma_{\mathbf{XY}}$	0.650	0.810	2.65	2.970	0.930	0.780	0.200	0.270

The data of Figure 65(a) were then normalized to account for the different values of V_j^{mix} . The data so obtained are presented in Figure 66; they indicate that (1) the acoustic characteristics of the tested nozzles are not significantly affected by a change in the velocity ratio, and (2) the considerable variation in the data observed in Figure 67(a) is mainly due to the different values of V_j^{mix} .

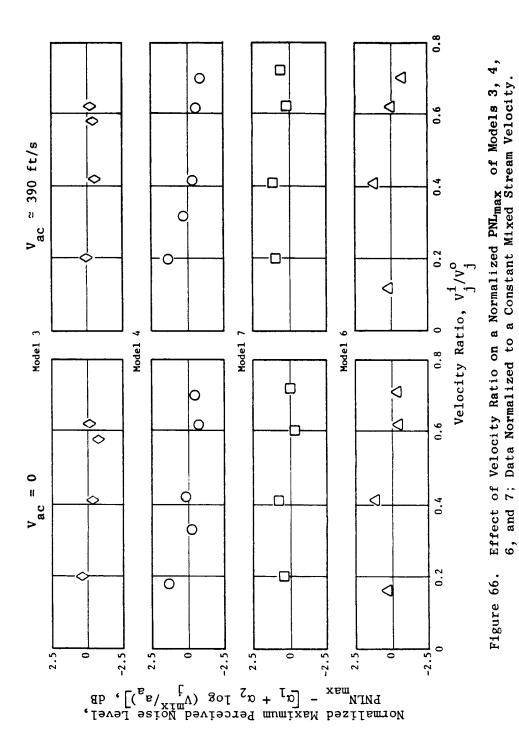
The spectral characteristics (at θ_I = 60°, 90°, 130°) and the PNL directivities obtained during the velocity ratio study with Models 3 and 4 are presented in Figures 67 and 68. Reduction in aft angle acoustic data at velocity ratios that correspond to the smallest V_j^{mix} of the test series is indicated. Similar data were obtained with the other coannular Models 6 and 7 of this study.

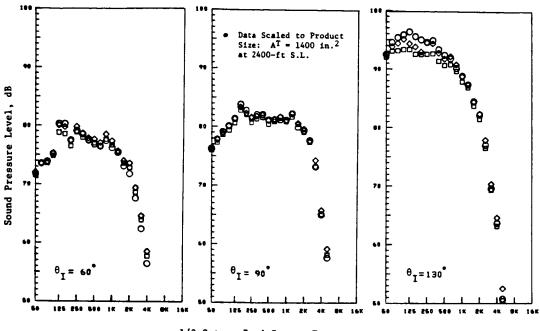
An optimum velocity ratio, at which V_{j}^{mix} is a minimum, can be estimated from the definition of the mixed velocity as follows:

$$V_{j}^{mix} = \frac{\rho^{o}A^{o}(V_{j}^{o}) + \rho^{i}A^{i}(V_{j}^{i})}{\rho^{o}A^{o}V_{j}^{o} + \rho^{i}A^{i}V_{j}^{i}}^{2}$$

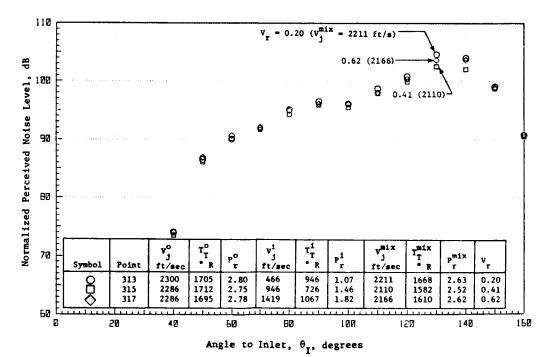
This can be rewritten as

$$V_{j}^{mix} = V_{j}^{o} \frac{1 + \rho_{r} A_{r} V_{r}^{2}}{1 + \rho_{r} A_{r} V_{r}}$$
 where $\rho_{r} = \frac{i}{\rho^{o}}$, $V_{r} = \frac{V_{j}^{i}}{V_{j}^{o}}$, $A_{r} = \frac{A^{i}}{A^{o}}$



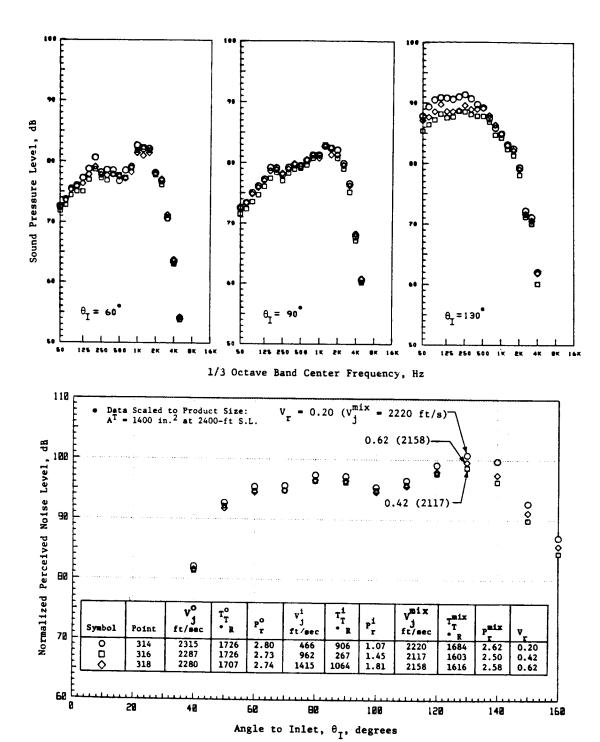


1/3 Octave Band Center Frequency, Hz



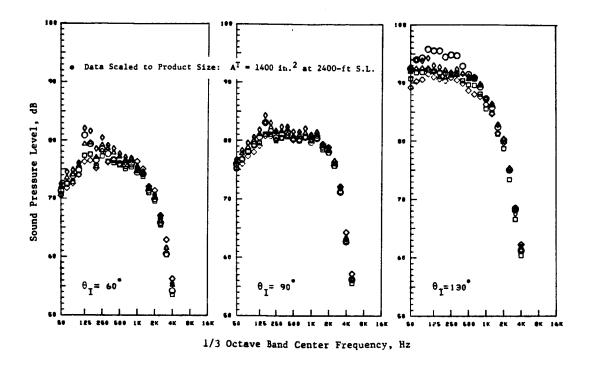
(a) Static

Figure 67. Spectral Characteristics and PNL Directivities of Model 3 Obtained at Different Velocity Ratios with $v_j^o \simeq 2300$ Ft/Sec.



(b) Flight, $v_{ac} \approx 390 \text{ ft/s}$

Figure 67. Spectral Characteristics and PNL Directivities of Model 3 Obtained at Different Velocity Ratios with $V_j^o \simeq 2300$ Ft/Sec (Concluded).



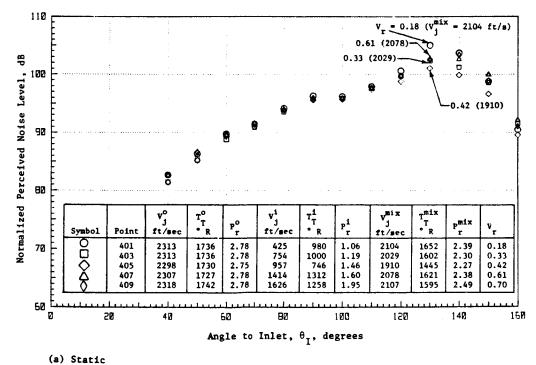
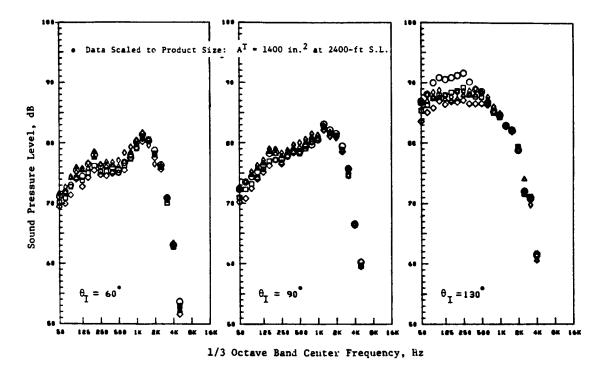
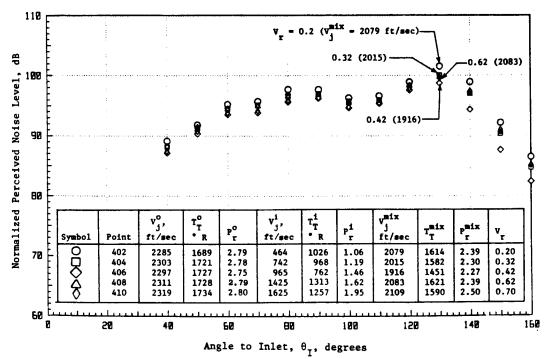


Figure 68. Spectral Characteristics and PNL Directivities of Model 4 Obtained at Different Velocity Ratios with $v_j^o \simeq 2300$ Ft/S.





(b) Simulated Flight, $V_{ac} \simeq 390 \text{ ft/s}$

Figure 68. Spectral Characteristics and PNL Directivities of Model 4 Obtained at Different Velocity Ratios with $V_{j}^{o} \simeq 2300$ Ft/S (Concluded).

For a given V_j^O and A_r and assuming that ρ_r is a constant (over the test series of this study ρ_r ranges from 1.5 to 2.0), the above expression for V_j^{mix} can be shown to result in a minimum value at

$$v_r = \frac{(1 + \rho_r A_r)^{1/2} - 1}{\frac{\rho_r A_r}{r}}$$

which approximates for small values of $\rho_r A_r$ to

$$V_r = 1/2 - 1/8 p_r A_r$$

For example, this expression defining V_r for a minimum in V_j^{mix} will yield $V_r = 0.38$ to 0.45 for the test nozzles of this study ($A_r = 0.2$ to 0.53, and $\rho_r = 1.5$ to 2.0). This confirms the data presented in Figure 65(b).

In conclusion, this study indicates that the noise levels of nozzles with $A_r < 1$ are not significantly affected by a variation in the inner-to-outer velocity ratio. But, for a maximum benefit in the aft angle noise data, the coannular plug nozzles should be operated at a value of V_r which yields a minimum V_1^{mix} .

5.1.4 Special Remarks Regarding Shock Noise and Shock Noise Control for High-Radius-Ratio Coannular Plug Nozzles

In the previous subsections, the general acoustic characteristics associated with coannular plug nozzles have been discussed for static and simulated flight conditions. This subsection describes other significant experimental results obtained in the General Electric anechoic test facility regarding the effectiveness of a convergent-divergent flowpath in shock noise control, the effect of temperature on shock noise, and the influence of downstream shock structure on the resultant shock noise signature and its control.

5.1.4.1 Influence of Contouring for Coannular Plug Nozzles

During earlier static experiments with high-radius-ratio-coannular plug nozzles (Reference .2), it was observed that substantial forward quadrant shock noise reduction was achieved with reference to a conical nozzle. Figure 69 is a summary illustration of this experimental observation. The data show that, at the same shock strength parameter, $\beta^{\rm eff}$ (as defined in the figure), there exists a forward quadrant lift in the static data due to the simulated flight, indicating that the coannular plug nozzle still results in shock cell generated exhaust nozzle noise. As pointed out earlier in flight, this shock-associated noise is amplified in the forward quadrant in simulated flight and thus additional noise reduction could be obtained is the shock cell pattern could be further mitigated or eliminated altogether.

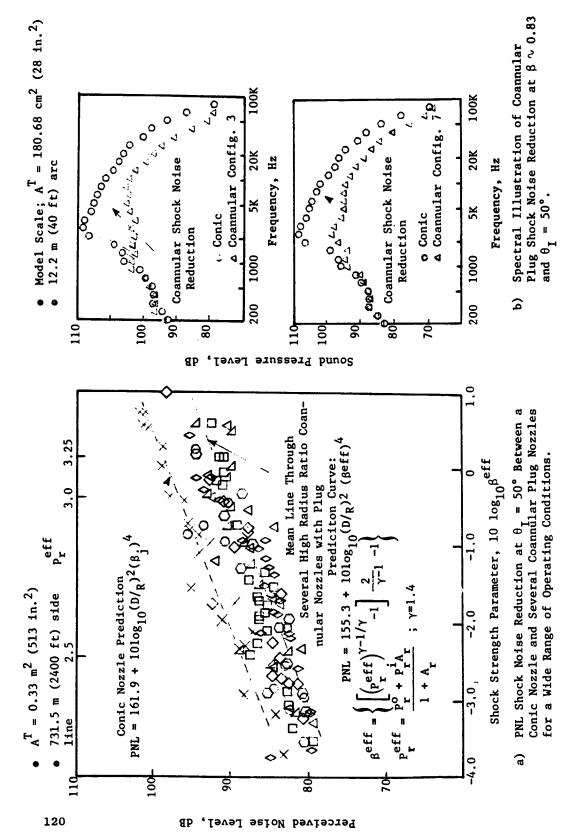


Illustration of Coannular Plug Nozzle Shock Noise Reduction Relative to Baseline Conical Nozzle (Reference 10). Figure 69.

As a first attempt toward obtaining reduced coannular nozzle shock noise, the outer shroud of the coannular plug nozzle was extended, holding the diameter of the shroud constant and obtaining a nozzle throat to exit plane area ratio which would correspond to convergent-divergent design area ratio for shock-free perfect expansion at $P_{\Gamma}^0 \sim 3.2$. No special contouring was performed for this nozzle. To evaluate the convergent-divergent effectiveness of this simplistic design, a series of heated coannular flow static tests was performed. For this test sequence, the area ratio of the nozzle was $A_{\Gamma} \sim 0.2$, and the inner stream was held at a pressure ratio of $P_{\Gamma} \sim 1.6$. For the heated outer flow, the outer pressure ratio was slowly varied from 2.5 to 3.6. Figure 70 illustrates the results.

The results shown in Figure 70 indicate that shock control was not obtained. The PNL values at $\theta_{\rm I}$ = 60° show that with a relatively fine grain outer nozzle pressure ratio variation over the design condition for perfect expansion no significant decrease in the forward quadrant shock noise was observed. These results indicate that in order to obtain perfect expansion for coannular or annular plug nozzle, proper care in the nozzle flowpath contour will be needed. It should be noted here that a subsequent contoured convergent-divergent nozzle has been successfully tested under a separate NASA contract effort.

5.1.4.2 Temperature Influences on Coannular Shock Noise

To determine the effect of temperature on coannular shock cell noise, two series of static acoustic tests have been performed. One test series was performed when both streams were heated, and the second series of tests was performed when both streams were at room temperature - but at nearly the same effective pressure ratio or shock strength condition. The results summarized in Figure 71 show two distinct levels of forward quadrant ($\theta_{\rm I}$ = 60°) shockcell noise. Figures 72 and 73 illustrate PNL directivity and SPL spectra for selected test points at a P_r^{mix} ~ 3.0. These results indicate that forward quadrant as well as aft quadrant radiated noise was substantially reduced for the room temperature jets. It should be noted that room temperature jets are a^{\intercal} considerably lower V_{1}^{mix} and, therefore, result in a lower jet mixing noise level. Nontheless, for circular nozzles, Tanna (Reference 19 has shown for similar temperature differences that there was no shock noise sensitivity with temperature. In actuality, the differences noted above could be temperature effects, jet mixing noise effects or a combination of these two effects on the shock noise. These diagnostic measurements are insufficient to separate each of the effects, but indicate that a temperature influence on the shock noise amplitude does exist. Further work in this important area is needed.

5.1.4.3 Influence of Downstream Shock Structure on the Resultant Shock Noise Signature

In specifying the levels of shock noise for any coannular or annular plug nozzle configuration, the shock region responsible for shock cell associated

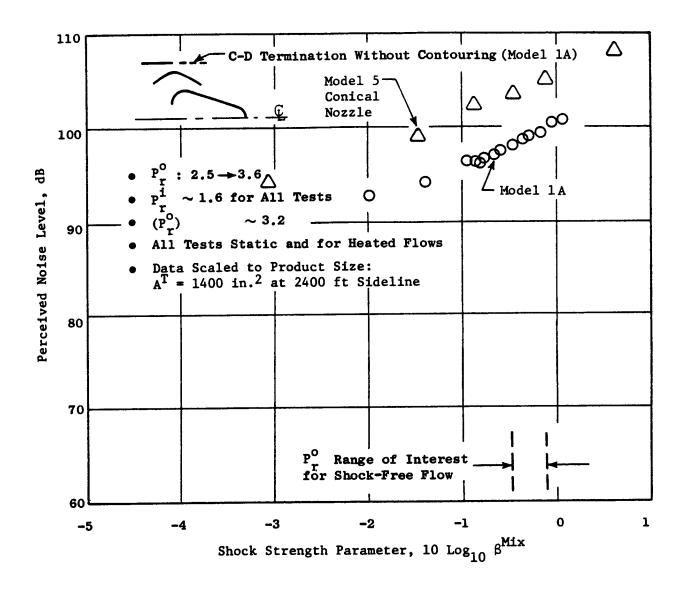


Figure 70. Coannular Plug Nozzle Acoustic Tests with Outer C-D Flow, Directivity Angle, $\theta_{\rm I}$ = 60°.

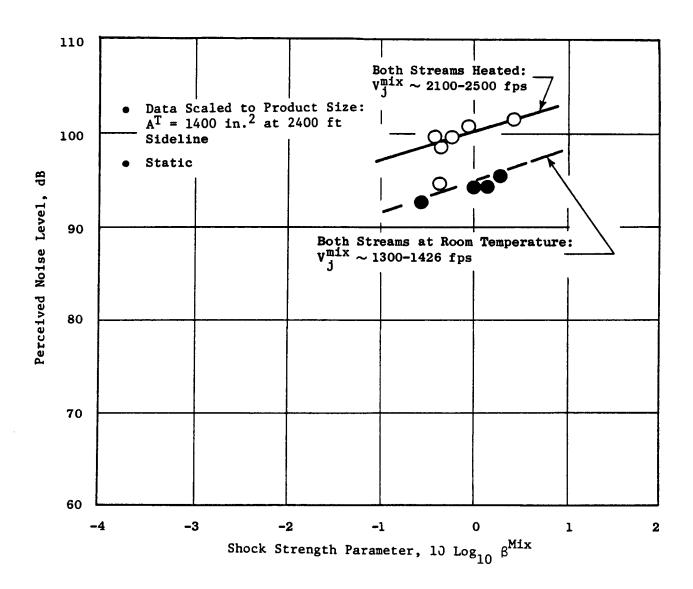
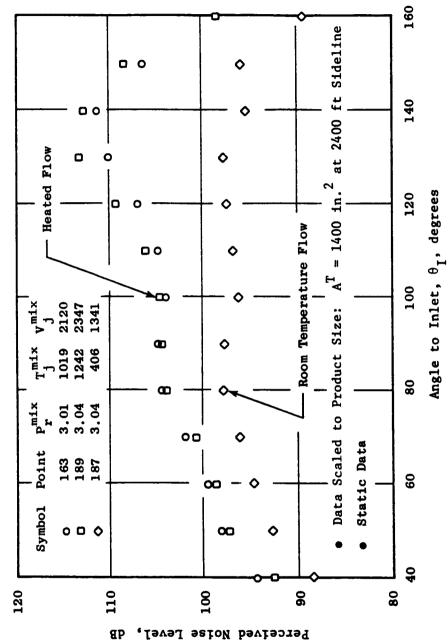
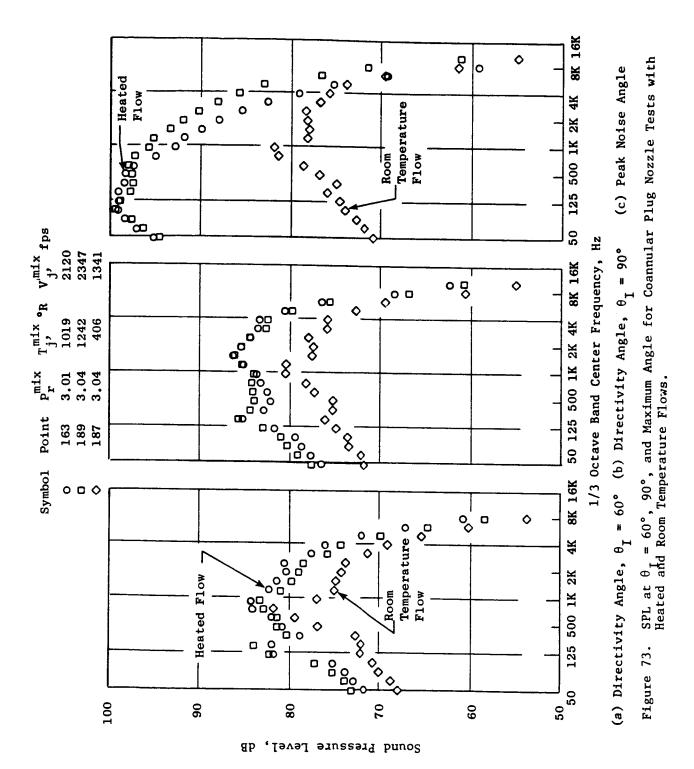


Figure 71. Coannular Plug Nozzle Data at $\theta_{\rm I}=60^{\rm o}$ for Heated and Room Temperature Jets Over the Same Pressure Ratio Range.



PNL Directivity for Coannular Plug Nozzle Tests with Heated and Room Temperature Flows. Figure 72.

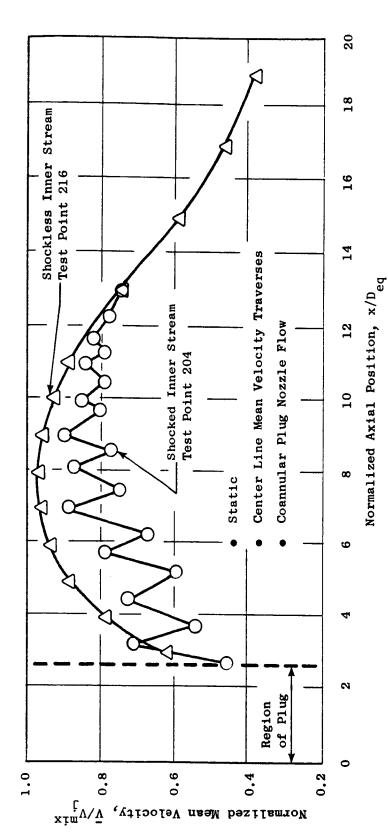


noise levels should be known. Earlier General Electric work (Reference 2) with coannular plug nozzles has indicated that the primary contribution to the shock cell noise may be due to the downstream shock cell structure. This is one reason why the coannular shock noise prediction procedure given in Section 5.3 uses D_{eq} as a characteristic dimension.

Test results obtained during the present investigation shed some light on this subject. Figure 74 shows GE laser velocimeter measurements of the axial mean velocity decay for two coannular plug nozzle experiments (Model 2). The first experiment is for LV measurements when the outer stream and inner streams of the coannular plug nozzle are operating at supercritical pressure ratios; the second experiment is for LV measurements when operating the outer annular plug nozzle at the same pressure ratio as the first experiment but with the inner stream operated at a subcritical pressure The laser velocimeter traces of the first experiment (both streams at supercritical conditions) show a strong series of shock cells, the second experiment (the inner stream subcritical) shows no downstream shock pattern -although the mean velocity is measured to be fully supersonic (well above the inner stream velocity). The flow in the outer stream of both of these experiments was set at the same pressure ratio condition and therefore at equal outer stream shock structures and strength. Figures 75 and 76 illustrate companion static and simulated flight acoustic measurements. The static and simulated flight measurements show significant shock noise reduction or control in the forward and aft quadrant. Although the peak angle noise was not influenced by the observed phenomenon, the flight case duration correction was significantly reduced. For a level flyover calculation, a 2.6 EPNdB reduction was realized. The above results can be rationalized to make the following observations:

- The control of the downstream shock cell structure as well as the shock cell structure in the vicinity of the plug (not all of the shock noise was gone for the second experiment) is important to annular and coannular plug nozzle shock control. This may imply that for coannular plug or annular plug nozzles more than a C-D termination will be required for a total shock control.
- Flow visualization and laser velocimeter measurements will be important to further quantify the physical phenomenon.
- A significant additional flight noise benefit (2.6 ΔΕΡΝdΒ for an unsuppressed annular/coannular nozzle) may be obtainable if total shock noise control can be achieved.

mix T	521	521
$egin{array}{lll} V_{j}^{ ext{mix}} & P_{r}^{ ext{mix}} & T_{r}^{ ext{mix}} \ fps & & e_{R} \end{array}$	2378 3.60 1521	2416 3.42 1621
v ^m tx fps	2378	
$\mathbf{T}_{\mathbf{T}}^{1}$ °R	794	844
$ m p_r^1$	3.21	1.57
$v_{\rm j}^{\rm t}$	2555 3.78 1697 1644 3.21	2563 3.78 1708 1109 1.57
Tr °R	1697	1708
$ m P_{f r}^{f o}$	3.78	3.78
vj fps	2555	2563
Test V _j Point fps	204	216



Laser Velocimeter Tests Showing Influence of Coannular Inner Stream on Downstream Shock-Cell Structure. Figure 74.

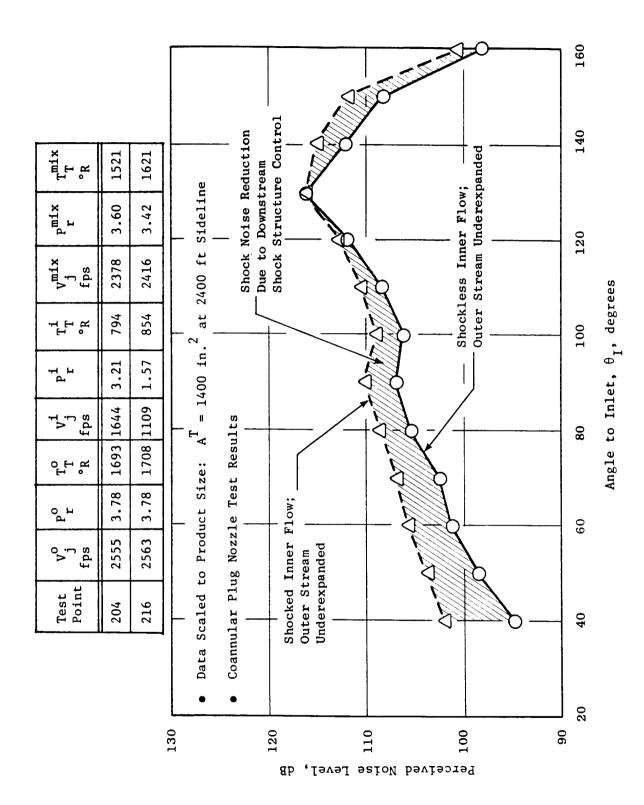
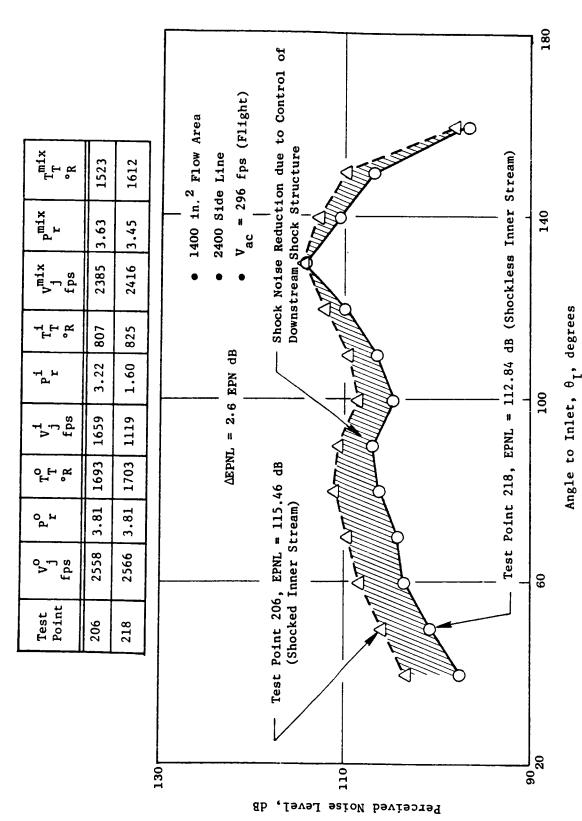


Figure 75. Influence of Coannular Plug Nozzle Shock-Cell Structure on PNL Directivity.



Influence of Coannular Plug Nozzle Downstream Shock Structure on PNL Directivity - Simulated Flight Acoustic Tests. Figure 76.

5.2 LASER VELOCIMETER TEST RESULTS

A laser velocimeter has been employed as a noninvasive diagnostic tool to measure the jet plume characteristics of conic and coannular plug nozzle models. A knowledge of the turbulent mixing characteristics of the jets gives an insight into the noise radiation from the jets. Extensive surveys of the mean and turbulent velocity measurements were performed for a wide range of flow conditions for the scale model nozzles. The following were the objectives for the LV measurements:

- 1. Determine the conic nozzle (Model 5) characteristics, at typical takeoff condition (viz, V_i = 2411 ft/sec, P_r = 3.17, T_T = 1700° R), at static and simulated flight conditions. This would serve as the baseline case for the purpose of comparison with other nozzles.
- 2. Determine the typical jet plume characteristics of an inverted velocity profile coannular plug nozzle operating at typical takeoff condition. Model 2 has been chosen for the purpose of illustration.
- 3. Evaluate the influence of the outer flowpath termination (i.e., convergent or convergent-divergent) on the shock cell structure of the plume and its influence on the shock cell noise.
- 4. Study the aerodynamic characteristics of the coannular plug nozzles operating at off-design pressure ratios (i.e., under- or over-expanded nozzles).
- 5. For a set of prescribed outer and inner stream conditions, evaluate the effect of nozzle area ratio and outer stream radius ratio on the plume development.
- 6. Determine the influence, if any, of the struts in the outer flowpath on the turbulent velocity levels and the mean velocity distribution.
- 7. Study the effect of geometric misalignment of the coannular plug nozzle on the plume asymmetry. This factor arose during the course of LV testing.

5.2.1 Exhaust Plume Characteristics of a Conic Nozzle

Figure 77 shows the shock cell pattern for the conic nozzle (Model 5) at two radial locations (viz, at the centerline and at the tip of the nozzle) for static condition. The aerodynamic test conditions represent typical take-off conditions. One observes that there are seven shock cells within the first 10 diameters of the nozzle at $R/R_2^{\circ} = 0.0$, but only one shock cell at $R/R_2^{\circ} = 1.0$ which is due to the deceleration by the ambient air. The location of velocity maxima and minima at $R/R_2^{\circ} = 0.0$ and 1.0 occurs at the same X/D location indicating that the Mach disks of the shock cells are fairly normal to

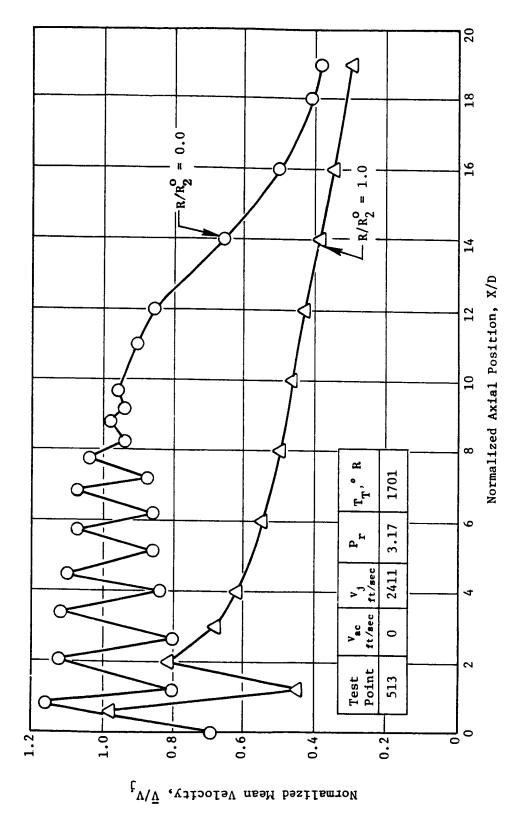


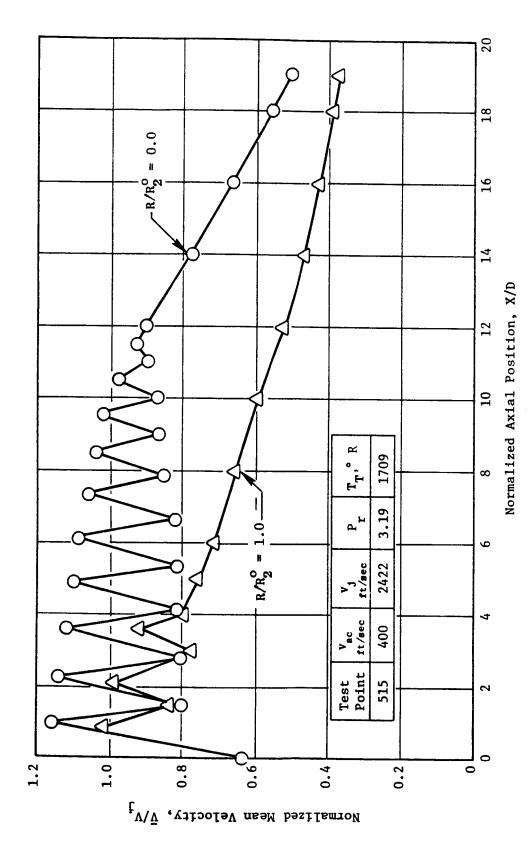
Figure 77. Mean Velocity Variation at Two Radial Locations for Conic Nozzle (Model 5) at Typical Takeoff Conditions (Static).

the jet axis. The shock strength parameter, β (= M_j^2 -1) is equal to 1.0 for P_r = 3.21. The average shock cell spacing for the first three cells equals 1.3 D, but equals only 1.03 D for the last three cells. The Fisher and Harper-Bourne model for shock noise (with L = 1.1 β D) yields an average shock cell spacing of 1.1 D, which seems to be a mean of the measured values.

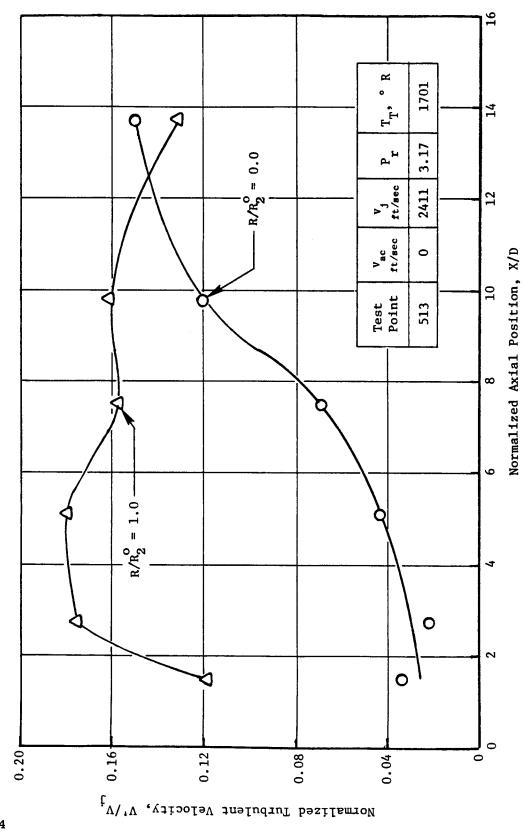
Figure 78 shows the shock cell pattern for the conic nozzle operating at a typical takeoff condition at two radial locations for a free-jet velocity of 400 ft/sec. There are nine shock cells along the jet axis compared to seven shock cells in the static case, and there are two shock cells at $R/R_2^0 = 1.0$ compared to one shock cell in the static case. This indicates the stretching of the shock cell pattern (i.e., the length of the jet plume over which the jet is locally supersonic) which is due to the reduced shear in simulated flight. Also, the average shock cell spacing for the first three cells equals 1.42 D compared to 1.3 D in the static case, indicating that the free jet has stretched each shock cell, as well. However, the Mach disks of the shock cells are still normal to the jet axis and do not seem to be influenced by the free jet. A slower decay is perceived in the mean velocity of the jet at both radial locations downstream of the shock cell pattern for the free-jet case as compared to the static case. Again, this is due to the reduction in shear by the free jet.

Next, the characteristics of turbulent velocity distribution for the conic nozzle for static and free-jet cases are studied. Figure 79 shows the axial variation of the turbulent velocity at two radial locations for the same static case as in Figure 77. Note that within the potential core (i.e., for X < 4D) the turbulent velocity remains within 4% of the jet exit velocity along the nozzle centerline and increases steadily through the transition and the fully developed regions of the jet to a peak value of 15% of Vj at $X \simeq 14D$. But, the turbulent velocity at R/R9 = 1.0 is about 12% of V_1 at one diameter downstream of the nozzle exit plane and rises rapidly to a peak value of about 18% and remains within a variation of 2% for about 10 diameters, then drops after that. The turbulent velocity at the nozzle tip is fairly high compared to the centerline value, since it is in the middle of the shear layer of the jet where the turbulent shear stresses are maximum. Recall that the turbulent shear stress is directly proportional to the square of the turbulent velocity. As one moves downstream, the shear layer widens and, hence, the radial velocity gradients at $R/R_2 = 1.0$ are reduced, resulting in lower turbulent velocity.

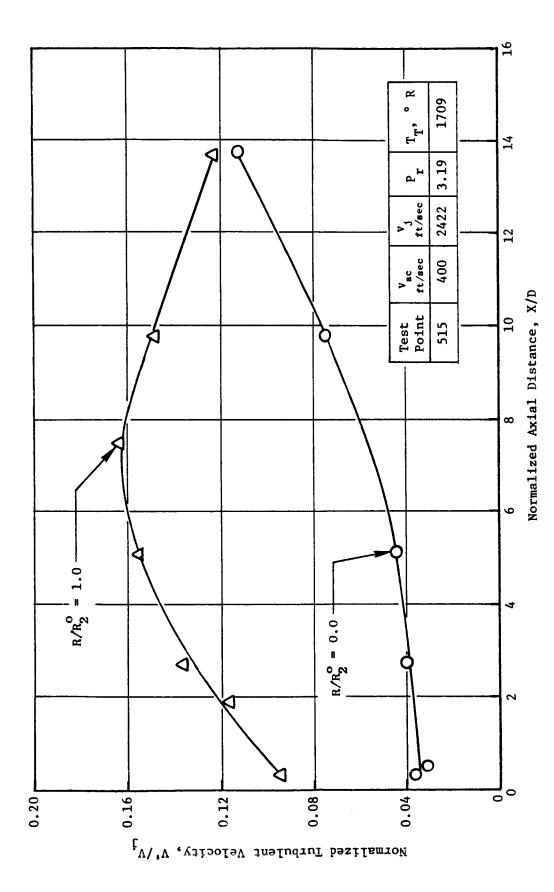
Figure 80 shows the axial variation of the turbulent velocity for the conic nozzle at two radial locations as in Figure 79, in the presence of a free jet at 400 ft/sec. It is worth mentioning that the turbulent velocities at the nozzle centerline are at about the same level as in the static case for locations within the potential core. Nevertheless, downstream of the potential core, the turbulent velocities are lower in simulated flight indicating a reduction in the turbulent shear stresses due to the free jet. At $R/R_2^{\rm O}=1.0$, the turbulent velocity reaches a peak value of 16% of $V_{\rm j}$ in comparison to 18% of $V_{\rm j}$ in the static case – yet another indication of reduced shear due to the free jet.



Mean Velocity Variation at Two Radial Locations for Conic Nozzle (Model 5) at Typical Takeoff Conditions (Flight, V $_{ac} \simeq 400~Ft/Sec)$. Figure 78.



Turbulent Velocity Variation at Two Radial Locations for Conic Nozzle (Model 5) at Typical Takeoff Conditions (Static). Figure 79.



Turbulent Velocity Variation at Two Radial Locations for Conic Nozzle (Model 5) at Typical Takeoff Conditions (Flight, $v_{ac} \approx 400~Ft/Sec)$. Figure 80.

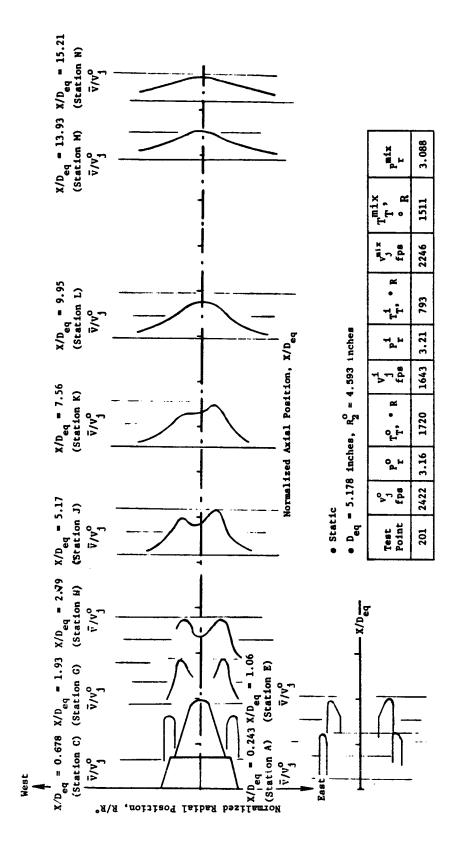
Thus, Figures 77 through 80 indicate that a free jet reduces the shear, thereby reducing the decay of the mean velocity and at the same time reducing the turbulent eddy velocity. The noise radiated by the jet depends both on the mean velocity decay and on the levels of turbulent velocities. Since a free jet reduces both parameters, it cannot be definitely concluded whether a flight enhances noise production or reduces it.

5.2.2 Exhaust Plume Characteristics of a Coannular Plug Nozzle

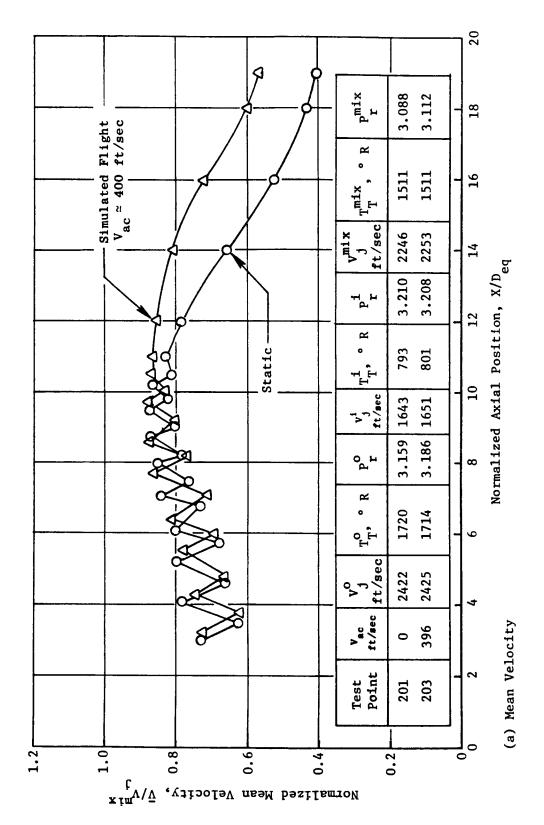
Model 2 is a geometrically scaled version of an AST/VCE coannular plug nozzle (Aⁱ/A^o = 0.2, R_r^o = 0.853) with a C-D termination on the outer flowpath and convergent termination on the inner flowpath but having no struts in the outer flowpath. Model 2 has been chosen to illustrate the typical plume characteristics of dual flow nozzles with an inverted velocity profile. ure 81 shows the plume development for a typical takeoff condition (Test Point 201). The radial variation of the mean velocity at various axial stations is illustrated. Station A (at $X/D_{eq} = 0.243$) is upstream of the inner flow exit plane, and the presence of only the outer stream is noted. At Station C (at $X/D_{eq} = 0.678$), the inner flow has appeared. The inverted velocity profile is clearly demonstrated at Stations G, H, and J (respectively at $X/D_{eq} = 1.93$, 2.79 and 5.17). Also observed is the profile asymmetry about the jet axis at Stations H, J, and K (respectively at $X/D_{eq} = 2.79$, 5.17, and 7.56) which can be traced to the geometric misalignment of the nozzle (Section 5.2.8). For stations downstream of Station K (i.e., for $\rm X/D_{eq}$ > 7.56), the radial profile resembles that of a conic nozzle, thus indicating that the dual flow character is maintained up to an axial distance of approximately 7.5 D_{eq} from the nozzle exit plane.

Figure 82(a) shows the influence of the free jet on the centerline mean velocity distribution for Model 2. Unlike the conic nozzle, the shock pattern is not noticeably stretched by the free jet. This can be attributed to the differences in geometry and the resulting plume development of the two nozzles. There are seven shock cells under simulated flight conditions and eight at static conditions over the same distance, indicating that each shock cell is stretched slightly. For X greater than 10 $D_{\rm eq}$, a slower decay for the simulated flight case is seen when compared to that of the static case. A similar observation was made earlier with the conic nozzle data.

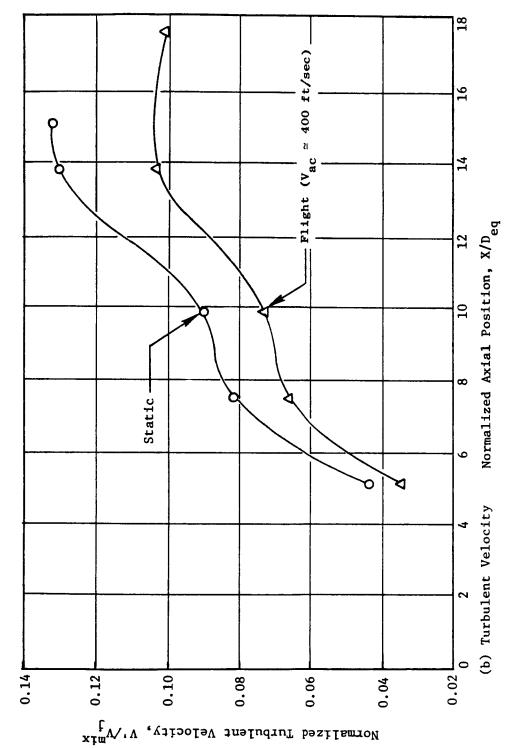
The axial variation of the turbulent velocities of the coannular plug nozzle under static and simulated flight conditions are presented in Figure 82b. Similar to Figures 79 and 80 for the conic nozzle, the data of Figures 82b indicates that the free jet has reduced the fluid shear. Hence, in simulated flight, the turbulent velocities are lower for all axial locations, and the peak turbulent velocity reached is about 3% lower compared to the static case. As in the case of the conic nozzle, the centerline turbulent velocities reach their peak value at X \approx 14 Deg.



Radial Variation of Mean Velocity Profiles at Various Axial Locations Downstream of AST/VCE Model 2 Nozzle. Figure 81.



Velocity Distribution for Coannular Plug Nozzle (Model 2) at Typical Influence of Simulated Flight on the Centerline Mean and Turbulent Takeoff Conditions. Figure 82.



Influence of Simulated Flight on the Centerline Mean and Turbulent Velocity Distribution for Coannular Plug Nozzle Model 2 at Typical Takeoff Conditions (Concluded). Figure 82.

5.2.3 Influence of Outer Flowpath Termination on the Flow Characteristics

AST/VCE Models 2 and 3 are compared in this section to isolate the effect of the outer flowpath termination on the shock structure and the resultant influence on shock cell noise. Models 2 and 3 are identical coannular plug nozzles, except that Model 2 has a convergent-divergent termination on the outer flowpath designed for $P_{\rm r}^0=3.20$ whereas Model 3 has a convergent termination on the outer flowpath.

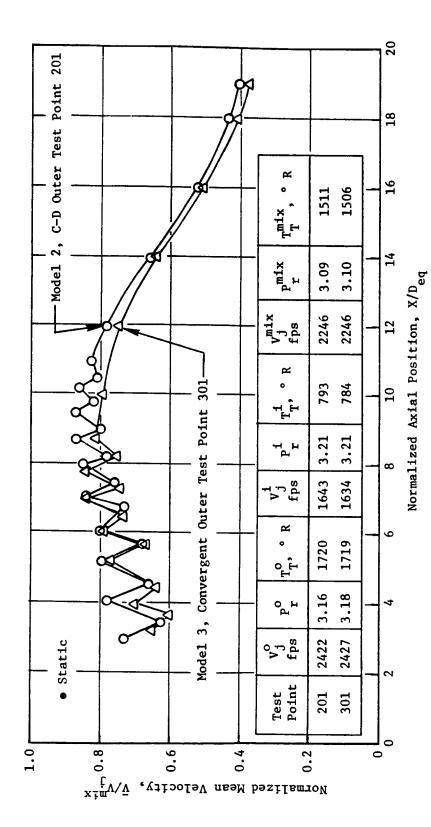
Figures 83 and 84 compare the axial variation of the mean velocity and the shock cell structure for Models 2 and 3 for static and free-jet cases, respectively. The outer stream pressure ratio for both models equals 3.2 which corresponds to the design pressure ratio of Model 2. However, since the inner stream is supersonic ($P_T^1 = 3.21$) and has a convergent termination for both models, the presence of a very strong shock cell pattern can be seen. As noted in Figure 81, the dual flow interaction exists for X < 7.5 D_{eq} over which most of the shock cell pattern exists. Thus, it is not possible to isolate the influence of the outer flowpath termination and only qualitative trends can be extracted. Table IV shows the peak SPL and the corresponding peak shock frequency at $\theta_I = 50^\circ$ for a scaled total flow area of 1400 in. 2 and an extrapolated sideline distance of 2400 ft. $\theta_I = 50^\circ$ is chosen because shock noise is the dominant noise component in the front quadrant.

Table IV indicates that for the static and free-jet cases, Model 3 generates more shock noise than Model 2 indicating that, although a C-D termination on the outer flowpath did not eliminate shock noise, it resulted in a small reduction in the shock noise compared to a convergent termination on the outer flowpath.

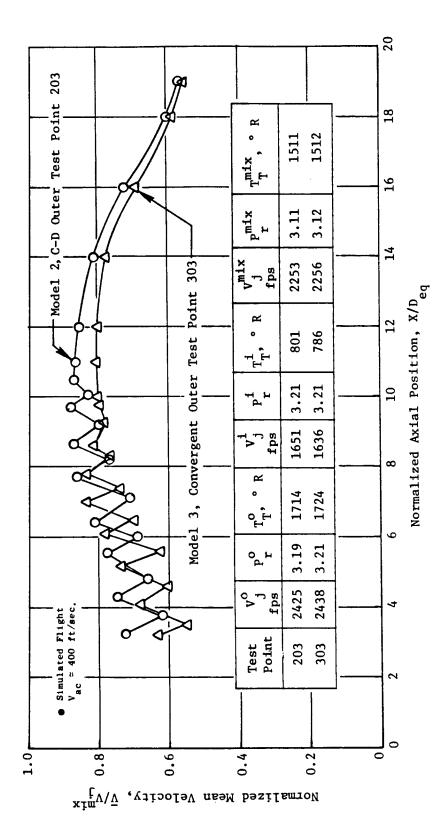
Table IV. Shock Noise Characteristics of Models 2 and 3.

- Data scaled to 1400 in.² flow area and 2400 ft sideline distance.
- $\bullet \quad \theta_{\mathrm{T}} = 50^{\circ}.$

Model	Test Point	Outer Nozzle Termination	v_{ac} , ft/sec	SPL _D , dB	f _D , Hz
2	201	C-D	0	87.1	250/315
3	301	С	0	87.2	250
2	203	C-D	400	91.2	315
3	303	С	400	93.4	315



Influence of Outer Flowpath Termination on the Centerline Mean Velocity at the Design Condition for the Outer Stream (Static). Figure 83.



Influence of Outer Flowpath Termination on the Centerline Mean Velocity at the Design Condition for the Outer Stream (Flight, $v_{ac} \simeq 400 \; \text{Ft/Sec})$. Figure 84.

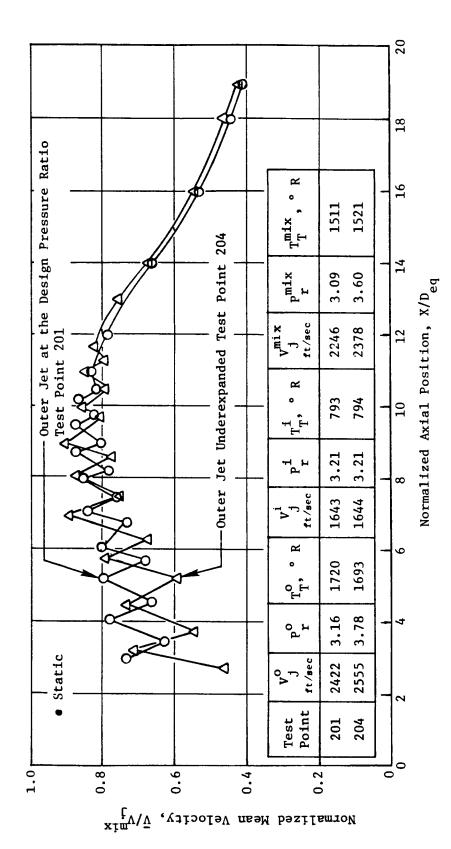
5.2.4 <u>Influence of Under/Overexpansion of Inner/Outer Streams</u> On the Cell Structure

AST/VCE Model 2 has been employed in order to study the influence of off-design pressure ratios on the flow characteristics. Figure 85 compares the centerline mean velocity axial variation when the outer stream is operated at the design pressure ratio of 3.2 and then at an underexpanded pressure ratio of 3.78, keeping the inner stream pressure ratio at 3.21 for both cases. Since the inner stream is supersonic, it is observed that a strong shock cell pattern occurs for both cases. Yet, because of the difference in the outer stream pressure ratios, the inner stream senses a different static pressure in both cases. The inner stream shock cell pattern begins with a compression wave for $P_T^O = 3.2$ whereas it begins with an expansion wave for $P_T^O = 3.78$. There are eight shock cells for the two cases, but the average shock cell spacing when $P_T^O = 3.16$ equals 0.9 D_{eq} and equals 1.08 D_{eq} when $P_T^O = 3.78$. This indicates that an underexpanded outer nozzle effectively lengthens the supersonic region of the inner jet.

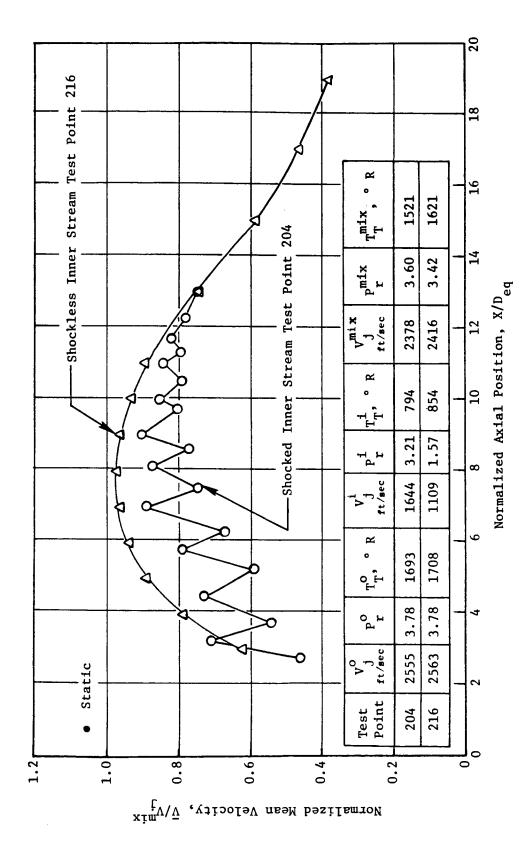
Figure 86 compares the centerline mean velocity variation when the inner stream is operated at subcritical and supercritical pressure ratios keeping the outer stream at an underexpanded pressure ratio (viz, $P_{\rm r}^{\rm O}=3.78$). There is an absence of shock pattern on the centerline when the inner stream is subsonic (i.e., $P_{\rm r}^{\rm i}=1.57$) and the outer stream is highly supersonic. And, when the inner stream is operated supersonically ($P_{\rm r}^{\rm i}=3.21$), a shock pattern exists consisting of eight shock cells with an average shock cell spacing of 1.08 $D_{\rm eq}$. For this reason, the inner stream pressure ratio is seen as a critical parameter in determining the occurrence of the shock cell structure of the jet plume and, consequently, the shock associated broadband noise.

To verify the above statement, the PNL directivities for the two test points are compared in Figure 87. Note that the front quadrant noise, which is dominated by shock noise (if present), for the case when $P_r^1=3.21$ is higher compared to the case when $P_r^1=1.57$. Also, the peak noise levels for both cases are seen to be identical, since the specific thrusts (i.e., V_{mix}) for both cases are about the same. At other aft angles, the PNL's are higher for $P_r^1=3.21$. The spectral distribution is analyzed next. Figures 88 through 90 show the spectral content at $\theta_1=60^\circ$, 130° (peak angle), and 140° , respectively. Notice the effective broadband shock noise suppression obtained at $\theta_1=60^\circ$ by having a subcritical inner stream pressure ratio. At the peak angle (i.e., $\theta_1=130^\circ$), the PNL's agree and so do the spectra as seen in Figure 89. The tones observed at $\theta_1=60^\circ$ and 140° have been analyzed using narrowband data with a band width of 10 Hz and are attributed to shock screech and reflections off the exhaust stack of the facility. Acoustic data after removing these tones still showed about a 4-5 PNdB reduction in the front quadrant noise level by employing the inner stream shock control.

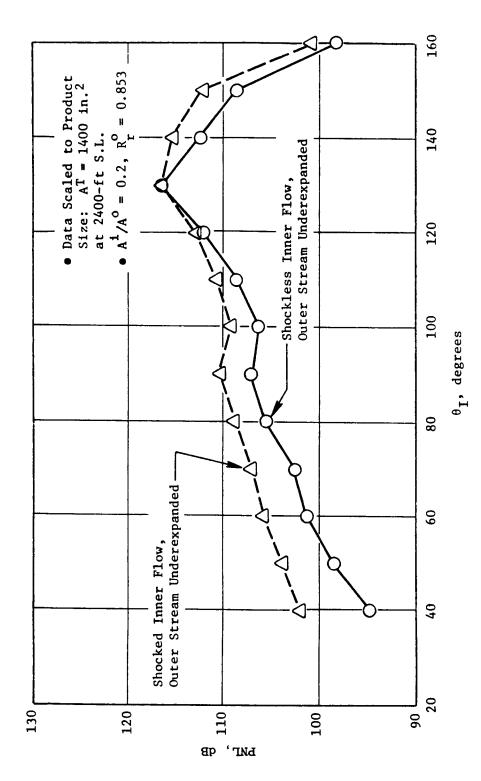
To further study the influence of the outer stream on the centerline mean velocity, variation of the centerline mean velocity is compared in Figure 91 for three outer stream conditions (viz., outer stream at design P_r^0 , under-



Effect of Underexpansion of the Outer Flow on the Mean Velocity Distribution Along the Centerline of Model 2 for Supersonic Inner Jet. Figure 85.



Effect of Inner Stream Pressure Ratio on the Degradation of Shock Cell Structure for Model 2. Figure 86.



PNL Directivity Showing Front Quadrant Shock Noise Suppression By Inner Stream Control. Figure 87.

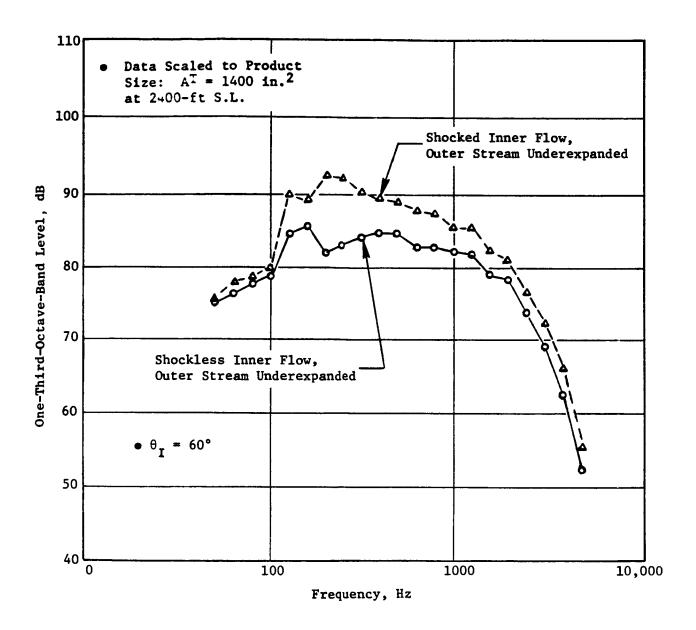


Figure 88. Front Quadrant (θ_{I} = 60°) Spectra Showing the Broadband Shock Noise Suppression By Holding the Inner Stream Pressure Ratio Subsonic.

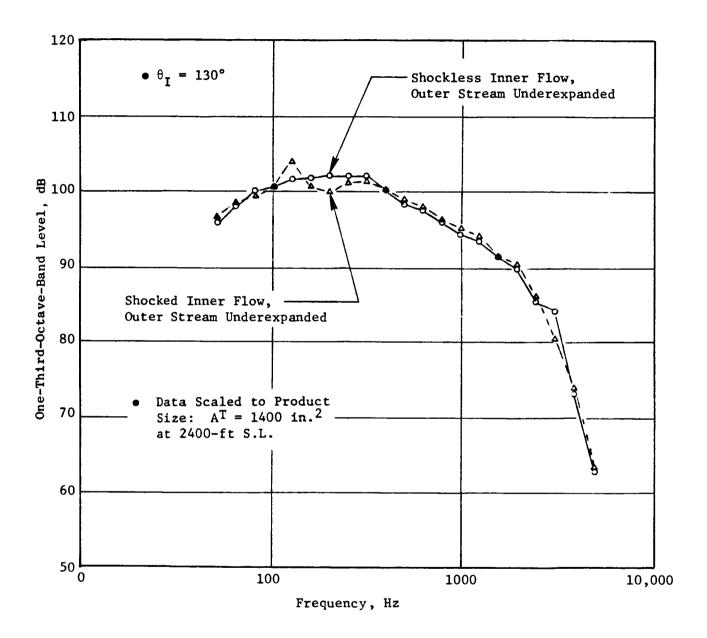


Figure 89. Peak Angle Spectra Showing Negligible Influence of the Variations in the Inner Stream Pressure Ratio.

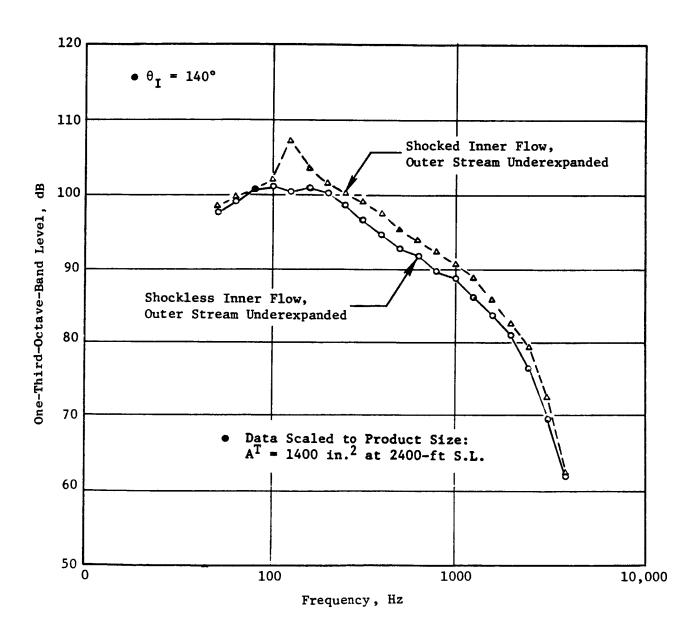
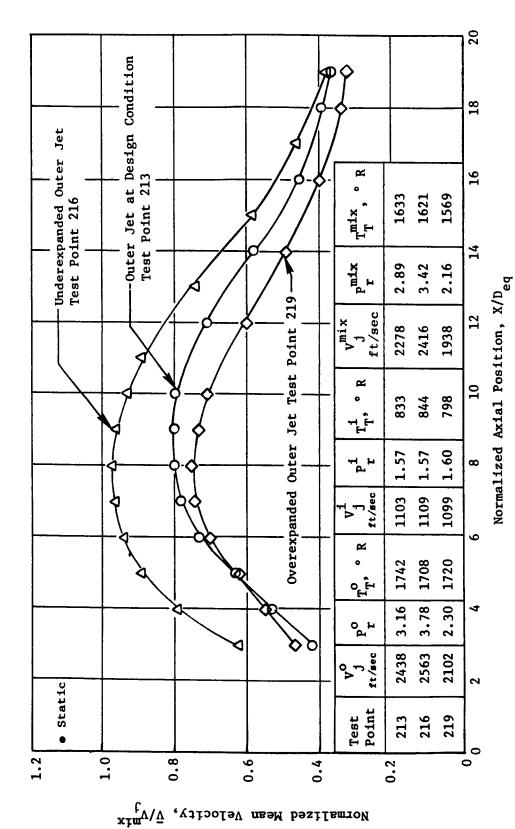


Figure 90. Aft Angle ($\theta_{\rm I}$ = 140°) Spectra for Shocked and Shockless Inner Flow.



Influence of Operating the Outer Stream at the Design and Off-Design Pressure Ratios on the Centerline Mean Velocity; Inner Stream Subsonic. Figure 91.

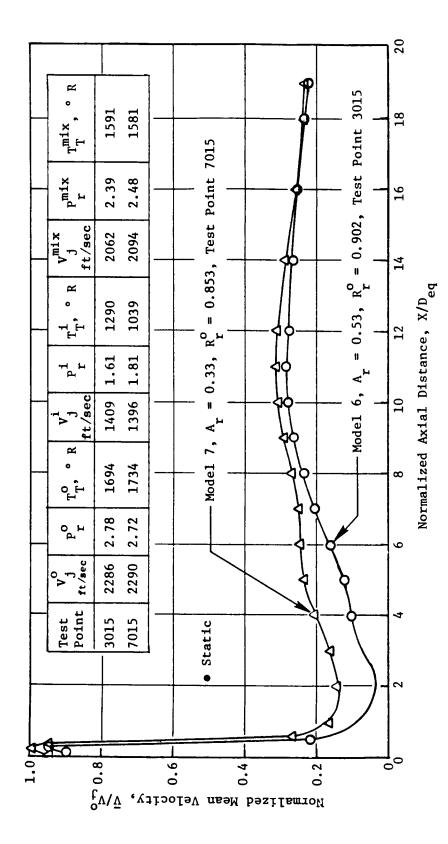
expanded outer stream, and overexpanded outer stream), keeping the inner stream subsonic. On the normalized basis also (i.e., $\overline{V}(V_j^{mix})$, the centerline velocity is maximum when the outer jet is underexpanded and is followed by a fully expanded case and then by the overexpanded case. This is due to the higher momentum transfer possible when the outer jet is underexpanded ($V_j = 2563 \text{ ft/sec}$) and is followed by the fully expanded case ($V_j = 2438 \text{ ft/sec}$) and then by the overexpanded case ($V_j = 2100 \text{ ft/sec}$).

5.2.5 Effect of Area Ratio and Outer Stream Radius Ratio on the Exhaust Plume Development

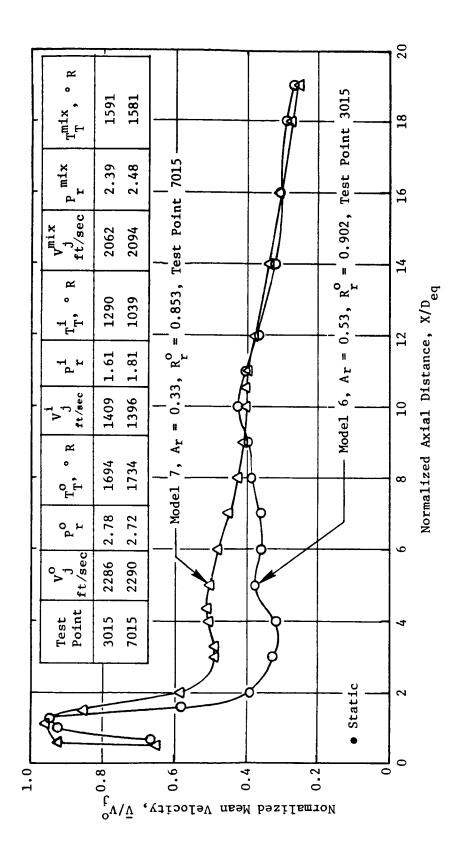
Nozzle area ratio and outer stream radius ratio are two important geometric parameters that determine the plume growth of coannular plug nozzles. AST/VCE Model 6 is a coannular plug nozzle with an area ratio (A^{i}/A^{o}) of 0.53 and an outer stream radius ratio (R_{r}^{o}) of 0.902, whereas AST/VCE Model 7 has an area ratio of 0.33 and an outer stream radius ratio of 0.853. As a result, a comparison of the plume characteristics of Models 6 and 7 will show the combined effect of A^{i}/A^{o} and R_{r}^{o} on plume growth.

Figure 92 shows the axial variation of the mean velocity at the radial location corresponding to the midpoint of the outer stream. The aerodynamic conditions of the outer and inner streams for both models are well matched and the mass-averaged conditions are obviously different due to the differences in the area ratio. It can be seen in Figure 92 that the traverse begins at $X/D_{eq} = 0$, since this traverse is at the midpoint of the outer stream. As soon as the inner stream appears (at $X/D_{eq} = 0.6$), the mean velocity for Model 6 drops below that of Model 7. As X/D_{eq} increases, the dual-flow character becomes prominent, and the differences between the two models increase. However, for X > 10 D_{eq} , the two traverses approach one another, indicating the diminishing effect of individual stream geometric parameters. Figure 93 shows the mean velocity axial variation at the radial location corresponding to the midpoint of the inner stream, which emerges at X = 0.6 D_{eq} . As seen in Figure 92, mean velocity for Model 6 drops in comparison to Model 7 for X between D_{eq} and 10 D_{eq} . For X > 10 D_{eq} , the two models have identical mean velocity decays. These trends are explained below.

Model 6 has a higher area ratio and a higher outer stream radius ratio compared to Model 7. As R_T^0 increases, the outer jet grows thinner and will suffer greater shear on either side due to the increased velocity gradient. Consequently, as R_T^0 increases, the outer jet will decay faster and will radiate less noise. However, for a given outer stream flow area, increasing R_T^0 would demand a larger engine diameter and subsequent weight and drag penalties. For this reason, a tradeoff between the potential noise benefit and weight and drag penalties has to be struck. Next, for given inner and outer stream flow conditions, as A^1/A^0 increases, the amount of inner flow available to slow down the outer flow increases, causing the outer flow to decelerate faster and to radiate less noise. However, as the amount of inner flow increases, the mass-averaged velocity (which is also the specific thrust defined



Influence of Area Ratio and Outer Stream Radius Ratio on the Mean Velocity Distribution at the Midpoint of the Outer Stream. Figure 92.



Influence of Area Ratio and Outer Stream Radius Ratio on the Mean Velocity Distribution at the Midpoint of the Inner Stream. Figure 93.

as thrust of the system per unit mass flow rate) decreases and hence the noise benefit has to be evaluated on an equal specific thrust basis.

5.2.6 Effect of Velocity Ratio on Plume Growth

Velocity ratio (V^{1}/V^{0}) is an important parameter which determines the velocity gradient between the inner and outer streams and, hence, determines the amount of shear exerted by the inner stream on the outer stream.

LV measurements were taken on Model 2 at two velocity ratios ($v^i/v^o = 0.52$ and 0.60), keeping the outer stream conditions the same. Figures 94 through 96 show the radial profiles at three X/D_{eq} stations. At $X/D_{eq} = 1.0$ (Figure 94), the inner stream has just appeared, and the inner streams at the two velocity ratios can be viewed distinctly; the outer streams do not show any effect due to the differences in the inner stream velocities. At $X/D_{eq} = 1.81$ (Figure 95), the plume has come closer to the jet centerline since the plug radius has been reduced. The peak outer velocities have remained at the same levels as at $X/D_{eq} = 1.0$, and again the inner streams have no influence on the outer streams. Instead, the inner streams have decelerated compared to their corresponding values at $X/D_{eq} = 1.0$. At $X/D_{eq} = 2.54$ (Figure 96), the absence of the plug is observed. The inner velocities have further reduced. Though the peak values of the outer velocities have not changed, the outer stream has grown wider indicating the momentum loss of the outer stream to the ambient.

The observations noted above can be explained by utilizing the principles of momentum transfer. It has been noted that the inner stream does not seem to have a noticeable influence on the outer stream, and the variation of velocity ratio does not alter the basic plume characteristics of the outer stream. Shear stress exerted by the inner stream on the outer stream is a dynamic quantity and depends on the ratio of the momentums of the two streams. Velocity ratio is a kinematic quantity which only determines the velocity gradient, not the shear stress.

One can define the momentum ratio as

Momentum Ratio =
$$\frac{w^i}{w^o} \frac{v^i}{v^o}$$

The pertinent aero conditions for the two test points considered are listed as follows:

Test Point	V ^O ft/sec	P° r	TO T	Vi ft/sec	P _r	T _T	V ^{mix} ft/sec	P ^{mix} r	T ^{mix} T
210	2085	2.29	1698	1258	2.07	701	1907	2.21	1484
219	2102	2.30	1720	1099	1.60	798	1938	2.16	1569

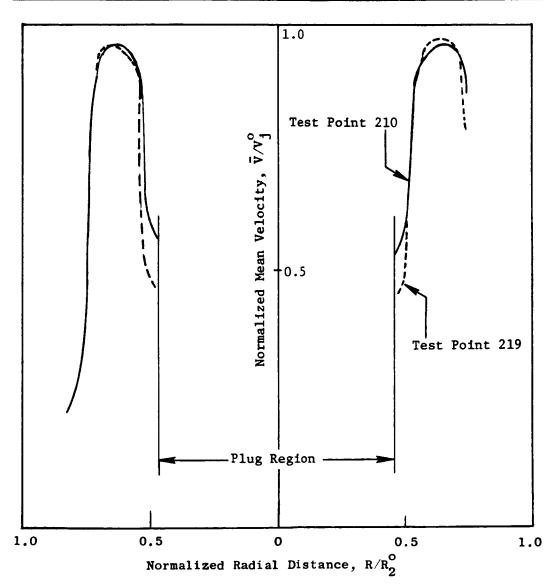


Figure 94. Radial Profile of the Mean Velocity at $\rm X/D_{eq} = 1.0$ Showing the Appearance of the Inner Stream at Two Velocities.

Test Point	v ^o j ft/sec	P ^o r	To T	v ⁱ j ft/sec	P ⁱ r	T _T	v ^{mix} j ft/sec	P ^{mix}	TT R
210	2085	2.29	1698	1258	2.07	701	1907	2.21	1484
219	2102	2.30	1720	1099	1.60	798	1938	2.16	1569

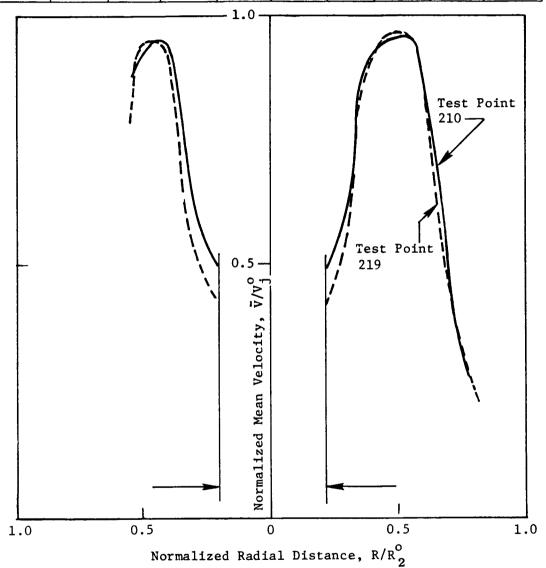
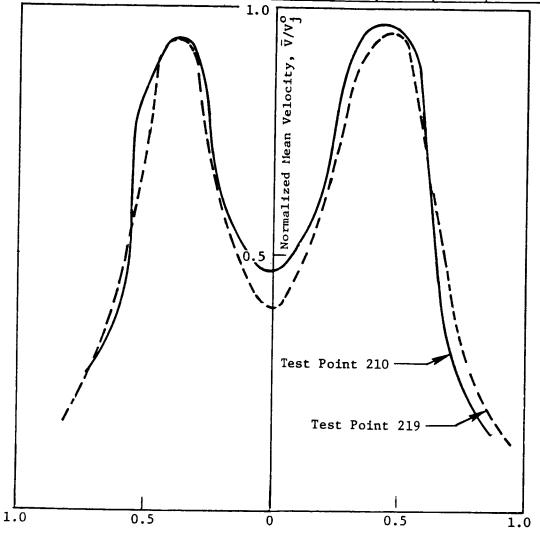


Figure 95. Radial Profile of the Mean Velocity at $\rm X/D_{eq}$ = 1.81 Showing the Deceleration of the Inner Streams Without Affecting the Outer Streams.

Test Point	v ^o j ft/sec	P°	T ^O T	v ⁱ j ft/sec	P _r i	T _T	v ^{mix} j ft/sec	P ^{mix}	TT R
210	2085	2.29	1698	1258	2.07	701	1907	2.21	1484
219	2102	2.30	1720	1099	1.60	798	1938	2.16	1569



Normalized Radial Distance, R/R_2^0

Figure 96. Radial Profile of the Mean Velocity at $\rm X/D_{eq}$ = 2.54 Showing the Disappearance of the Plug and Negligible Influence on the Outer Streams.

Test Point	v ^o fps	w ^o pps	v ⁱ fps	w ⁱ pps	v^i/v^o	$\frac{\overset{\text{w}^{i}}{\circ} \overset{\text{v}^{i}}{\circ}}{\overset{\text{o}}{\vee} \overset{\text{o}}{\circ}}$	V ^{mix} j ft/sec
210	2085	494	1258	135	0.603	0.165	1907
219	2102	492	1099	96	0.523	0.102	1938

Thus, the momentum of the inner stream available to retard the outer stream is small for both test points, but this fact is not evident when one looks at the velocity ratio only. In order to have the inner stream exert noticeable influence on the outer stream, the momentum ratio has to be increased without increasing the velocity ratio. This means the inner-to-outer stream area ratio has to be increased and/or the inner-to-outer stream density ratio has to be increased. When the inner-to-outer stream momentum ratio is increased, the mass-averaged velocity reduces. Therefore, the tradeoff between increased shear benefit and lowered specific thrust, due to the increase in the momentum ratio, has to be further evaluated.

5.2.7 Influence of Struts on Mean and Turbulent Velocities

AST/VCE Models 1A and 2 are identical coannular plug nozzles except that Model 1 has eight struts in the outer flowpath while Model 2 has none. So, a comparison of the LV measurements of the mean and turbulent velocities just downstream of the outer nozzle exit plane (i.e., before the inner flow emerges) would indicate the influence of the struts, if any. Figure 97 indicates the distribution of the mean and turbulent velocities at $X/D_{eq} = 0.29$ on either side of a strut location for Models 1A and 2. Because there are no significant differences in the mean velocity distribution and in addition, the levels of turbulent velocities are the same for both models, it can be concluded that the struts do not significantly alter the mean and turbulent velocity characteristics of the nozzle.

5.2.8 A Rationale for the Observed Flow Asymmetry

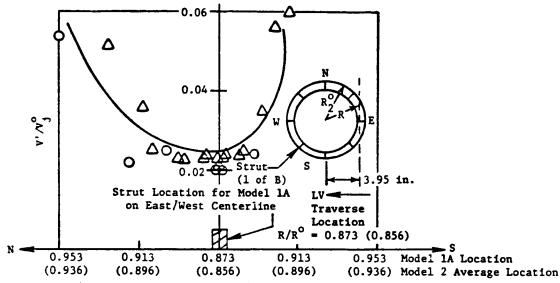
During the course of the LV testing, certain radial profiles of coannular plug nozzles showed asymmetric velocity distribution about the jet centerline. Figures 98 and 99 show the LV radial traverses at $X/D_{eq}=7.5$ for AST/VCE Models 1A and 2, respectively. The asymmetry in peak velocity is 32% for Model 1 and 28% for Model 2. The measurement of the annular gap showed that the higher velocity occurred on the side with the larger annular height. Thus, the velocity profile asymmetry can be traced to the geometric misalignment of the nozzle hardware which results in unequal circumferential annular gap distribution. Since the outer nozzle in the two cases is operating at a supercritical condition, the velocity on the side of larger annular height is expected to be higher from a continuity consideration. An exactly opposite trend should be expected for subcritical exit conditions.

Symbol	Model	Struts	Test Point	v ^o j ft/sec	P° r	TO T	v ⁱ j ft/sec	P _r i	T _T	vmix j ft/sec	P ^{mix}	TT R
0	1A	Yes	116	2177	2.47	1715	1332	2.14	755	2004	2.37	1518
Δ	2	No	222	2170	2.45	1718	1341	2.17	756	1999	2.35	1520

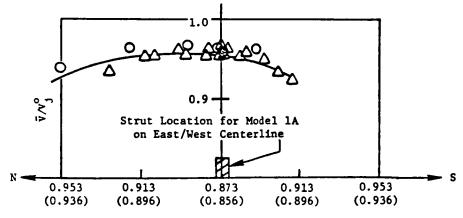
• Static

$$D_{eq} = 5.238, R_2^0 = 4.593$$

(a) Turbulence Intensity Profiles from Histogram Data

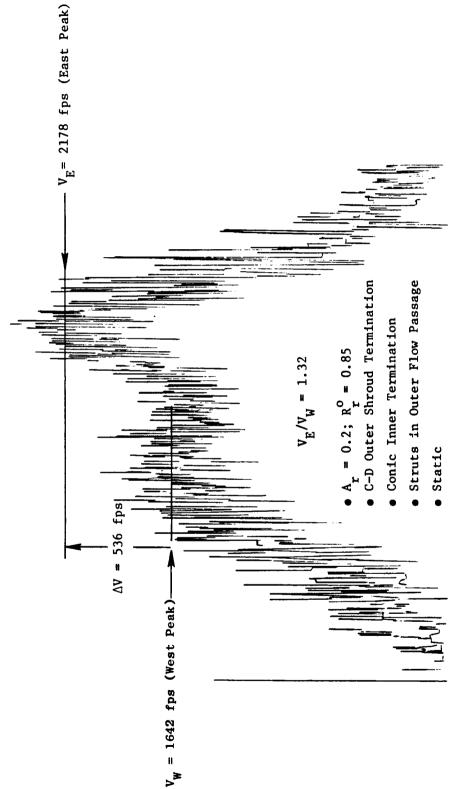


(b) Mean Velocity Profiles from Histogram Data



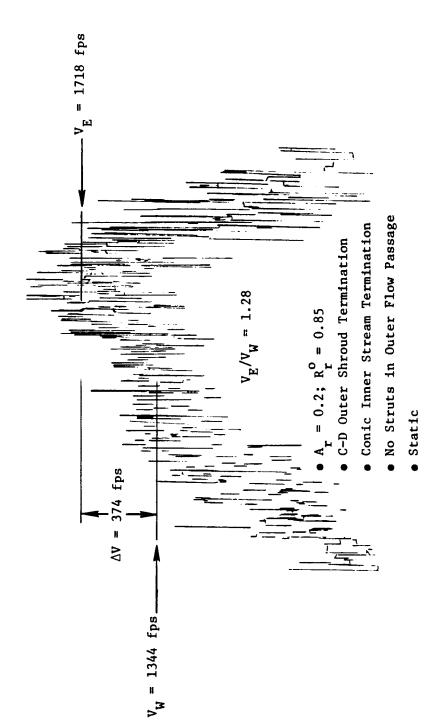
North/South Chordwise Radial Location, R/R_2^0

Figure 97. Strut Effect on Downstream Velocity Flow Profiles at $X/D_{eq} = 0.294$.



T _T ° R	1481
p ^{mfx}	3.15
$v_{\rm j}^{\rm mix}$ ft/sec	2239
${ m T}_{ m T}^{ m t}$ ° R	682
$rac{ ext{p}^{ ext{1}}}{ ext{r}}$	3.19
$v_{\rm j}^{\rm i}$ ft/sec	1636
${ m T}_{ m T}^{ m o}$	1678
P ^o r	3.23
$v_{\rm j}^{ m o}$ ft/sec	2411
Test Point	101

Radial Traverse Showing Mean Velocity Flow Asymmetry for Model 1A for X/D $\,\simeq\,$ 7.5 with Annular Gap Not Centered Before Figure 98.



T ^{mfx} ° R	1520
p ^{mix} r	2.35
vmix j ft/sec	1999
${ m T}_{ m T}$	156
$rac{\mathbf{p_{t}^{t}}}{\mathbf{r}}$	2.17
$v_{\rm j}^{\rm i},$ ft/sec	1341
T ^o ° R	1718
P ^O r	2.45
v ^o j ft/sec	2170
Test Point	222

Radial Traverse Showing Mean Velocity Flow Asymmetry for Model 2 for X/D $\,\approx\,7.5$ with Annular Gap Not Centered Before Test. Figure 99.

In order to minimize the velocity asymmetry, the above tests were rerun with the annular gap measured and monitored to within $\pm 3\%$ variation circumferentially. Figures 100 and 101 show the rerun LV radial traverses at $X/D_{eq}=7.5$ for Models 1A and 2, respectively. Observe that for Model 1A, the velocity asymmetry now reduced from 32% to 21%. Similarly, for Model 2 the velocity asymmetry was reduced from 28% to 3%. The presence of struts in Model 1A and their absence in Model 2 could be the cause for the observed differences in the measured velocity asymmetry even after the gap was monitored. The annular gaps were measured and adjusted to within $\pm 3\%$ when the models were cold. It is hypothesized that an unequal thermal expansion of the struts could have caused Model 1A to be excessively misaligned during the rerun than Model 2 and thus account for the different reductions in the velocity asymmetry that were observed even after the nozzles were carefully aligned.

5.2.9 Summary of Observations

The extensive deployment of the laser velocimeter to measure the mean and turbulent velocities of the jet plumes of scale model conic and coannular plug nozzles has yielded valuable information regarding the mixing characteristics, which in turn has enabled one to understand the noise characteristics of these nozzles.

The principal conclusions of this study are summarized below:

- The measured conic nozzle shock characteristics at typical takeoff condition agree well with those predicted by the empirical model of Fisher and Harper-Bourne. Due to reduction in shear, the free jet stretched the shock cell pattern and reduced the decay rate of the mean velocity, as well as the levels of the turbulent velocities. Thus it is not known whether a free jet reduces shock noise or enhances it.
- The radial profile measurement at various axial stations for a coannular plug model showed that the dual-flow character is maintained for about 7 to 8 Deq lengths downstream of the nozzle exit. The turbulent velocity distribution for a coannular plug nozzle shows similar characteristics to those of a conic nozzle.
- 3. Although a C-D flowpath for the outer stream did not eliminate shock cell pattern, it resulted in a small reduction in the shock noise when compared to a convergent flowpath for the outer stream.
- 4. A study of the under-/overexpanded nozzles showed that an underexpanded outer nozzle effectively lengthens the supersonic region of the inner jet. Even so, the outer stream pressure ratio does not significantly alter the basic aerodynamic characteristics of the inner jet which are mainly functions of the inner stream pressure ratio. The inner stream pressure ratio is a critical parameter in determining the shock cell distribution, as evidenced by the LV and

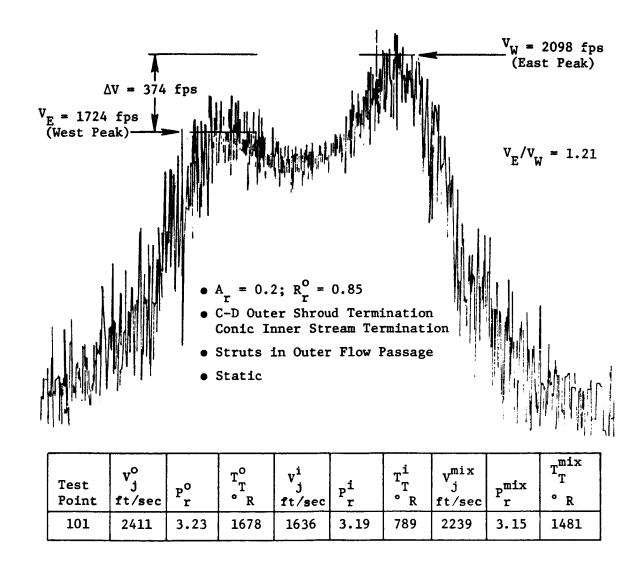
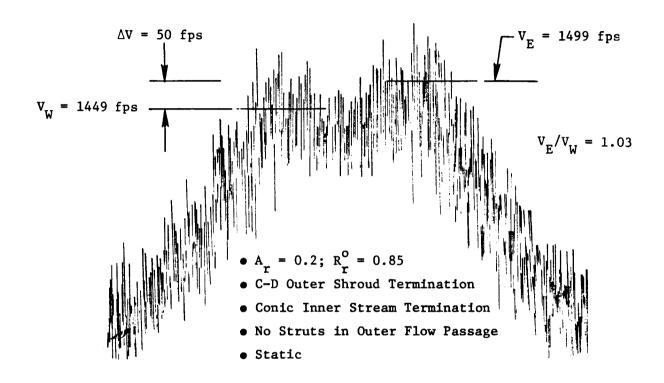


Figure 100. Radial Traverse Showing Mean Velocity Flow Asymmetry for Model 1A for $\rm X/D_{eq} \approx 7.5$ with Annular Gap Centered Before Test.



Test Point	v ^o j ft/sec	P°r	TO T R	v ⁱ j ft/sec	P _r i	T _T	V ^{mix} j ft/sec	P ^{mix}	Tmix T
222	2170	2.45	1718	1341	2.17	756	1999	2.31	1520

Figure 101. E-W Radial Traverse Showing Mean Velocity Flow Asymmetry for Model 2 for X/D $_{\rm eq}$ $\stackrel{\simeq}{=}$ 7.5 with Annular Gap Centered Before Test.

acoustic measurements. If the inner stream is operated at subcritical pressure ratios, considerable amount of shock noise can be eliminated with no loss of specific thrust.

The tones in the data are due to shock screech and reflections off of the exhaust stack of the facility. Acoustic data without the tones still show about a 4 to 5 PNdB reduction in the front quadrant noise level by employing a subcritical inner stream.

- The combined influence of the area ratio (A^i/A^o) and the outer stream radius ratio (R_r^o) was studied using Models 6 and 7. With an increase in A^i/A^o and R_r^o , the outer jet is sheared more strongly which results in its faster decay. However, an increase in R_r^o and A^i/A^o would require a larger engine size and would reduce the specific thrust. Thus, an optimum value of R_r^o and A^i/A^o has to be found.
- 6. For the nozzle considered (Model 2, $A^{1}/A^{0} = 0.2$, $R_{T}^{0} = 0.853$), the inner stream does not have a noticeable influence on the outer stream, and a variation of the velocity ratio does not alter the basic plume characteristics of the outer stream. In order to have the inner stream exert a noticeable influence on the outer stream, the momentum ratio has to be increased without increasing the velocity ratio. There exists a tradeoff between the increased shear benefit and the lowered specific thrust due to the increase in the momentum ratio, which needs to be evaluated by further studies.
- 7. The struts in the outer flowpath of the coannular plug nozzle did not significantly alter distribution of the mean and turbulent velocities.
- 8. Due to a geometric misalignment, flow asymmetry was noted in some coannular plug nozzles. This was minimized in the nozzle having no struts by an accurate pretest alignment. For the case of the nozzle with struts, a careful alignment did not eliminate the asymmetry but did reduce it.

5.3 A UNIQUE COANNULAR PLUG NOZZLE JET NOISE PREDICTION PROCEDURE

5.3.1 The Basic Concept of the Procedure

An effort has been made to develop a semiempirical spectral prediction method for coannular plug nozzles operated in the inverted velocity profile mode that will account for the various noise-generating mechanisms. The effects of flight on coannular jet mixing and shock noise are also predicted. The guidelines for such a prediction method are that it be based on the physics of the flow and on its noise-radiating characteristics, and yet remain simple. The M*G*B model* (Reference 20) developed under a DOT contract incorporated the source spectrum, convective amplification, and fluid

shrouding effects in order to predict the noise from eddies which are then integrated to obtain the jet mixing noise spectra for any jet/jets issuing from any given nozzle configuration. However, the M*G*B model calculates the aerodynamic properties at each slice of the jet before it can predict the radiated noise from that slice. This requires a verified aerodynamic model before verifying the acoustic model. As very little detailed aerodynamic data is available for coannular nozzles, it was considered useful to develop a semi-empirical model using the large acoustic data bank (References 2 and 9) but still use the same physical concepts incorporated in M*G*B.

5.3.2 An Outline of the Prediction Procedure

The prediction procedure consists of two modules:

- Coannular jet mixing noise prediction
- Coannular shock noise prediction.

Extensive details of these modules, the computer program, user's manual, sample input/output cases and comprehensive comparisons with the data can be found in Reference 21. Hence, only a brief description of the method is given here.

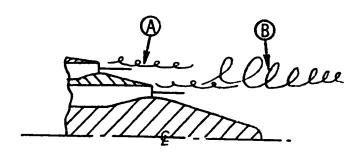
A. Coannular Jet Mixing Noise

The method developed to predict the jet mixing noise spectrum identified the noise spectrum as being made up of:

- Source spectrum due to small-scale, random turbulence eddy fluctuations
- 2. Convective amplification and Doppler shift due to convecting eddies
- 3. Fluid shielding or flow shrouding of the eddies by the mean flow.

1. Source Spectrum

The first step in the prediction procedure was to determine the source spectrum. The spectrum at 90° is the source spectrum, as there are no convection effects or fluid shielding effects. The parameters needed to define this spectrum were the characteristic velocity and length scales at this angle of emission. From coannular plug nozzle acoustic data (Reference 2), the 90° spectrum (made "lossless," i.e., air attenuation added to the measured data) showed two distinct regions (Figure 102) which were appropriately defined as



Coannular Plug Nozzle Configuration

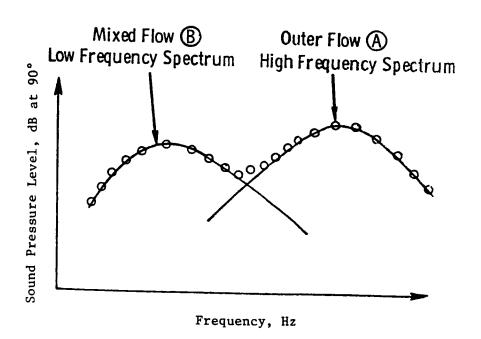


Figure 102. Source Spectrum Modeling of High and Low Frequency Regions of a Coannular Plug Nozzle.

the low and high frequency regions of the spectrum. The high frequency portion of the spectrum was identified as being generated by the outer flow before it merged with the inner flow, while the low frequency portion of the spectrum was identified as being generated by the mixed flow.

Having identified the outer jet as the probable source of the high frequency portion of the source spectrum, the fully expanded outer jet velocity and the hydraulic diameter as defined below were chosen as the characteristic velocity and length scales for this portion of the source spectrum (Figure 103).

$$D_{hyd}^{o} = \frac{4 \times \text{outer jet noise radiating area}}{\text{outer jet noise radiating perimeter}}$$

$$= 2 h^{o} \left(1 + R_{r}^{o}\right) \tag{1}$$

Similarly, for the low frequency portion of the source spectrum, the massaveraged velocity and the diameter based on the total flow area as defined below were chosen as the characteristic velocity and length scales, respectively.

$$v_{j}^{\text{mix}} = \frac{v_{j}^{\text{o}} w + v_{j}^{\text{i}} w^{\text{i}}}{w^{\text{o}} + w^{\text{i}}}$$
(2)

$$D_{eq}^{T} = \frac{4}{\pi} (A^{o} + A^{i})^{1/2}$$
 (3)

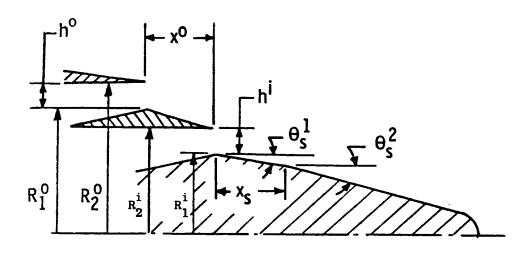
Having chosen the characteristic velocity and length scales, the large acoustic data base for coannular jets was used once again to determine the normalized source spectrum (normalized sound pressure level as a function of the Strouhal number).

For the low frequency noise portion of the source spectrum, the peak Strouhal number was observed to be correlated by:

$$\left[\frac{f_p^{LF} \quad D_{eq}^T}{v_i^{mix}}\right] \left[\frac{T_T^{mix}}{T_a}\right]_{eff} = 0.9$$
(4)

where

$$\left[\frac{T_{T}^{\text{mix}}}{T_{a}}\right]_{\text{eff}} = 0.65 \frac{T_{T}^{\text{mix}}}{T_{a}} + 0.35$$
(5)



 $R_r = Radius Ratio (R_1/R_2)$

h = Step Height

A = Area

D_{eq} = Equivalent Circular Diameter based on A

 Θ_{s} = Ramp Angle of Inner Plug

R = Radius

x = Distance

Superscripts

o = Outer Flow Region

i = Inner Flow Region

Figure 103. Schematic of Nozzle Configuration and Definition of Parameters.

and

$$T_{T}^{mix} = \frac{T_{T}^{o} \quad w^{o}_{+} T_{T}^{i} \quad w^{i}}{w^{o}_{+} \quad w^{i}}$$
 (6)

Equation 5 assumes that the total temperature profiles are similar, while Equation 4 was observed to predict the same peak Strouhal number for conic nozzles as did the SAE method (Reference 22).

However, for the high frequency portion of the source spectrum, the peak Strouhal number was found to correlate by

$$\left[\frac{f_{p}^{HF}}{v_{j}^{o}}\right]\left[\frac{T_{T}^{o}}{T_{a}}\right]_{eff} = 1.18$$
(7)

where

$$\left[\frac{T_{T}^{o}}{T_{a}}\right]_{eff} = 0.65 \frac{T_{T}^{o}}{T_{a}} + 0.35$$
 (8)

Normalization of the SPL's for the low frequency portion of the source spectrum was found to be best using the conventional Lighthill and Hoch velocity and density dependence laws (References 23 and 24),

i.e.,
$$SPLN^{LF}(f) = SPL^{LF}(f) - c_1 \log_{10} \left(v_j^{mix} / c_a \right)$$

$$- 10 \log_{10} \left(\rho_j^{mix} / \rho_a \right)^{\omega} - 10 \log_{10} \frac{(A^T)}{\left(\frac{A^T}{R^2} \right)}$$
(9)

where

$$c_{1} = \begin{cases} 75, & \text{for } V_{j}^{\text{mix}}/C_{a} \leq 2.0 \\ \\ 80, & \text{for } V_{j}^{\text{mix}}/C_{a} > 2.0 \end{cases}$$

The shape of the normalized low frequency portion of the spectrum was based on that of a conic nozzle, while the absolute level was based on coannular plug nozzle jet mixing noise data. Figure 104 shows the normalized low frequency source spectrum for coannular jet mixing noise.

Similarly, for the high frequency spectrum, normalized SPL's were determined in addition to the velocity, density, and area terms, the high frequency portion of the source spectrum was influenced by velocity ratio, radius ratio, and area ratio that are incorporated into the normalizing factor. These influences were empirically derived from the existing data base. The best correlation was obtained by

$$SPLN^{HF}(f) = SPL^{HF}(f) - 80 \log_{10} (v_j^o/C_a)$$

$$-10 \log_{10} (\rho_j/\rho_a)^\omega - 10 \log_{10} (A^o/R^2)$$

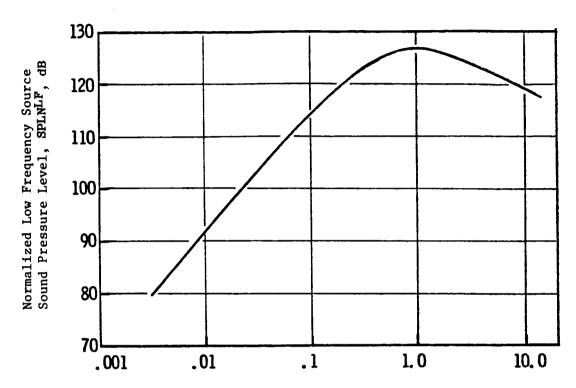
$$+ 50 \log_{10} (R_r^o) - 10 \log_{10} (1+A_r^i)$$

$$-15 \log_{10} (4.42 v_r^2 - 4.56 v_r + 2.15)$$
(10)

where

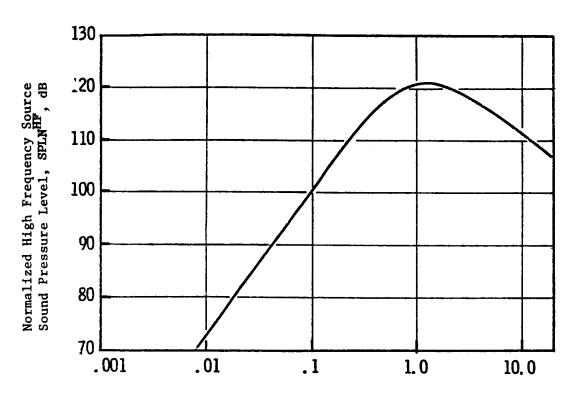
$$v_r = v_j^i/v_j^o$$
 and $A_r = A^i/A^o$

Figure 105 shows the normalized high frequency portion of the source spectrum. Using Figures 104 and 105, the low and high frequency portions of the source spectrum are determined, and the frequency (1/3-octave band) at which they intersect determines the end of the low frequency and the beginning of the high frequency portion of the source spectrum for any set of inner and outer jet flow conditions.



Effective Low Frequency Stroubal Number, $S^{LF} = ((f \cdot D_{eq}^{T})/(V_{j}^{mix} - V_{ac})) \cdot (T_{T}^{mix}/T_{a}) \text{ eff}$

Figure 104. Normalized Low Frequency Spectrum at $\theta_{\rm I}$ = 90°.



Effective High Frequency Strouhal Number,

$$S^{\mathrm{HF}} = ((f \cdot D_{\mathrm{hyd}}^{\mathrm{o}})/(V_{\mathrm{j}}^{\mathrm{o}} - V_{\mathrm{ac}})) \cdot (T_{\mathrm{T}}^{\mathrm{o}}/T_{\mathrm{a}})_{\mathrm{eff}}$$

Figure 105. Normalized High Frequency Spectrum at $\theta_{\rm I}$ = 90°.

2. Convective Amplification and Doppler Shift Due to Convecting Turbulence Eddies

As the turbulence eddies (that generate the mixing noise) are convected relative to the observer, there is an amplification of the source spectrum in the aft angles as well as an attenuation in the forward quadrant. The amplification/attenuation is a function of the eddy convection Mach number and the observer angle relative to the inlet. The convection Mach number of the turbulence eddies is given by

$$M_c = 0.5 \left[0.55 + \frac{0.39}{V_r}\right] \left(V_j/a_o\right) \text{ for } V_r < 1.0$$
 (11)

= 0.55
$$v_j/a_0$$
 for $v_r \ge 1$

The above expression was derived from the M*G*B (Reference 10) prediction method as applied to inverted velocity profile coannular jets, though the constants are different because here the jet exit mean velocity is used in place of the local mean velocity.

The expression for the convective amplification is

$$\Delta SPL_{CA} = N \left(10 \log_{10} C\right)$$
 (12)

where

$$c = \left[\left(1 + M_c \cos \theta_I \right)^2 + \alpha^2 M_c^2 \right]^{1/2}$$
 (12a)

and $\alpha = 0.325$ (from Reference 25).

Balsa's theory (References 26 through 28) for a conic jet suggests that the constant N, which is a function of $\theta_{\rm I}$, is 3 for angles to the inlet less than the cutoff angle for shielding and N = 7 for angles greater than the cutoff value. Based on these limits, the variation of N with $\theta_{\rm I}$ was determined from the data base and is shown in Figure 106. The convection Mach numbers were calculated using $V_{\rm j}^{\rm mix}$ and $V_{\rm j}^{\rm o}$ for the low and high frequency portions of the spectrum, respectively.

The frequency shift associated with eddies moving relative to the observer is given by

$$f_{\theta_{T}} = \frac{f_{90}^{\circ}}{C} \tag{13}$$

This shifts the peak frequency to lower frequencies in the front quadrant and to higher frequencies in the aft quadrant.

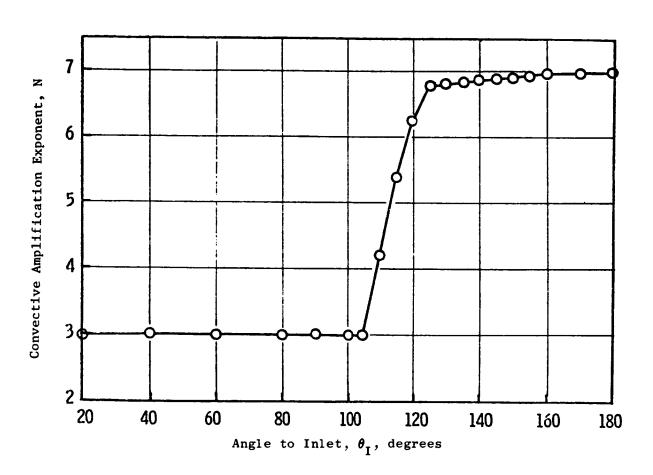


Figure 106. Variation of Convective Amplification Exponent with Angle.

3. Mean Flow Shrouding

Since Balsa's formulation (References 26 through 28) of the shielding function must be applied to each eddy in the flow depending on its location and the observer angle, a simplification is necessary for a semiempirical prediction procedure. Gliebe (Reference 29) derived such a simplification for conic nozzles by assuming average flow properties to replace the local flow properties, thereby determining the shielding function and shielding factor (defined later) as functions of observer angle, characteristic flow, geometric properties, and source frequency. Comparisons made by Gliebe with data based on these assumptions showed reasonable agreement. Based on the initial success of Gliebe, his expressions for shielding function and shielding factor were scrutinized, refined, and extended here to apply to coannular plug nozzles.

From Balsa's analysis, the reduction in noise level radiated by a single eddy (Figure 107) due to fluid shielding is given by

$$(\Delta dB)_{\text{shielding}} = -2K_1 g^2 (R)^{1/2} dR$$
 (14)

where the shielding function is given by

$$\frac{\left|\frac{g^{2}}{c}\right|^{1/2}}{c} = \sqrt{\frac{\left[\left(1+M_{c} \cos \theta_{I}\right)^{2} (c/c_{a})^{2}\right] - \cos^{2} \theta_{I}}{c^{2}}}$$
(15)

where

$$(c/c_a)^2 = 0.65 T_j/T_a + 0.35$$

Instead of defining the shielding function for each eddy, an average shielding function was defined where M_{C} and C are now based on the characteristic mean velocity of the flow rather than the local mean velocity. The constant of proportionality defined as the shielding factor is then given by

$$H\left(\frac{fD}{C_a}\right) = \frac{\Delta SPL \ (f)_{shielding}}{\left(\frac{fD}{C_a}\right) \ \left(2\pi C_a\right) \left(\frac{g}{C}\right)}$$
(16)

Using the data base on coannular jets (Reference 2), a curve of $H(fd/C_a)$ versus (fD/C_a) (Figure 107) was determined which could be used with the low and high frequency portions of the spectrum. From theory, it can be shown that shielding of the noise from the turbulence eddies by the mean flow occurs only for negative values of the shielding function g^2 . So, the angle θ_C at

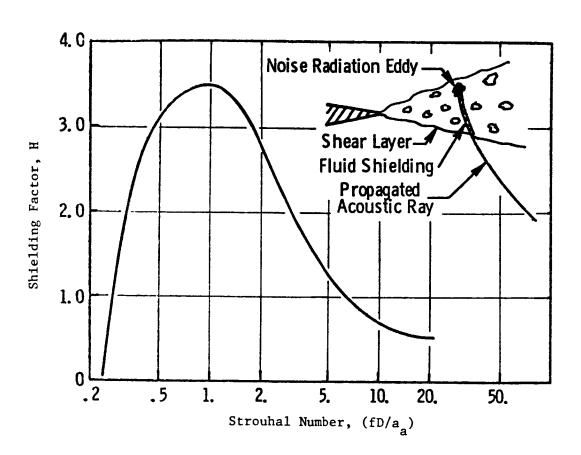


Figure 107. Variation of Shielding Factor with Strouhal Number for Coannular Plug Nozzles.

at which there will no longer be any shielding is obtained by equating the argument of Equation 15 to zero. This yields

$$\theta_{\rm C} = \cos^{-1} \left[\frac{-1}{({\rm C/C_a}) + {\rm M_C}} \right] \tag{17}$$

Using Equations 14 through 17, the reduction in SPL as a function of frequency due to fluid shielding can be obtained for the low and high frequency portions of the spectrum at each angle.

Therefore, the procedure for predicting the jet mixing noise spectrum at any angle for coannular plug nozzles can be summarized as follows: the source spectrum is predicted using Equations 1 through 10 and Figures 104 and 105. Next, convective amplification effects on both noise level and frequency are added using Equations 11 through 13 and Figure 106. Finally, the fluid shielding effects are determined using Equations 14 through 17 and Figure 107. Then, the predicted sound pressure level at any frequency and angle is given by

where

The spectrum thus predicted is lossless (i.e., no air attenuation effects are included). It can then be scaled, extrapolated, or simply converted to standard day conditions using the appropriate procedures.

4. Extension of the Jet Mixing Noise Prediction Method to Flight

In order to apply the prediction method described above to engines in flight, the method had to be extended to include the effects of the flight velocity on the noise signature.

Based on comparisons of conic free-jet data transformed to flight with static data, the Strouhal numbers corresponding to peaks in the low and high frequency components of the coannular jet mixing noise source spectrum are modified from those for static predictions as follows:

$$\left(\frac{f_{p}^{LF} D_{eq}^{T}}{v_{j}^{mix}}\right) \left(\frac{T_{T}^{mix}}{T_{a}}\right)_{eff} \left(\frac{v_{j}^{mix} - v_{ac}}{v_{j}^{mix}}\right) = 0.9$$
(19)

and

$$\left(\frac{f_p^{HF} p_{hyd}^o}{v_j^o}\right) \left(\frac{T_T^o}{T_a}\right)_{eff} \left(\frac{v_j^o - v_{ac}}{v_j^o}\right) = 1.18$$
(20)

where V_{ac} = aircraft velocity.

Also, the static source spectrum levels for the low and high frequency components of the spectrum are reduced, respectively, by

$$SPL^{LF} = 20 \log_{10} \frac{v_{j}^{mix} - v_{ac}}{v_{j}^{mix}}$$

$$SPL^{HF} = 20 \log_{10} \frac{v_{j}^{o} - v_{ac}}{v_{j}^{mix}}$$
(21)

Hence, using Equations 19 through 21, the in-flight source spectrum can be predicted.

To predict the flight spectra at any other angle, the same procedure outlined for static prediction is used with the convection Mach numbers now based on relative velocity, i.e.,

$$M_{c}^{LF} = \frac{1}{2} \left[0.55 + \frac{0.39}{V_{r}} \right] \left(\frac{V_{j}^{mix} - V_{ac}}{a_{a}} \right) \text{for } V_{r} < 1.0$$

$$= 0.55 \left(V_{j}^{mix} - V_{ac}/a_{a} \right) \text{ for } V_{r} \ge 1.0$$

$$M_{c}^{HF} = \frac{1}{2} \left[0.55 + \frac{0.39}{V_{r}} \right] \left(\frac{V_{j}^{0} - V_{ac}}{C_{a}} \right) \text{ for } V_{r} < 1.0$$

$$= 0.55 \left(V_{j}^{0} - V_{ac} \right) / C_{a} \text{ for } V_{r} \ge 1.0$$

It must be noted that, except at $\theta_{\rm I}=90^{\circ}$, this method of predicting flight noise spectra does not use the conventional method of applying flight-effect corrections to the static spectra in the form of a relative velocity exponent as used by several existing methods (References 30 and 31). As a result, the controversy of deciding the correct value of the exponent is avoided, as suggested in Reference 32.

B. Coannular Shock Noise Prediction

The spectral shock cell noise prediction method due to Fisher and Harper-Bourne (FHB) (Reference 33) for conic nozzles has been shown to adequately predict the shock noise from conic nozzles over a wide range of operating conditions. Therefore, the approach taken in formulating a shock noise spectral prediction method for coannular nozzles was to use the FHB method with appropriate modifications.

As the shock structure in any under- or overexpanded jet of a given cross-sectional area is a function of the operating pressure ratio and independent of the reservoir temperature, it would only seem logical to determine the effective pressure ratio of the coannular jet stream based solely on the pressure ratios of the two streams and not on the temperature at which they are operated. This is accomplished by equating the ideal thrust from an equivalent jet to the sum of the ideal thrusts produced by the inner and outer streams of a coannular plug nozzle configuration. By so doing, the effective pressure ratio $P_{\rm r}^{\rm eff}$ of the equivalent jet can be determined and hence the effective (isentropic) Mach number. The effective shock strength parameter is then defined as:

$$\beta^{\text{eff}} = \sqrt{(M_{j}^{\text{eff}}^{2} - 1)}$$
 (23)

where

$$M_i^{eff^2} = 2/\gamma - 1 \left[(P_r^{eff})^{\gamma - 1/\gamma} - 1 \right]$$
 (24)

and
$$P_r^{eff} = P_r^0 + P_r^i A_r/1 + A_r$$

The equation used in the FHB method for the average shock cell length \boldsymbol{L} is

$$L = KBD \tag{25}$$

where

D = Exit diameter of the nozzles

 β = Shock strength parameter as defined by Equation 23

K = 1.1 for conic nozzles

For coannular jets, based on examining the data (Reference 17), the constant K was determined to be a function of $\beta^{\rm eff}$ and is given by

$$\kappa = 0.48 \text{ Beff} + 0.54$$
 (26)

Since the turbulence eddies convecting past the shock cells generate the broadband noise, the convection velocity associated with these eddies must be identified. The mass-averaged total temperature T_T^{mix} , and the effective jet Mach number, M_j^{eff} , are used to determine the effective jet velocity which is substituted in the FHB method in order to predict the peak frequency for the shock-associated broadband noise.

With these changes, the FHB method was used to predict shock noise for coannular jets. On comparing the predicted spectra with that measured, the shape was found to be satisfactory but the level was 6 dB too high. Since this was consistently observed for several operating conditions, the level was lowered by 6 dB and viewed as the noise from two shock cells instead of eight.

1. Extension of Shock Noise Prediction to Flight

The only difference between the static and flight shock noise is the dynamic effect and the Doppler shift. The dynamic effect is given by

$$(SPL_F - SPL_S) = 40 \log_{10} (1 + M_{ac} Cos \theta_I)$$
 (27)

which amplifies the shock noise in the front quadrant while it mitigates it in the aft quadrant. The Doppler shift is given by

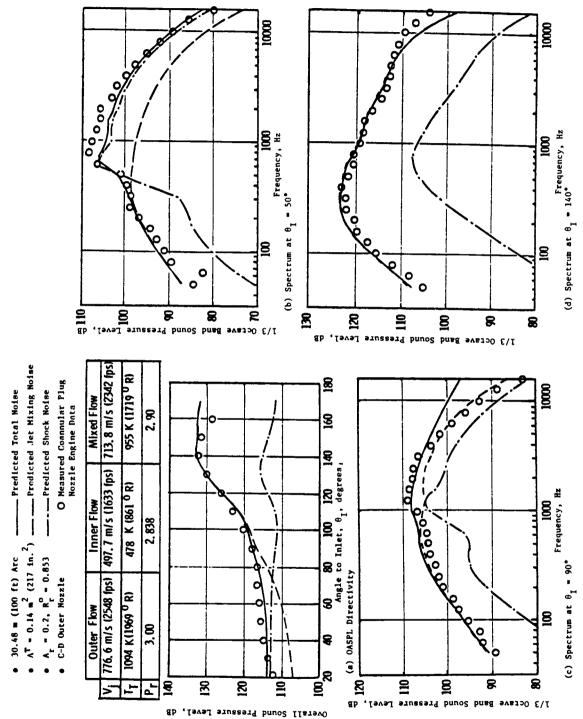
$$f_F = f_S/(1 + M_{ac} \cos \theta_I)$$
 (28)

With these two changes, the flight peak noise spectra are predicted.

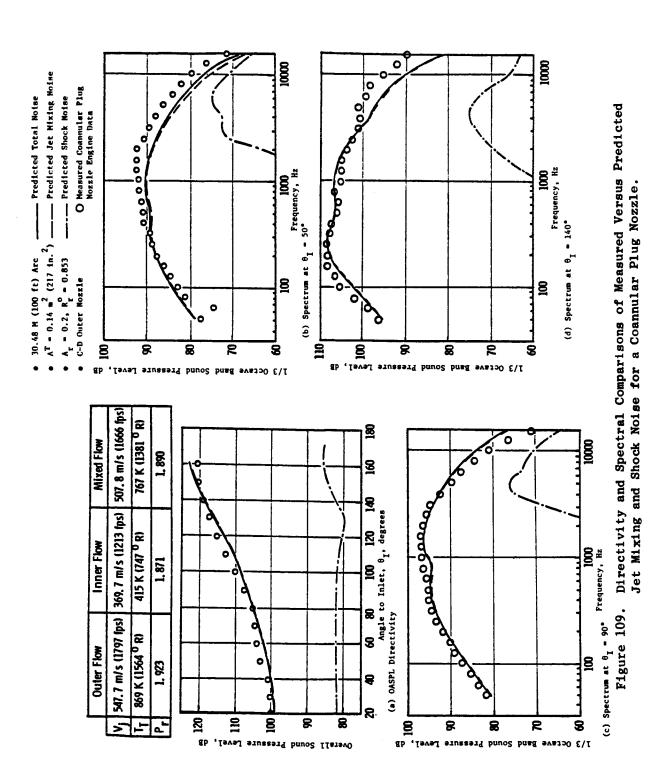
The coannular jet mixing noise and shock noise routines have been programmed, and the total spectra can be predicted for varying operation conditions and nozzle geometries.

5.3.3 Comparison of Data and Predictions

Comparison of the measured OASPL directivity and the 1/3-octave sound pressure spectra for a coannular plug nozzle configuration with $A_{\rm T}=0.2$ and $R_{\rm T}=0.853$ and mounted on a VCE engine (Reference 9) operating in the inverted velocity profile mode is shown in Figures 108 and 109. A good agreement is observed in both the noise level and in the shape of the directivity and sound pressure spectra. Additional verification is provided by Figures 110 through 112 that compare the measured and predicted variations of normalized PNLmax with 10 log $V_{\rm j}^{\rm mix}$ for coannular nozzles of different area ratios and radius ratios.



Directivity and Spectral Comparisons of Measured Versus Predicted Jet Mixing and Shock Noise for a Coannular Plug Nozzle. Figure 108.



- $A^{T} = 0.907 \text{ m}^{2} (1400 \text{ in.}^{2})$
- 731.5 m (2400 ft) Sideline
- C-D Outer Nozzle
- F_{ref} = 22,819 N (5130 lb)

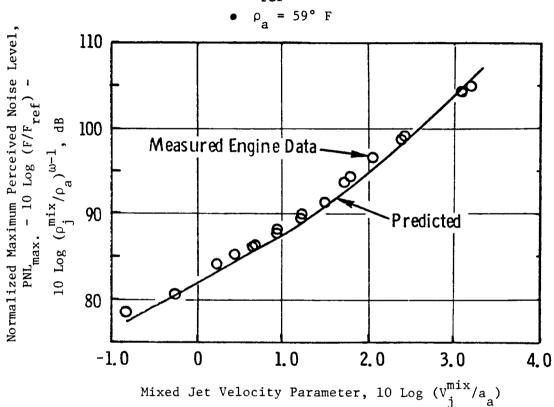


Figure 110. Comparison of Measured Versus Predicted Variation of Normalized PNL for a Coannular Nozzle, Area Ratio = 0.2, Outer Radius Ratio = 0.853.

- $A^T = 0.907 \text{ m}^2 (1400 \text{ in.}^2)$
- 731.5 m (2400 ft) Sideline
- C-D Outer Nozzle
- $F_{ref} = 22,819 \text{ N (5130 lb)}$
- $\rho_a = 59^{\circ} F$

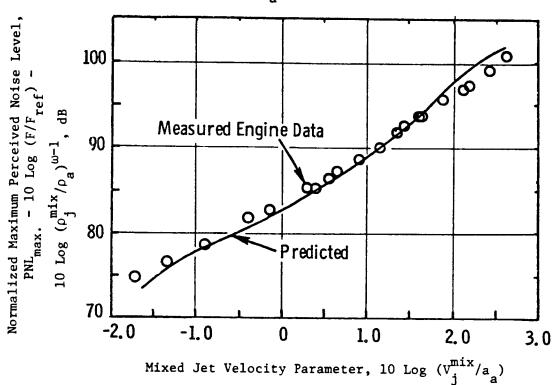
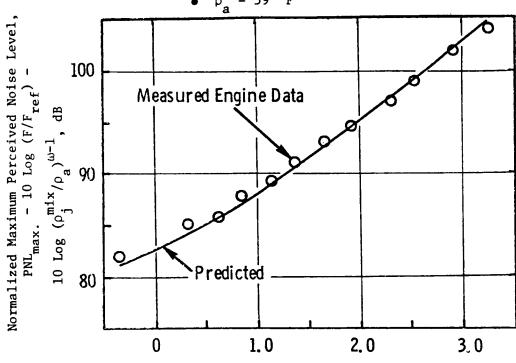


Figure 111. Comparison of Measured Versus Predicted Variation of Normalized PNL for a Coannular Nozzle, Area Ratio = 0.475, Outer Radius Ratio = 0.853.

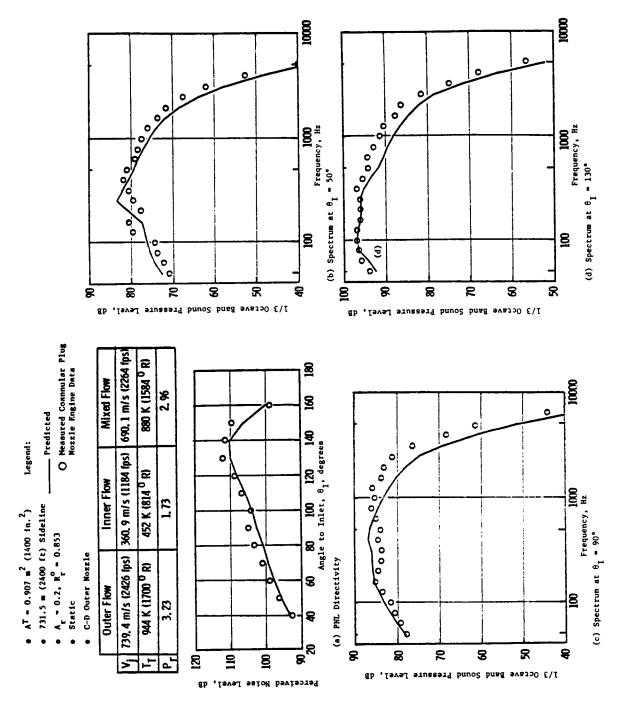
- $A^{T} = 0.907 \text{ m}^{2} (1400 \text{ in.}^{2})$
- 731.5 m (2400 ft) Sideline
- C-D Outer Nozzle
- F_{ref} = 22819 N (5130 lb)
- $\rho_a = 59^{\circ} F$



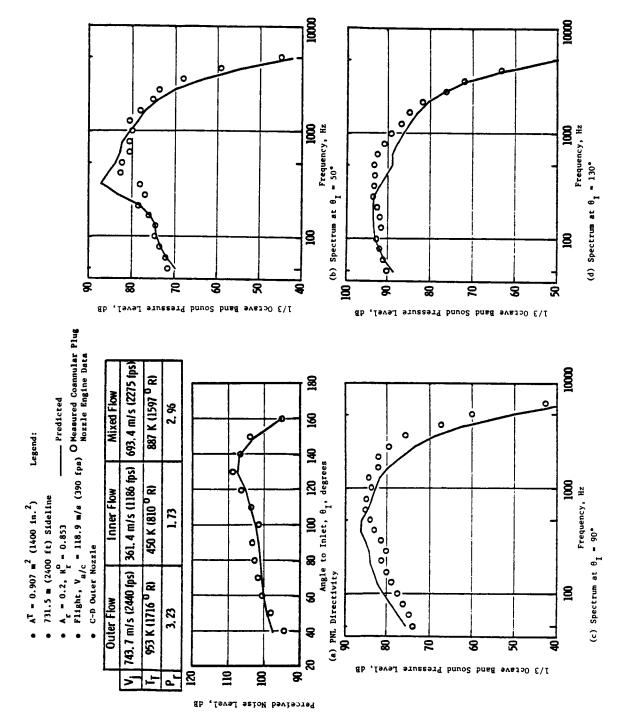
Mixed Jet Velocity Parameter, 10 Log (V_{j}^{mix}/a_{a})

Figure 112. Comparison of Measured Versus Predicted Variation of Normalized PNL for a Coannular Nozzle, Area Ratio = 0.2, Outer Radius Ratio = 0.875.

To assess the in-flight prediction method, the free-jet data from coannular plug nozzles were flight transformed (Reference 28) and then compared with the predictions. This is shown in Figures 113 and 114. A good agreement is observed in both the spectrum shape and directivity. It must be indicated here that the data used for comparisons with the predictions were from engine and scale model tests, and they were not a part of the data base that was used during the development of the prediction method.



Directivity and Spectral Comparisons of Measured Versus Predicted Jet Mixing and Shock Noise for a Coannular Plug Nozzle (Static). Figure 113.



Directivity and Spectral Comparisons of Measured Versus Predicted Jet Mixing and Shock Noise for a Coannular Plug Nozzle (Flight). Figure 114.

6.0 CONCLUSIONS AND RECOMMENDATIONS

6.1 CONCLUSIONS

All six high-radius-ratio coannular plug nozzles, along with a reference conical nozzle, were successfully tested in a simulated flight environment. The nozzle models tested are candidate exhaust nozzle configurations for General Electric designs for VCE and AST applications. The nozzle geometric variables included: an outer nozzle radius ratio variation (0.853 and 0.902), an inner-to-outer nozzle area ratio variation (0.2 and 0.53), nozzles with and without struts, and a first simple attempt at a convergent-divergent termination on the outer stream nozzle for further shock noise control. Most of the tests were conducted at elevated exhaust nozzle temperatures (up to 1760° R) and at high nozzle pressure ratios (up to 3.6).

The significant results are:

- In a simulated flight environment, the unsuppressed coannular plug nozzle maintained its general favorable noise reduction features. At typical takeoff sideline engine operation, 5 PNdB jet noise reduction and 6 PNdB shock noise reduction were measured relative to a reference baseline conical nozzle at the same specific thrust and effective nozzle pressure ratios.
- Outer stream radius ratio and inner-stream-to-outer-stream area ratio nozzle geometry parameters were found to influence the simulated flight acoustic signature similar to what has been observed and reported from General Electric static acoustic test results (Reference 2):
 - A higher outer stream radius ratio results in greater noise on. reduction.
 - At low specific thrust values, an increase in area ratio will increase the jet noise.
 - Inner-stream-to-outer-stream velocity ratio is an important design parameter.
- The control of the downstream coannular shock structure can have a significant effect on the forward quadrant and aft quadrant radiated noise. Up to 2.5 AEPNdB noise reduction for a coannular plug nozzle configuration was estimated.
- To obtain shock control through the use of a convergent-divergent termination, great care will be necessary in the aerodynamic contour of the annular nozzle passages.

• A unique coannular plug nozzle spectral jet mixing and shock noise prediction method was evaluated. The method is based on the physics of the flow and its noise-radiating characteristics and is the first method to model the source spectrum, eddy spectrum, and fluid shrouding in a simplistic fashion. The predicted data were found to be in good agreement with measurements over a range of operating velocities. The general methodology might be easily extended to nozzles of other complex geometries.

6.2 RECOMMENDATIONS

Based on the studies conducted during this contract effort, the following items warrant future investigations:

- Continue systematic simulated flight acoustic experiments to evolve design criterion for shock control of coannular and annular plug nozzles using convergent-divergent terminations.
- Improve evaluation of the influence of temperatures on forward quadrant coannular plug nozzle shock noise.
- Evaluate the flight acoustic spectrum of a candidate multielement, mechanically suppressed coannular plug nozzle.
- Further evolve the spectral acoustic prediction methodology of this study to other complex nozzle geometries for supersonic and subsonic exhaust nozzle applications.

7.0 NOMENCLATURE

A Cross-Sectional Area

a Radius of Free Jet

AST Advanced Supersonic Transport

C Convective Amplification Factor (see Equation 12a)

c Speed of Sound

CDR Comprehensive Data Report

Corr_R Refraction Correction

Corr_T Turbulent Absorption Correction

Corr90° Turbulent Absorption Correction at $\theta_{I} = 90^{\circ}$

dB Decibel

D Diameter

D_R Directivity Factor

D_I Directivity Factor

EPNL Effective Perceived Noise Level, EPNdB

F Thrust

Fref Reference Thrust, 5130 pounds

f Frequency

FTFSDR Flight Transformed, Full-Scale Data Reduction Computer Program

g Fluid Shielding Function (see Equation 15)

H Shielding Factor

h Annular Step Height

Hz Hertz

I_o, I_i Bessel Functions

Jo, Ji Bessel Functions

K Constant in Equation (25)

k Wave Numbers, 2π f/c

L Shock Separation Distance

l Path Length

M Mach Number

M_C Convection Mach Number

N Convective Amplification Exponent

OAPWL Overall Sound Power Level

OASPL Overall Sound Pressure Level

P Pressure

Pr Defined = P_T/P_a

PNL Perceived Noise Level

PNLN Normalized PNL, Defined as PNL - 10 log (F/F_{ref}) $(\rho/\rho_a)^{w-1}$

R Radius

RH Relative Humidity

R_r Radius Ratio

S Strouhal Number

SL Side Line

SPL Sound Pressure Level

SPLN Normalized SPL

T Temperature

U_c Convection Velocity

u Turbulent Velocity

V Ideally Expanded Velocity

VCE Variable Cycle Engine

W Weight Flow Rate

```
Axial Distance
X
               Distance Along Plug (see Figure 103)
Xs
                Geometric Acoustic Length, ka sin 0, see Figure 8
x
                Geometric Acoustic Length, ka \{\cos^2 \theta - (1 - M \cos \theta)^2\}^{1/2},
y
                see Figure 8
                Bessel Functions
Yo, Yi
                Turbulence Parameter
                Linear Regression Coefficients
a1, a2
                Shock Strength Parameter
β
                Specific Heat Ratio
                Eddy Viscosity
                Angle Measured Relative to Exhaust Centerline
                cos^{-1} [1/(1 x M)], see Figure 8
θc
                Angle Measured Relative to the Inlet Centerline
θΙ
                Plug Angles, Defined in Table III
                Absolute Viscosity Coefficient
                Density
                Standard Error of Estimate, dB
σyn
                Density Exponent
Subscripts
                Hub Dimension (defined in Table III)
1
2
                Tip Dimension (defined in Table III)
                Aircraft
ac
                Initial
i
```

Ambient Conditions

Basic

Basic (Directivity)

С

Critical

E

East

eff

Effective

F

Flight

eq

Equivalent

fj

Free Jet

hyd

Hydraulic

max

Maximum

p

Peak

j

Based on Ideal Jet Conditions

r

Ratio

S

Static

T

Stagnation Condition

W

West

60°

Evaluated at $\theta_{\rm I} = 60^{\circ}$

θ

Value at Angle $\theta_{\mathbf{I}}$

Superscripts

eff

Effective

HF

High Frequency

i

Inner Stream

LF

Low Frequency

0

Outer Stream

mix

Fully mixed Conditions

T

Based on Total Area

_

Mean Quantity

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APPENDIX A

AERODYNAMIC TEST CONDITIONS FOR THE ACOUSTIC TESTS

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	mix J /SEC	33 4 4 5 5 5 5 5 5 5 5 5 5 5 5 5 5 5 5 5	dB 140	85.8 889.9 883.9 883.9 895.6 895.6 895.6 895.6 895.6 895.6 895.6 895.6 895.6 895.6 895.6 895.6 895.6 895.6
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	i W /SEC	667.2 667.2 667.2 667.2 667.2 668.8	2400 VE TO 110	88 99 99 99 99 99 99 99 99 99 99 99 99 9
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	t V j FT/SE	888 1388 1388 1388 138 138 138 138 138 1	(FULL INGLE R 70	00 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0
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Ð	v J FT/SEC	2 4 2 4 2 4 2 4 2 4 2 4 2 4 2 4 2 4 2 4	A B	- 1 - 1 - 1 - 1 - 1 - 1 - 1 - 1 - 1 - 1
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OTAL =	mix J /SEC	22246 22253 22253 22378 22378 22378 22186 22186 22278 22278 22278 22283 22283 22283 22283	dB 140	110.5 1110.5 1110.5 100.0 100.0 100.0 100.0 100.0 100.0 100.0 100.0 100.0 100.0 100.0
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SIZ	T T DEG	<u> </u>	DE L O DE	
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18.05	1 W B/SEC	995.22 995.23 995.23 995.23 995.23 955.23	2400 VE TO	105.2 105.2 110.6 110.6 100.5 100.3 100.3 100.4 100.6
. = H31	EC L	699998-1838-1639 69999999-1838-1639 69999999-1838-1838	SIZE, RELATI 90	106.9 110.2 110.2 110.2 110.3 110.3 110.4 103.4 104.5 104.5 104.5 104.5 104.5 104.5 104.5 104.5 105.8
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NNER =	r P	33.22 33.22 33.22 22.28 22.28 22.28 22.28 22.28 23.22 25.07	S	998 987 988 989 989 989 989 989 989 989
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ARE	T T DEG R	1720 1714 1697 1693 1713 1723 1709 1709 1709 1709 1700 1700 1700 1700	HH *	667 667 667 667 667 667 667 667 667 667
8	o _ r	23.30 23.30 23.30 23.30 23.30 23.30 23.30 23.30	P amb PSIA	
- MODEL	v ac T/SEC	296 396 396 295 396 297 396 396 396 396 396	T amb DEG R	5514.4 5504.0 5506.8 5508.1 5508.1 5508.1 5508.1 5508.1 5518.2 5518.3 5518.4 5518.4 5518.4 5518.4
NOZZLE -	TEST POINT F1	201 202 203 204 205 207 208 210 211 211 215 216 216 219	TEST POINT	202 203 203 204 205 205 207 203 203 213 213 214 215 215 216

	i v /v j j	35815 0.53 41040 0.62	OAPWL dB	173.7 181.2
	L B	35815 41040	0.0	
	mix V f FT/SEC	1936 3 1999 4	dB 140	97.4 106.6
		1553 1520	LINE), DEGREES 130	100.9 107.4
	P T T 1	1112 95.3 2.18 1553 1341 135.7 2.35 1520	PNL (FULL SIZE, 2400 FT SIDE LINE), dB ANGLE RELATIVE TO INLET, DEGREES 50 70 90 110 120 130 1	99.8
	1 W F	95.3	2400 I	97.4 102.5
	1 J SEC LE	112 341 1	L SIZE, RELATI 90	97.7 100.6
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	1 1 T V DEG R FT/SEC	8 12 756	PNI 50	12.4
	7 <u>1</u> 1	1695 2094 499.7 1.61 1718 2170 524.6 2.16		
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NTIN	T T EG R	1695 1718		
, 2 7	ه د		P amb PSIA	
NOZZLE - HODEL 2 CONTINUED	V ac FT/SEC	396 2.31 0 2.44	T amb DEG R	509.7 506.6
NOZZLE	TEST POINT	221 222	TEST	221 222

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1400.00]	F LB	61543 621605 62172 74903 749312 743812 52864 53011 33788 43586 43588 43586 43588 43586 43586 43586	0 M
TOTAL =	mix V j FT/SEC	2244 2254 2256 2370 2377 2377 2186 2186 2187 2210 2210 2109 2165	£ 1001111000000000000000000000000000000
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18.05;	i W B/SEC	200.3 201.2 200.2 200.2 198.4 198.4 145.1 145.1 138.7 138.7 138.7 96.1 96.1	2 TO 1 10 1 10 1 10 1 10 1 10 1 10 1 10 1 1
OUTER =	v J /SEC L	1628 1628 1628 1620 1620 1640 1254 1256 1256 1256 1419 1419	L SIZE RELAT 90 106.2 106.2 108.6 108.6 108.6 108.7 102.3 101.7 101.3 101.3
.50	I T EG R FT	784 7786 7757 7757 7757 7757 7751 7751 7751 767 767	
NNER = 3	i P T F DE	22. 22. 22. 23. 24. 24. 27. 29. 20. 20. 20. 20. 20. 20. 20. 20. 20. 20	02 01 01 01 02 02 02 02 02 02 02 02 02 02 02 02 02
ZE - IN	W /SEC	581.0 582.0 582.0 5862.0 3366.8 337.8 529.7 529.7 529.7 529.3 11.2 529.3 11.2 520.3 11.2 520.3 11.2 520.3 11.2 520.3 11.2 520.3 11.2 520.3 11.2 520.3 11.2 520.3 11.2 520.3 11.2 520.3 11.2 520.3 520.	L B B B B B B B B B B B B B B B B B B B
MODEL SI	EC LB	22 2 2 2 2 2 2 2 2 2 2 2 2 2 2 2 2 2 2	LVM 111 111 111 111 111 111 111 111 111 1
REA [MC	V J F FT/S		Z D
A	T T DEG 1	17 19 17 17 17 17 1691 1691 17 15 17 15 17 15 17 16 17	
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- MODEL	V ac FT/SEC	289 289 289 385 385 385 385 386 347	ED 00-000000/-0/-00-0-
NOZZLE	TEST POINT	301 3002 3002 3005 3005 3005 3005 3005 3005	TEST 301 302 303 303 304 305 306 307 308 308 309 301 311 311 311 311 311 311 311

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1400.00]	ล 8	34344 34182 35858 38694 38124 40492 43763 43763	V O	
TOTAL =	mix V J FT/SEC	2104 2079 2028 2028 1909 1915 2078 2107 2107	dB 140	107.8 103.0 101.3 104.9 99.3 107.5 108.7
SIZE - TO	× «	1652 1601 1782 1782 1620 1620 1521 1589	LINE), DEGREES 130	109 1 105 6 1 105 6 1 105 6 1 105 6 1 105 105 105 105 105 105 105 105 105 1
FULL SI	mix mix r T r DEG	2.2.2.2.2.2.2.2.2.2.2.2.2.2.2.2.2.2.2.) FT SIDE INLET, D 120	104.7 102.6 102.6 102.5 104.7 103.1 105.0
18.05; 1	EC F	~3 B M D T B M B B -	3 H H	102.0 100.7 100.2 100.3 100.3 102.6 102.9
11	0	425 464 7754 742 10957 10957 1414 1626 1626 1625	SIZE, RELATIV 90	4.00013 4.00013 10012 1010 1010 1010 1010 1010 101
, outer	t V J FT/SEC		FE NG 1	00000000000000000000000000000000000000
= 9.50	T T DEG R	980 1026 1000 968 746 762 1313 1258 1258	PNL A 50	99999999999999999999999999999999999999
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SIZE -	o W LB/SEC	2667.1 2667.1 2662.1 2662.1 2666.0 2666.3	LVM LI	882 4 4 5 5 5 6 8 8 8 8 8 8 8 8 8 8 8 8 8 8 8 8
[MODEL	o V j FT/SEC	2313 2285 2313 2303 2297 2297 2310 2318	NF L	0-wwoomewww
AREA [0 H	1736 1689 1736 1730 1727 1727 1728 1742	HH **	2043 2043 2043 2043 2043 2043 2043 2043
=	o T. r DEG	22.78 22.78 22.78 11.22.78 12.75 11.22.78 11.22.78	P amb PSIA	44444444444444444444444444444444444444
MODEL	V ac FT/SEC	3870 3870 3870 3870 3870	T amb DEG R	49993.2 4993.2 4992.7 4993.7 4993.7
NOZZLE -	TEST POINT FT	100849377 100849377 1008497 10084	TEST	-008498797 -008498797 -008498797
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1400.0	F E3	62758 362718 362318 36851 362165 36471 333997 333997 295543 295543 29568 406643 406643 42885 432885	V 0	811111111111111111111111111111111111111
OTAL =	mix V j FT/SEC	2064 2082 1728 1726 2038 2038 2006 1684 1671 1684 1895 2062 2079 2087	dB 140	4 10 10 10 10 10 10 10 10 10 10 10 10 10
SIZE - T	mix T EG R	1283 1305 1305 1305 1305 1305 1305 1418 1428 1418 1501 1501 1501 1583	DE C	7.5.601 7.501 7.501 7.501 7.501 7.501 7.50
FULL S	P T T O	20000000000000000000000000000000000000	FT SIDE INLET, 120	4.000 4.
11.22 ;	i ₩ B/SEC	433.6 423.6 425.0 291.9 777.4 50.1 187.3 180.5 161.6 161.6 201.8	, 2400 IVE TO	106.1 108.0 98.0 96.0 96.2 102.2 100.7 96.6 96.7 96.7 99.7 102.1 102.3
OUTER =	v J /SEC L	1615 1615 1242 1242 1242 1242 1242 1409 1409 1616 1616	LL SIZE E RELAT	96.99 104.99 96.99 96.59 96.57
.85	1 T EG R FT	761 676 676 685 685 687 709 609 747 719 719 719 726 290 2913	PNL (FU ANGL 70	001 003 003 003 003 003 003 003 003 003
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ZE - IN	w /SEC	23.25.33.25.33.35.55.5	LBM	10000000000000000000000000000000000000
DEL SI	SEC LB	######################################	LVM	2.75 2.75 2.75 2.75 2.75 2.39 2.75 2.75 2.75 2.75 2.75 2.75 2.75 2.75
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	o T r DEG	21 169 30 168 30 168 20 173 229 170 778 173 80 173 775 169 775 169 775 169 775 169 775 169 775 169 775 169 775 169	P F amb PSIA	00000000000000000000000000000000000000
MODEL 6	ac EC	0.09 0.09 0.09 0.09 0.09 0.09 0.09 0.09	G B R R	00000000000000000000000000000000000000
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00.00]	د. <u>۱</u>	3082 3082 3082 3082 3082 3082 3082 3082	OAP	182 182 174 174 175 175 177 178 178 178 178 178 178 178 178 178
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I [MODEL	v J FT/SEC	24331 24331 2009 2309 2309 2009 2009 2009 2009 2009	N P	
AREA	T T EG R	1741 1741 1730 1730 1730 1714 1715 1717 1713 1713 1713 1713	RH ~	00000000000000000000000000000000000000
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- MODEL	V ac T/SEC	386 398 399 399 399 399 399 399 399 399 399	T amb DEG R	288.33 288.33 288.33 288.33 288.33 288.33 288.33 288.33 288.33 288.33 288.33
NOZZLE -	TEST POINT F	7001 7002 7009 7009 7009 7010 7011 7012 7018 7019	TEST	7001 7002 7003 7004 7005 7006 7010 7011 7013 7014 7016 7019

NOZZLE - MODEL 7 CONTINUED

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mix T EGR F	1367 1286 1280 1280 1280 1280 1280 1586 1586 1586	LINE), 130 130 130 104.7 106.3 106.3 107.6 111.0 1113.4 1112.6 110.6 110.6 110.6 110.6 110.6 110.6 110.6 110.6 110.6 110.7
r J	www.v-r-ww.v.v.v.w.v.w.v.w.w.w. www.u-r-c-o	NLET DE 120 D 120
1 P		E 100 F 100
EC LB/	200 200 200 200 200 200 200 200 200 200	SIZE, 90 100000000000000000000000000000000000
v V J FT/SH	121 172 173 173 173 173 173 173 173 173 173 173	ANGLE R ANGLE R 70 995.2 995.3 97.6 86.5 87.6 87.6 905.2 111111111111111111111111111111111111
T T DEG R	8845 88565 1065 1007 1007 1007 1007 1358 1355 1355 1355 1355 1355	PNL 0.042 = 4.00 = 0.0
ન 다	2.27 2.27 2.27 2.27 2.27 2.27 2.27 2.27	75872 2000 000 000 000 000 000 000 000 000
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o j /SEC 1	2076 2000 2000 2010 2010 2010 11694 11694 11694 11694 2248 2248 2248 2284 2282 2284 2385 2385	
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***	3/SEC	,		30.1	31.2	30.5	30.1	77.9	, c	20.0		05.2	49.3	48.6	10.1	15.5	-20	185.5	82.3	2400	To	110	03.	'n.		0.50	5.	æ	ه ف	٠, د د		. 6	•	٠,	0 u		107.5	۲,	•
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OAPWL

E LINE), DEGREES 130

SIDE ET, I

332490 33058 33058 16290 16748 50122 50122 50217 4984 14984 14988 14984

mix V j FT/SEC

mix T DEG

75.6 175.6 181.9 181.9 181.9 182.2 182.2 181.6 181

NOZZLE - MODEL 7	ODEL	. 7 co	ONTINUED	JED												
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APPENDIX B

AERODYNAMIC TEST CONDITIONS FOR THE LASER VELOCIMETER TESTS

Table B-I. Aerodynamic Test Matrix for Laser Velocimeter Tests.

	Ī		Outer												
	Test					Inner				Mixed					
Node 1	Point	Va/c ft/sec	7°	ft/sec	TT R	. I	yo lbm/sec	Př	yi ft/sec	. 17	.Tk	yi lbe/sec	FE/sec	.TH	₩ ¹ /₩°
1	101	0	3.17	2416	1706	1263	10.52	3.19	1636	790	567	2.16	2283	1146	0.677
	101A	0	1.00	0	592	592	0	3.21	1635	785	563	2.18	1635	563	
	103	400	3.17	24.27	1722	1274	10.47	3.19	1633	787	565	2.17	2291	1155	0.673
	113A		1.00	0	571	571	0	1.62	1104	784	683	1.08	1104	683	_
	116	0	2.43	2175	1735	1361	8.00	2.06	1348	801	650	1.40	2051	1276	0.620
	1194	0	2.26	2092	1733	1406	7.44	1.02	262	1013	1008	0.18	2050	1403	0.125
	1505	0	1.45	1447	1720	1367	4.53	1.45	898	666	599	1.01	1348	1393	0.621
	1506	400	1.45	1453	1733	1579	4.51	1.43	899	693	626	0.97	1355	1413	0.619
1A	101A	0	3.19	2A53	1722	1272	10.54	3.22	1648	789	565	3.09	2270	1114	0.672
	116A	0	2.45	2184	1721	1366	7.96	2.20	1338	735	587	2.15	2003	1511	0.613
2	201	0	3.17	2A 35	1733	1283	10.44	3.24	1629	774	553	3.10	2250	1118	0.669
	203	400	3.18	2425	1716	1268	10.52	3.20	1630	782	561	3.04	2247	1112	0.672
	204	0	3.79	2551	1686	1186	12.65	3.23	1638	784	561	3.06	2372	1069	0.642
	204	400	3.76	2559	1707	1203	12.48	3.24	1635	780	557	3.08	2376	1080	0.639
	210	0	2.28	2088	1710	1384	7.56	2.09	1264	700	567	2.10	1909	1210	0.605
	212	400	2.28	2084	1704	1379	7.57	2.11	1265	693	560	2.13	1904	1203	0.607
	213	0	3.15	24.30	1734	1286	10.37	1.58	1097	817	71.7	1.43	2268	1229	0.451
	215	400	3.18	ZA 29	1722	1273	10.50	1.61	1104	789	696	1.49	2265	1213	.0.455
	21.6	0	3.79	2551	1689	1187	12.65	1.59	1098	806	706	1.45	2402	1151	0.430
	219	0	2.27	2085	1713	1388	7.52	1.59	109 Z	800	701	1.46	1924	1283	0.524
	221	400	2.29	2089	1704	1377	7.61	1.61	1094	783	683	1.50	1925	1270	0.524
	222	•	2.42	2165	1727	1376	7.98	2.18	1324	731	585	2.14	1987	1213	0.612
	222A	0	2.45	21.54	1713	1360	8.12	2.18	1334	739	591	2.13	2000	1510	0.611
3	301	0	3.17	2434	1732	1282	10.44	3.22	1634	782	560	3.10	2250	1119	0.671
	303	400	3.17	2433	1731	1281	10.44	3.21	1634	784	562	3.09	2251	1119	0.671
5	513	0	3.14	24.28	1735	1288	11.67		_		-	_	2428	1288	
	515	400	3.19	2431	1720	1270	11.90	-	_	-	-		2431	1270	
•	30091	0	2.26	2088	1726	1401	4.64	1.37	795	611	556	2.26	1664	1145	0.361
	3011	400	2.23	2070	1722	1402	4.58	1.39	796	587	534	2.37	1636	1127	0.385
- 1	3015R	۰	2.78	2292	1704	1306	5.74	1.63	1430	1317	1156	1.98	2073	1278	0.628
	3016	300	2.73	2290	1727	1332	5.60	1.63	1426	1294	1136	2.00	2063	1290	0.623
l	3017	400	2.74	2293	1726	t329	5.62	1.63	1431	1303	1144	1.99	2068	1291	0.624
	3018	٥	2.70	2285	1736	1343	5.52	1.96	1630	1259	1049	2.47	2082	1258	0.713
	3020	400	2.70	2281	1730	1338	5.53	1.95	1622	1255	1047	2.46	2078	1254	0.711
7	7009	•	2.25	2085	1730	1406	7.46	1.48	868	592	529	2.59	1772	1195	0.416
	7011	400	2.25	2090	1737	1412	7.45	1.48	869	593	530	2.59	1775	1199	0.416
l	7015	۰	2.73	2295	1735	1338	9.04	1.43	1406	1036	874	2.54	2100	1244	0.613
ļ	7016	300	2.71	2287	1734	1340	8.98	1.83	1411	1043	861	2.53	2094	1247	0.617
	7017	400	2.70	2292	1746	1351	8.91	1.83	1410	1042	880	2.53	2097	1255	0.615
ŀ	7018	0	2.67	2269	1729	1341	8.86	2.73	1643	901	676	4.08	2072	1134	0.724
	7019	300	2.68	2267	1721	1333	8.91	2.74	1650	905	679	4.08	2073	1130	0.728
	7020	400	2.67	2274	1736	1347	8.84	2.78	1644	888	663	4.18	2071	1129	0.723

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16. Abstract																									
This report summarizes the experimental and analytical results of a scale-model simulated flight acoustic exploratory investigation of high-radius-ratio coannular plug nozzles with inverted velocity and temperature profiles. In all, six coannular plug nozzle configurations and a baseline convergent conical nozzle were tested for simulated flight acoustic evaluation in General Electric's Anechoic Free-Jet Acoustic Test Facility. The nozzles were tested over a range of test																									
									high speed aircraft. The or	conditions that are typical of a Variable Cycle Engine for application to advance high speed aircraft. The outer stream radius ratio for most of the configurations															
									was 0.853, and the inner-stream-to-outer-stream area ratio was tested in the range																
									of 0.2 to 0.54. Other varia	of 0.2 to 0.54. Other variables investigated were the influence of bypass struts,															
									a simple noncontoured convergent-divergent outer stream nozzle for forward quadrant																
									inner-stream-to-outen stream	shock noise control, and the effects of varying outer stream radius ratios and															
inner-stream-to-outer-stream velocity ratios on the flight noise signatures of the nozzles. It was found that in simulated flight, the high-radius-ratio coannular plug nozzles maintain their jet noise and shock noise reduction features previously observed in static testing. The presence of nozzle bypass struts will not significantly affect the acoustic noise reduction features of a General Electric-type																									
								nozzle design. A unique coannular plug nozzle flight acoustic spectral prediction method was identified and found to predict the measured results quite well. Special laser velocimeter and acoustic measurements were performed which have given new insights into the jet and shock noise reduction mechanisms of coannular plug																	
																	nozzles with regard to identifying further beneficial research efforts.								
																	17. Key Words (Suggested by Author(s))	18. Distribution State							
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Supersonic jet noise reduction; Variable cycle engine; Acoustic flight noise simu-																									
								lation; Laser velocimeter me	easurements																

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